

Ten on Ten Student Notebook with

38 प्रश्न: 28: कसविशेषरूपमा हाम्रो देशमा कतिपय दिनमा तः गर्जोस्मा 1

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To find P.I

① $y'' + y = x \Rightarrow (D^2 + 1)y = x \Rightarrow y = CF + PI$

② P.I = $\frac{1}{(D^2 + 1)} x = (D^2 + 1)^{-1} x = (1 - D^2 + D^4 - \dots) x$
 $= x$

③ $\frac{1}{(D^2 + 9)} \sin ax = \frac{-x \cos ax}{2a}$

for polar form $\frac{\partial F}{\partial r} - \frac{d}{d\theta} \left(\frac{\partial F}{\partial \theta} \right) = 0$

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Syllabus

UNIT-I
Laplace transforms and their properties,
Laplace transform of derivatives and integrals,
Inverse Laplace transformation & Convolution
theorem. Impulsive function, Application of,
Laplace transform in solving linear differential
equations with constant coefficients and
with variable coefficients.
Fourier transforms: Definition & Properties
of Fourier sine, cosine and complex transforms
development of the Fourier integral from the
Fourier series - simple Applications.
Application to the solution of simple
boundary value problems by Laplace and
Fourier transforms.

UNIT-II
(84) Calculus of Variations - Functional, Euler's
equation and application, Geodesics, Geodesics
on a sphere of radius 'a', isoperimetric (90)
problems, Variational problem with several
variables, approximate solⁿ of boundary (114)
value problems by Rayleigh - Ritz method,
Hamilton's principle, Lagrange equation and
applications

UNIT-III
(115) Integral equation of the first and second
kind of Fredholm and Volterra type, Solving
with separable kernels, characteristic no.s
and eigenfunctions, resolvent kernel.

Book: (1) Applied Mathematics for Engineers and Physicists
(2) by L.A. Pipe

Laplace Transform
:- Let $F(t)$ be a fⁿ defined for all the
value of $t (t \geq 0)$. then Laplace transform
of $F(t)$ is a fⁿ of a new variable
and given by -

$$L\{F(t), s\} = L\{F(t)\} = f(s) = F(s) = \int_0^{\infty} e^{-st} f(t) dt, s > 0$$

The Laplace transform $F(t)$ is said to
exist if the integral (1) converges for
some value of s , otherwise it doesn't
exist.

Remark:

- (1) L is called Laplace transform operator.
- (2) The operator of multiplying $F(t)$ by e^{-st}
and integrating b/n the limit 0 to ∞ .

Eg. $f(t) = 1$ - $L(1) = 1/s$

$$\left. \begin{aligned} L(e^{at}) &= \frac{1}{s-a} \\ L(e^{-at}) &= \frac{1}{s+a} \end{aligned} \right\} \text{H.W}$$

(i) $L(1) = \int_0^{\infty} e^{-st} (1) dt = \int_0^{\infty} e^{-st} dt$

$$= \left(\frac{e^{-st}}{-s} \right)_0^{\infty} = \frac{e^{-\infty}}{-s} - \frac{e^0}{-s}$$

$$= 0 + \frac{1}{s} = \frac{1}{s} \text{ Ans}$$

(ii) $L(e^{at}) = \int_0^{\infty} e^{-st} e^{at} dt$

$$= \int_0^{\infty} e^{-(s-a)t} dt = \left[\frac{e^{-(s-a)t}}{-(s-a)} \right]_0^{\infty}$$

$$= \frac{e^{-\infty}}{-(s-a)} - \frac{e^0}{-(s-a)} = 0 + \frac{1}{s-a} = \frac{1}{s-a}$$

(iii) $L(e^{-at}) = \int_0^{\infty} e^{-st} e^{-at} dt$

$$= \int_0^{\infty} e^{-(s+a)t} dt = \left[\frac{e^{-(s+a)t}}{-(s+a)} \right]_0^{\infty}$$

$$= \frac{e^{-\infty}}{-(s+a)} - \frac{e^0}{-(s+a)} = 0 + \frac{1}{s+a}$$

$= \frac{1}{s+a}$ Ans

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Properties of Laplace transform:

(i) Linearity Property of Laplace transform:
we know that by defⁿ of L.T -

$$L(f(t)) = \int_0^{\infty} e^{-st} f(t) dt$$

$$L(a_1 f_1(t) + a_2 f_2(t)) = a_1 L(f_1(t)) + a_2 L(f_2(t))$$

where a_1, a_2 are constant.

Proof: $L(a_1 f_1(t) + a_2 f_2(t)) = \int_0^{\infty} [a_1 f_1(t) + a_2 f_2(t)] e^{-st} dt$

$$= a_1 \int_0^{\infty} e^{-st} f_1(t) dt + a_2 \int_0^{\infty} e^{-st} f_2(t) dt$$

$$= a_1 L(f_1(t)) + a_2 L(f_2(t))$$

(2) The first shifting theorem or first transformation:-

statement: If $L[f(t)] = f(s)$, then $L[e^{at} f(t)] = f(s-a)$ where a is real or complex const

Proof: $\because L[f(t)] = f(s) = \int_0^{\infty} e^{-st} f(t) dt$

$$L[e^{at} f(t)] = \int_0^{\infty} e^{-st} e^{at} f(t) dt$$

$$= \int_0^{\infty} e^{-(s-a)t} f(t) dt = \int_0^{\infty} e^{-rt} f(t) dt$$

$= f(r) = f(s-a)$ (where $r = s-a$)

(3) Laplace transform of some element function:-

To find L.T. of function $F(t) = t^n$ where n be any real no. > -1

Proof:- Given that $L[F(t)] = f(s)$ - (1)
by definition of Laplace transform -

$$L[f(t)] = \int_0^{\infty} e^{-st} f(t) dt$$

$$L[t^n] = \int_0^{\infty} e^{-st} t^n dt \quad (\text{putting } y=st)$$

$$y=st \Rightarrow \frac{dy}{dt} = s \Rightarrow dt = \frac{dy}{s}$$

Limit: No change

$$L[t^n] = \int_0^{\infty} e^{-y} \left(\frac{y}{s}\right)^n \frac{dy}{s}$$

$$= \frac{1}{s^{n+1}} \int_0^{\infty} e^{-y} y^n dy$$

$\int_0^{\infty} e^{-y} y^n dy$ Gamma f^n

$$L(t^n) = \frac{n!}{s^{n+1}}$$

$$L[e^{at} t^n] = \frac{n!}{(s-a)^{n+1}} \quad (\text{use first shifting theorem})$$

Since the integral is convergent if $s > a$, and divergent if $s < a$. we must take $s > a$ for the L.T

Problem $f(t) = \sin at$

①

$$L[\sin at] = \int_0^{\infty} e^{-st} \sin at dt$$

We know that $e^{iat} = \cos at + i \sin at$
 $e^{-iat} = \cos at - i \sin at$

$$\begin{aligned} \therefore \sin at &= \frac{e^{iat} - e^{-iat}}{2i} \\ &= \int_0^{\infty} e^{-st} \left(\frac{e^{iat} - e^{-iat}}{2i} \right) dt \\ &= \frac{1}{2i} \left[\int_0^{\infty} e^{-(s-ia)t} dt - \int_0^{\infty} e^{-(s+ia)t} dt \right] \end{aligned}$$

$$= \frac{1}{2i} \left[\frac{-e^{-(s-ia)t}}{(s-ia)} + \frac{e^{-(s+ia)t}}{(s+ia)} \right]_0^{\infty}$$

H.W

① $L[\sin at] = \frac{a}{s^2 + a^2}$

② $L[\cos at] = \frac{s}{s^2 + a^2}$

③ $L[\sinh at] = \frac{a}{s^2 - a^2}$

④ $L[\cosh at] = \frac{s}{s^2 - a^2}$

Solⁿ

③ $\sinh at = \frac{e^{at} - e^{-at}}{2}$

$$\begin{aligned} \therefore L[\sinh at] &= \int_0^{\infty} e^{-st} \left\{ \frac{e^{at} - e^{-at}}{2} \right\} dt \\ &= \frac{1}{2} \int_0^{\infty} e^{-(s-a)t} dt - \frac{1}{2} \int_0^{\infty} e^{-(s+a)t} dt \\ &= \frac{1}{2} \left[\frac{-e^{-(s-a)t}}{(s-a)} \right]_0^{\infty} - \frac{1}{2} \left[\frac{-e^{-(s+a)t}}{(s+a)} \right]_0^{\infty} \\ &= \frac{1}{2} \left[0 + \frac{1}{s-a} \right] - \frac{1}{2} \left[0 + \frac{1}{s+a} \right] \\ &= \frac{1}{2} \left[\frac{1}{s-a} - \frac{1}{s+a} \right] = \frac{1}{2} \frac{s+a - s+a}{(s-a)(s+a)} \\ &= \frac{1}{2} \cdot \frac{2a}{(s-a)(s+a)} = \frac{a}{(s-a)(s+a)} \\ &= \frac{a}{s^2 - a^2} \end{aligned}$$

⑤

$$L\left[\frac{e^{at} - 1}{a}\right]$$

$$\begin{aligned} L\left[\frac{e^{at} - 1}{a}\right] &= \int_0^{\infty} e^{-st} \left\{ \frac{e^{at} - 1}{a} \right\} dt \\ &= \frac{1}{a} \int_0^{\infty} e^{-(s-a)t} dt - \frac{1}{a} \int_0^{\infty} e^{-st} dt \\ &= \frac{1}{a} \left[\frac{-e^{-(s-a)t}}{(s-a)} \right]_0^{\infty} - \frac{1}{a} \left[\frac{-e^{-st}}{s} \right]_0^{\infty} \\ &= \frac{1}{a} \left[\frac{1}{s-a} \right] - \frac{1}{a} \left[\frac{1}{s} \right] \end{aligned}$$

$L\left[\frac{e^{at} - 1}{a}\right] = \frac{1}{a} \left[\frac{1}{s-a} - \frac{1}{s} \right]$

$$\cos 2t = \cos^2 t - \sin^2 t = 2\cos^2 t - 1 = 1 - 2\sin^2 t$$

$$= \frac{1}{a} \left[\frac{1}{s-a} - \frac{1}{s} \right]$$

$$= \frac{1}{a} \left[\frac{s - st + a}{s(s-a)} \right] = \frac{1}{s(s-a)}$$

- (6) $L[\cos^2 t]$ (8) $L[\cos^2 t/2]$
- (7) $L[\sin^2 t]$ (9) $L[\sin^2 t/2]$

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[4]

Integral Transform:-

Suppose that the f^n $K(s,t)$ is a known f^n of two variable s & t . Then Integral $\int_{-\infty}^{\infty} K(s,t) F(t) dt$ — (1)

the eqn (1) is convergent then eqn (1) is called integral transform of the given f^n . It is denoted by $\bar{F}(s)$ or $T\{f(t)\}$ or $f(s)$

Thus we can write as
$$\bar{F}(s) = \int_{-\infty}^{\infty} K(s,t) F(t) dt$$

the f^n $K(s,t)$ is known as the Kernel of the transformation, has s as a parameter (real or complex) independent of t .

Remark: If we choose
$$K(s,t) = \begin{cases} 0 & ; t < 0 \\ e^{-st} & ; t > 0 \end{cases}$$

then $f(s)$ is known as Laplace transform

Prove that —

Que

$$\sin at = \frac{e^{iat} - e^{-iat}}{2i}$$

(1) $L[e^{at} t^n] = \frac{n!}{(s-a)^{n+1}}$

$L[e^{at} t^n] = \int_0^{\infty} e^{-st} e^{at} t^n dt$

(Directly by First shift part 2?)

(2) $L[e^{at} \cosh bt] = \frac{s-a}{(s-a)^2 - b^2}$

(3) $L[e^{at} \sinh bt] = \frac{b}{(s-a)^2 - b^2}$

(4) $L[e^{at} \sin bt] = \frac{b}{(s-a)^2 + b^2}$

(5) $L[e^{at} \cos bt] = \frac{s-a}{(s-a)^2 + b^2}$

Directly may be same a/c to first shifting property

(6) if $L\{f(t)\} = F(s)$. Show that

(i) $L[\sinh at f(t)] = \frac{1}{2} [F(s-a) - F(s+a)]$

(ii) $L[\cosh at f(t)] = \frac{1}{2} [F(s-a) + F(s+a)]$

(7) Evaluate —

(i) $\sinh 2t \sin 3t$ Ans: $\frac{128}{s^4 + 10s^2 + 169}$

(ii) $\cosh 3t \cos 2t$ Ans: $\frac{2s(s^2-5)}{s^4 - 10s^2 + 169}$

Prove that —

(i) $L[t \sin at] = \frac{2as}{(s^2 + a^2)^2}$

(ii) $L[t \cos at] = \frac{s^2 - a^2}{(s^2 + a^2)^2}$

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[5]

Transform of Periodic function

If $f(t)$ is a periodic function with period T i.e. $f(t) = f(t+T)$. then prove that
$$L\{f(t)\} = \int_0^T e^{-st} f(t) dt$$

By definition -

$$L[f(t)] = \int_0^{\infty} e^{-st} f(t) dt$$

$$= \int_0^T e^{-st} f(t) dt + \int_T^{2T} e^{-st} f(t) dt + \int_{2T}^{3T} e^{-st} f(t) dt + \dots$$

∴ the second integral $t = u+T$ in the third integral $t = u+2T$ and so on.

we get -

$$\int_0^T e^{-st} f(t) dt + \int_0^T e^{-s(u+T)} f(u+T) du + \int_0^T e^{-s(u+2T)} f(u+2T) du + \dots$$

$$\Rightarrow \int_0^T e^{-st} f(t) dt + e^{-sT} \int_0^T e^{-su} f(u+T) du + e^{-s2T} \int_0^T e^{-su} f(u+2T) du + \dots$$

$$\Rightarrow \int_0^T e^{-st} f(t) dt + e^{-sT} \int_0^T e^{-su} f(u) du + e^{-s2T} \int_0^T e^{-su} f(u) du + \dots$$

$$= \left(1 + e^{-sT} + e^{-s2T} + \dots \right) \int_0^T e^{-st} f(t) dt$$

$$= \frac{1}{1 - e^{-sT}} \int_0^T e^{-st} f(t) dt$$

$$= \frac{\int_0^T e^{-st} f(t) dt}{(1 - e^{-sT})}$$

Ques Find the Laplace Transform for the given fn

$$f(t) = \begin{cases} \sin \omega t & 0 < t < \pi/\omega \\ 0 & \pi/\omega < t < 2\pi/\omega \end{cases}$$

(Ans: $\frac{\omega e^{-s\pi/\omega} + \omega}{(1 - e^{-2s\pi/\omega})(s^2 + \omega^2)}$)

Solⁿ -

$$L[f(t)] = \int_0^{\infty} e^{-st} f(t) dt$$

$$= \int_0^{\pi/\omega} e^{-st} f(t) dt + \int_{\pi/\omega}^{2\pi/\omega} e^{-st} f(t) dt$$

$$= \int_0^{\pi/\omega} e^{-st} \sin \omega t dt + \int_{\pi/\omega}^{2\pi/\omega} e^{-st} \cdot 0 dt$$

type-I putting $\sin \omega t = \frac{e^{i\omega t} - e^{-i\omega t}}{2i}$

we have $\frac{1}{2i} \int_0^{\pi/\omega} e^{-st} (e^{i\omega t} - e^{-i\omega t}) dt$

$$= \frac{1}{2i} \left[\int_0^{\pi/\omega} e^{-(s-i\omega)t} dt + \int_0^{\pi/\omega} e^{-(s+i\omega)t} dt \right]$$

$$= \frac{1}{2i} \left[\left[\frac{e^{-(s-i\omega)t}}{i\omega - s} \right]_0^{\pi/\omega} - \left[\frac{e^{-(s+i\omega)t}}{s+i\omega} \right]_0^{\pi/\omega} \right]$$

$$= \frac{1}{2i} \left[\frac{e^{-s\pi/\omega} e^{i\pi} - 1}{i\omega - s} - \frac{e^{-s\pi/\omega} e^{-i\pi} - 1}{i\omega + s} + \frac{1}{s+i\omega} \right]$$

$$= \frac{1}{2i} \left[\frac{-e^{-s\pi/\omega} + 1}{i\omega - s} + \frac{e^{-s\pi/\omega} - 1}{i\omega + s} + \frac{1}{s+i\omega} \right]$$

$$= \frac{1}{2i} \left[\frac{-i\omega e^{-s\pi/\omega} + i\omega - i\omega e^{-s\pi/\omega} + i\omega + i\omega - s - i\omega}{-(\omega^2 + s^2)} \right]$$

$$= \frac{1}{2i} \left[\frac{-e^{-s\pi/\omega} [s+i\omega] - 2s}{(\omega^2 + s^2)} \right]$$

type-II

let $I = \int_0^{\pi/\omega} e^{-st} \sin \omega t dt$

$$= \sin \omega t \frac{e^{-st}}{-s} + \int \frac{\omega \cos \omega t e^{-st}}{s} dt$$

$$= \frac{\omega}{s} \left[\cos \omega t \frac{e^{-st}}{-s} - \int \frac{\omega \sin \omega t e^{-st}}{s} dt \right]$$

$$I = \sin \omega t e^{-st} - \omega \cos \omega t e^{-st} - \frac{\omega^2}{s^2} \int e^{-st} \sin \omega t dt$$

$$= \frac{-s \sin \omega t e^{-st}}{s} - \frac{\omega \cos \omega t e^{-st}}{s^2} - \frac{\omega^2}{s^2} I$$

$$\left(1 + \frac{\omega^2}{s^2}\right) I = \frac{-s \sin \omega t e^{-st}}{s} - \frac{\omega \cos \omega t e^{-st}}{s^2}$$

$$I = \frac{-s^2}{s^2 + \omega^2} \left(\frac{s \sin \omega t e^{-st}}{s} + \frac{\omega \cos \omega t e^{-st}}{s^2} \right)$$

$$[I]_0^{\infty} = \frac{-s^2}{s^2 + \omega^2} \left[\frac{-\omega}{s} e^{-s^2/\omega} - \frac{\omega}{s^2} \right]$$

$$= \frac{s^2 \omega}{s^2 (s^2 + \omega^2)} [e^{-s^2/\omega} + \omega]$$

$$= \frac{\omega}{(s^2 + \omega^2)} [\omega + e^{-s^2/\omega}]$$

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[6]

Transform of some special functions
 $J_0(x)$ and $J_1(x)$ where

$$J_0(x) = 1 - \frac{x^2}{2^2} + \frac{x^4}{2^2 \cdot 4^2} - \frac{x^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots \quad (1)$$

and $J_1(x) = -J_1(x)$ and $L[F(t)] = x f(x) - f(0)$

Bessel function

To prove $L[J_0(x)] = \frac{1}{\sqrt{s^2 - 1}}$

$$L[J_1(x)] = 1 - \frac{3}{\sqrt{s^2 - 1}}$$

Solⁿ - $L[J_0(x)] = L[1] - L\left[\frac{x^2}{2^2}\right] + L\left[\frac{x^4}{2^2 \cdot 4^2}\right]$

$$- L\left[\frac{x^6}{2^2 \cdot 4^2 \cdot 6^2}\right] + \dots$$

$$= \frac{1}{s} - \frac{1}{2^2} \frac{L^2}{s^3} + \frac{L^4}{2^2 \cdot 4^2 \cdot 5^2} - \frac{L^6}{2^2 \cdot 4^2 \cdot 6^2} \frac{1}{s^7} + \dots$$

$$= \frac{1}{s} - \frac{1}{2 \cdot 2^3} + \frac{1}{2} \frac{3}{4 \cdot 2^5} - \frac{1}{2} \frac{3}{4} \frac{5}{6} \frac{1}{s^7} + \dots$$

$$= \frac{1}{s} \left[1 - \frac{1}{2s^2} + \frac{1}{2} \frac{3}{4s^4} - \frac{1}{2} \frac{3}{4} \frac{5}{6} \frac{1}{s^6} + \dots \right]$$

$$\frac{1}{s} [1 + \beta^2]^{-1/2} = \frac{1}{s} \left[\frac{1 + \beta^2}{1 + \beta^2} \right]^{1/2} = \frac{1}{s} \left[\frac{1 + \beta^2}{1 + \beta^2} \right]^{1/2}$$

$$[J_1(x)] = 0 \left[\frac{2x}{2^2} - \frac{4x^3}{2^2 \cdot 4^2} + \frac{6x^5}{2^2 \cdot 4^2 \cdot 6^2} + \dots \right]$$

$$= \frac{x}{2} - \frac{x^3}{2^2 \cdot 4} + \frac{x^5}{2^2 \cdot 4^2 \cdot 6} + \dots$$

$$L[J_1(x)] = L\left[\frac{x}{2}\right] - L\left[\frac{x^3}{2^2 \cdot 4}\right] + L\left[\frac{x^5}{2^2 \cdot 4^2}\right] - \dots$$

$$= \frac{1}{2} L[1] - \frac{L^3}{2^2 \cdot 4} + \frac{L^5}{2^2 \cdot 4^2 \cdot 6} - \dots$$

$$= \frac{1}{2} - \frac{3}{2 \cdot 4} + \frac{7}{2} \frac{3}{4} \frac{5}{6} - \dots$$

$$= \frac{1}{8} \left[1 - \frac{1 \cdot 3}{2 \cdot 4} + \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6} + \dots \right]$$

$$= \frac{1}{8} \left[1 + \frac{1}{\beta^2} \right]^{-1/2} = \frac{1}{8} \frac{1}{\sqrt{1 + \beta^2}} = \frac{1}{8 \sqrt{1 + \beta^2}}$$

[7] Change of scale property:-

If $L[f(t)] = f(s)$ then -

$$L[f(at)] = \frac{1}{a} f\left(\frac{s}{a}\right)$$

Solⁿ - $L[f(at)] = \int_0^{\infty} e^{-st} f(at) dt$

Let $at = u \Rightarrow t = u/a \Rightarrow dt = du/a$

$$\Rightarrow \int_0^{\infty} e^{-su/a} \frac{f(u)}{a} du$$

$$= \int_0^{\infty} e^{-ru} \frac{f(u)}{a} du \quad (\text{where } r = s/a)$$

$$= \frac{1}{a} f(r) = \frac{1}{a} f(s/a) \quad (\because r = s/a)$$

[8] Prove that $L\left[\frac{\cos \sqrt{t}}{\sqrt{t}}\right] = \sqrt{\frac{\pi}{s}} e^{-1/4s}$

$$\therefore \frac{\cos \sqrt{t}}{\sqrt{t}} = \frac{1}{\sqrt{t}} \left[1 - \frac{(\sqrt{t})^2}{2!} + \frac{(\sqrt{t})^4}{4!} - \frac{(\sqrt{t})^6}{6!} + \dots \right]$$

H.W.
$$= \left[(t)^{-1/2} - \frac{t^{1/2}}{2!} + \frac{t^{3/2}}{4!} - \frac{t^{5/2}}{6!} + \dots \right]$$

$$\therefore L \left[\frac{\cos \sqrt{t}}{\sqrt{t}} \right] = \frac{\sqrt{\pi}}{\delta^{1/2}} - \frac{1}{2!} \frac{\sqrt{\pi}}{\delta^{3/2}} + \frac{1}{4!} \frac{\sqrt{\pi}}{\delta^{5/2}} - \dots$$

We know that $\sqrt{1/2} = \sqrt{\pi}$
 $\left(\sqrt{1/2} = \frac{1}{2} \sqrt{\pi} = \frac{1}{2} \sqrt{\pi}, \sqrt{9/2} = \frac{3}{2} \cdot \frac{1}{2} \sqrt{\pi} \right)$

$$= \frac{\sqrt{\pi}}{\delta^{1/2}} - \frac{\sqrt{\pi}/2}{\delta^{3/2}} + \frac{3/2 \cdot 1/2 \sqrt{\pi}}{\delta^{5/2}} - \dots$$

$$\frac{5/2 \cdot 3/2 \cdot 1/2 \sqrt{\pi}}{\delta^{7/2}} - \dots$$

$$= \frac{\sqrt{\pi}}{\delta^{1/2}} \left[1 - \frac{1}{2 \cdot 2! \cdot \delta} + \frac{3 \cdot 1}{2 \cdot 2 \cdot 4! \cdot \delta^2} - \frac{5 \cdot 3 \cdot 1}{\delta^{3 \cdot 2 \cdot 2 \cdot 2 \cdot 4!}} \dots \right]$$

$$= \frac{\sqrt{\pi}}{\delta} \left[1 - \frac{1}{(2\delta)} \cdot \frac{1}{2!} + \frac{3 \cdot 1}{(2\delta)^2 \cdot 4!} - \frac{5 \cdot 3 \cdot 1}{(2\delta)^3 \cdot 6!} + \dots \right]$$

$$= \frac{\sqrt{\pi}}{\delta} \cdot e^{-x/\delta}$$

(9) Find $L \left[\frac{\sin at}{t} \right]$ given that

$L \left[\frac{\sin t}{t} \right] = \tan^{-1} 1/\delta$

Sol^m Since, $L \left[\frac{\sin t}{t} \right] = \tan^{-1} 1/\delta$

\therefore By the change of scale property that if $L[f(t)] = f(s)$ then $L[f(at)] = 1/a f(s/a)$.

$\therefore L \left[\frac{\sin at}{at} \right] = L \left[\frac{\sin at}{t} \right]$

$\therefore L \left[\frac{\sin at}{t} \right] = \dots$

$a L \left[\frac{\sin at}{at} \right] = a \left[\frac{1}{a} f \left(\frac{s}{a} \right) \right] = a \left[\frac{1}{a} \tan^{-1} \left(\frac{1}{s/a} \right) \right]$

$= \tan^{-1} \left(\frac{a}{s} \right)$

Sol^m (b)

if $L[f(t)] = \frac{s^2 - s + 1}{(2s+1)^2 (s-1)}$ ST $L[f(2t)] = \frac{s^2 - 2s + 4}{4(s+1)^2 (s-2)}$

By change of scale property, we have

$L[f(2t)] = \frac{1}{2} f(s/2) = \frac{1}{2} \left[\frac{(s/2)^2 - s/2 + 1}{(2 \cdot s/2 + 1)^2 (s/2 - 1)} \right]$

$= \frac{1}{2} \left[\frac{s^2 - s + 4}{(s+1)^2 (s-2)} \right]$

$= \frac{s^2 - 2s + 4}{4(s+1)^2 (s-2)}$

$(-x - 1)(d) / (p - 20)$

Transform of Error Function:—

(10)

$erf(\sqrt{x}) = \frac{2}{\sqrt{\pi}} \int_0^{\sqrt{x}} e^{-t^2} dt$ — (ii)

$L[erf(\sqrt{x})] = \frac{1}{\delta \sqrt{\delta^2 + 1}}$

In eqⁿ (i), [(6)] - Applying Laplace transform them we have—

$L[J_0(x)] = L[1] + \frac{1}{2} L[x^2] + \frac{1}{2! \cdot 4^2} L[x^4] - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} L[x^6] + \dots$

$= \frac{1}{\delta} - \frac{1}{2^2} \frac{\sqrt{3}}{\delta^2} + \frac{1}{2^2 \cdot 4^2} \frac{\sqrt{5}}{\delta^4} - \frac{\sqrt{7}}{2^2 \cdot 4^2 \cdot 6^2 \delta^6} + \dots$

$= \frac{1}{\delta} \left[1 - \frac{1 \cdot 2}{2! \delta^2} + \frac{3 \cdot 1}{2! \cdot 4 \delta^4} - \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6 \delta^6} + \dots \right]$

$e^{x/2} = \frac{2}{\sqrt{\pi}} \int_0^{\sqrt{x}} e^{-t^2} dt$

$$= \frac{1}{\beta} \left[1 + \frac{1}{\beta^2} \right]^{-1/2}$$

$$= \frac{1}{\beta} \left[\frac{\beta^2 + 1}{\beta^2} \right]^{-1/2} = \frac{1}{\beta} \frac{\sqrt{\beta^2}}{\sqrt{\beta^2 + 1}} = \frac{1}{\sqrt{\beta^2 + 1}}$$

Now,

$$J_1'(x) = -J_1(x)$$

$$\Rightarrow \frac{-2x}{2^2} + \frac{4x^3}{2^2 \cdot 4^2} - \frac{6x^5}{2^2 \cdot 4^2 \cdot 6^2} + \dots = -J_1(x)$$

$$\Rightarrow J_1(x) = \frac{x}{2} - \frac{x^3}{2^2 \cdot 4} + \frac{x^5}{2^2 \cdot 4 \cdot 6} - \dots$$

$$L[J_1(x)] = L\left[\frac{x}{2}\right] - \frac{1}{2^2 \cdot 4} L[x^3] + \frac{1}{2^2 \cdot 4 \cdot 6} L[x^5] - \dots$$

$\Gamma_2 = 1!$
 $\Gamma_4 = 3!$
 $\Gamma_6 = 5!$

$$= \frac{1}{2} \frac{\Gamma_2}{\beta^2} - \frac{1}{2^2 \cdot 4} \frac{\Gamma_4}{\beta^4} + \frac{1}{2^2 \cdot 4 \cdot 6} \frac{\Gamma_6}{\beta^6} - \dots$$

$$= \frac{1}{2\beta^2} \left[1 - \frac{3}{2 \cdot 4 \cdot \beta^2} + \frac{15}{2 \cdot 4 \cdot 6 \cdot \beta^4} - \dots \right]$$

$$= \frac{1}{2\beta^2} \left[1 - \frac{3 \cdot 1}{4 \cdot \beta^2} + \frac{5 \cdot 3 \cdot 1}{6 \cdot \beta^4} - \dots \right]$$

In eqn (ii) $erf(\sqrt{x}) = \frac{2}{\sqrt{\pi}} \int_0^{\sqrt{x}} \left[1 - \frac{t^2}{1} + \frac{t^4}{2!} - \frac{t^6}{3!} + \dots \right] dt$

$$erf(\sqrt{x}) = \frac{2}{\sqrt{\pi}} \left[t - \frac{t^3}{1 \cdot 3} + \frac{t^5}{5 \cdot 2!} - \frac{t^7}{7 \cdot 3!} + \dots \right]_0^{\sqrt{x}}$$

$$= \frac{2}{\sqrt{\pi}} \left[\sqrt{x} - \frac{x^{3/2}}{3} + \frac{x^{5/2}}{5 \cdot 2!} - \frac{x^{7/2}}{7 \cdot 3!} + \dots \right]$$

Now, using Laplace transform

$$L[erf(\sqrt{x})] = \frac{2}{\sqrt{\pi}} L\left[\sqrt{x} - \frac{x^{3/2}}{3} + \dots\right]$$

$$L[erf(\sqrt{x})] = \frac{2}{\sqrt{\pi}} \left[\frac{\Gamma_{3/2}}{\beta^{3/2}} - \frac{\Gamma_{5/2}}{3 \cdot \beta^{5/2}} + \frac{\Gamma_{7/2}}{5 \cdot 2! \cdot \beta^{7/2}} - \frac{\Gamma_{9/2}}{7 \cdot 3! \cdot \beta^{9/2}} + \dots \right]$$

$$= \frac{2}{\sqrt{\pi} \cdot \beta^{3/2}} \left[\frac{\frac{1}{2} \sqrt{\pi}}{1} - \frac{3/2 \cdot \frac{1}{2} \sqrt{\pi}}{3 \cdot \beta} + \frac{5/2 \cdot 3/2 \cdot \frac{1}{2} \sqrt{\pi}}{5 \cdot 2! \cdot \beta^2} - \dots \right]$$

$$+ \frac{7/2 \cdot 5/2 \cdot 3/2 \cdot \frac{1}{2} \sqrt{\pi}}{7 \cdot 3! \cdot \beta^{3/2}} + \dots$$

$$erf(t) = \frac{2}{\sqrt{\pi}} \int_0^t e^{-x^2} dx$$

self: best pt

$$\therefore erf(t) = \frac{2}{\sqrt{\pi}} \int_0^t \left(1 - x^2 + \frac{x^4}{2} - \frac{x^6}{6} + \frac{x^8}{24} - \dots \right) dx$$

$$= \frac{2}{\sqrt{\pi}} \left[x - \frac{x^3}{3} + \frac{x^5}{5 \cdot 2} - \frac{x^7}{7 \cdot 6} + \frac{x^9}{9 \cdot 24} - \dots \right]_0^t$$

$$= \frac{2}{\sqrt{\pi}} \left[t - \frac{t^3}{3} + \frac{t^5}{5 \cdot 2!} - \frac{t^7}{7 \cdot 3!} + \dots \right]$$

$$\therefore L[erf(\sqrt{t})] = \frac{2}{\sqrt{\pi}} \left[\frac{\frac{1}{2} \sqrt{\pi}}{\beta^{3/2}} - \frac{3/2 \cdot \frac{1}{2} \sqrt{\pi}}{3 \beta^{5/2}} + \frac{5/2 \cdot 3/2 \cdot \frac{1}{2} \sqrt{\pi}}{5 \cdot 2! \cdot \beta^{7/2}} - \dots \right]$$

$$= \frac{1}{\beta^{3/2}} \left[1 - \frac{3}{2\beta} + \frac{3!}{2 \cdot \beta^2 \cdot 2!} - \dots \right]$$

$$= \frac{1}{\beta^{3/2}} \cdot \left(1 + \frac{1}{\beta} \right)^{-1/2} = \frac{1}{\beta^{3/2}} \frac{\sqrt{\beta}}{\sqrt{\beta+1}}$$

$$= \frac{1}{\beta \sqrt{\beta+1}}$$

self: Second shifting / translation theorem;

Also known as Heaviside shifting

$$L[f(t)] = f(s) \quad \text{and}$$

$$g(t) = \begin{cases} f(t-a) & t > a \\ 0 & t < a \end{cases}$$

then $L[g(t)] = e^{-as} f(s)$

Proof: - We have

$$\begin{aligned}
 h[g(t)] &= \int_a^\infty e^{-st} g(t) dt = \int_a^0 e^{-st} g(t) dt + \int_0^\infty e^{-st} g(t) dt \\
 &= \int_0^a e^{-st} dt + \int_0^\infty e^{-st} f(t) dt \\
 &= \int_0^\infty e^{-st} f(t+a) dt \\
 &= \int_0^\infty e^{-s(a+x)} f(x) dx \\
 &= e^{-as} \int_0^\infty e^{-sx} f(x) dx = e^{-as} f(s)
 \end{aligned}$$

$t-a=x$
 $dt=dx$
 $t=a+x$

$$\Rightarrow h[f(t-a)] = e^{-as} f(s)$$

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Piecewise Continuity: A f^n f is called piecewise continuous on an interval $a \leq t \leq b$. if the interval (a, b) can be subdivided into a finite number of intervals in each of which the f^n f is continuous and has finite right and left hand limits.

We note that the right and left hand limits of f at t are denoted by

$$\lim_{\epsilon \rightarrow 0} f(t+\epsilon) = f(t, +0) = f(t^+) \quad (\epsilon > 0)$$

$$\lim_{\epsilon \rightarrow 0} f(t-\epsilon) = f(t, -0) = f(t^-) \quad (\epsilon > 0)$$

* Laplace Transform of the Derivatives

Theorem (1) Let $F(t)$ be continuous $\forall t \geq 0$ and be of exponential order α as $t \rightarrow \infty$. If $\frac{d}{dt} F(t)$ is a class A. Then Laplace transformation of derivative exists. When $s > \eta$ and

$$L\left[\frac{d}{dt} F(t)\right] = s \cdot L[F(t)] - F(0)$$

Ques
I

$$L\left[\sqrt{t} - \frac{1}{\sqrt{t}}\right]^3 =$$

$$\text{Let } f(t) = \left(\sqrt{t} - \frac{1}{\sqrt{t}}\right)^3 = t^{3/2} - 3t^{1/2} + 3t^{-1/2} - t^{-3/2}$$

$$\therefore L[f(t)] = \frac{\Gamma^{5/2}}{s^{5/2}} - \frac{3\Gamma^{3/2}}{s^{3/2}} + \frac{3\Gamma^{1/2}}{s^{1/2}} + \frac{\Gamma^{-1/2}}{s^{-1/2}}$$

$$= \frac{3 \cdot \frac{1}{2} \sqrt{\pi}}{s^{5/2}} - \frac{3 \cdot \frac{1}{2} \sqrt{\pi}}{s^{3/2}} + \frac{3 \sqrt{\pi}}{s^{1/2}} + \frac{-2\sqrt{\pi}}{s^{-1/2}}$$

$$\therefore L[f(t)] = \frac{-3\sqrt{\pi}}{4 \cdot s^{5/2}} - \frac{3\sqrt{\pi}}{2 \cdot s^{3/2}} + \frac{3\sqrt{\pi}}{s^{1/2}} + \frac{2\sqrt{\pi}}{s^{-1/2}}$$

$$= \frac{\sqrt{\pi}}{4} \left[\frac{3}{s^{5/2}} - \frac{6}{s^{3/2}} + \frac{12}{s^{1/2}} + \frac{8}{s^{-1/2}} \right]$$

Proof of this

$$L\left[\frac{d}{dt} F(t)\right] = L[f'(t)] = \int_0^{\infty} e^{-st} f'(t) dt$$

$$= \int_0^{\infty} e^{-st} d(f(t))$$

By Integrating —

$$L[F'(t)] = \left[e^{-st} f(t) \right]_0^{\infty} - \int_0^{\infty} f(t) \left[e^{-st} (-s) \right] dt$$

$$= -f(\infty) + s \int_0^{\infty} e^{-st} f(t) dt$$

$$\therefore \boxed{L[F'(t)] = s f(s) - f(0)}$$

Hence Proved

$$L[F'(t)] = L\left[\frac{d}{dt} F(t)\right] = s L[F(t)] - F(0)$$

Proof

we divide the proof into two parts (two cases)

Case-I Let $F(t)$ be continuous $\forall t \geq 0$
Case-II Let $F'(t)$ be piecewise cont. $\forall t \geq 0$

Pf: Case-I. By definition of Laplace Transform

$$L[F(t)] = \int_0^{\infty} e^{-st} F(t) dt$$

Now, $L[F'(t)] = L\left[\frac{d}{dt} F(t)\right]$

$$= \int_0^{\infty} e^{-st} \frac{d}{dt} F(t) dt$$

$$= \left[e^{-st} F(t) \right]_0^{\infty} + s \int_0^{\infty} e^{-st} F(t) dt$$

$$= \lim_{t \rightarrow \infty} e^{-st} F(t) - F(0) + s \int_0^{\infty} e^{-st} F(t) dt$$

$\therefore F(t)$ is exponential order η , we can find the constant $m > 0$. Such that $|f(t)| \leq m e^{\eta t} \forall t \geq 0$ — (2)

$$|e^{-st} F(t)| \leq e^{-st} |F(t)|$$

$$\leq e^{-st} m e^{\eta t}$$

$$\leq m e^{-(s-\eta)t}$$

$$\therefore |e^{-st} F(t)| \leq m e^{-(s-\eta)t} \text{ — (3)}$$

Now, for $s > \eta$, as $t \rightarrow \infty$, $m e^{-(s-\eta)t} \rightarrow 0$
 \therefore From eqn (1), we see that $L[F'(t)]$ exist and given by

$$L[F'(t)] = s L[F(t)] - F(0)$$

Pf: Case-II
self

Self

Then the Result of the above thm takes the form

$$L[F'(t)] = sL[F(t)] - e^{-as}[F(a+0) - F(a-0)]$$

The quantity $F(a+0) - F(a-0)$ is known as the 'Jump at the discontinuity $x=a$ '.

III If $F(t)$ is discontinuous at $t = a_1, a_2, \dots, a_n$, we break up the integral into range from 0 to ∞ .

$$L[F'(t)] = \int_0^{a_1} e^{-st} F'(t) dt + \int_{a_1}^{a_2} e^{-st} F'(t) dt + \dots + \int_{a_n}^{\infty} e^{-st} F'(t) dt$$

Thm: II Let F & F'' be cts for all $t \geq 0$. & F'' be of exponential order η as $t \rightarrow \infty$. If F'' is of class A. then Laplace transformation of $F''(t)$ exist. when $s > \eta$ & given by

$$L[F''(t)] = s^2 L[F(t)] - sF(0) - F'(0)$$

Prf:- For, proving this theorem, we have two cases.

Case-I: let $F''(t)$ be cts $\forall t \geq 0$

Case-II: let $F''(t)$ be piece-wise cts $\forall t \geq 0$.

Self

Remark-

(I) Suppose, $F(t)$ is not continuous at $t=0$, but limit $t \rightarrow 0 \lim F(t)$, i.e

The right hand limit $F(0+0)$ exist, then we can prove that -

$$L[F'(t)] = sL[F(t)] - F(0+0)$$

(II) Suppose that $F(t)$ is not continuous at $t=a$, $(0 < a < \infty)$, but RHL $F(a+0)$ and LHL $F(a-0)$ exist,

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Find the L.T of

Problem: Solve the differential eqn

$$y'' + 3y' + 2y = 0 \quad y(0) = 1, y'(0) = 2$$

Solⁿ: Here, L.T. of this transformation is given by —

$$L(y'' + 3y' + 2y) = L(0) \quad \left(\frac{s+5}{s^2+3s+2} \right)$$

$$\Rightarrow L(y'') + 3L(y') + 2L(y) = L(0) \quad 2L(F(s))$$

$$\Rightarrow s^2 L(F(s)) - sF(0) - F'(0) + 3[sL(F(s)) - sF(0)] + 2F(s) = 0$$

$$\Rightarrow L[F(s)] [s^2 + 3s + 2] - s(1) - 2 - 3(1) = 0$$

$$\Rightarrow L[F(s)] (s^2 + 3s + 2) - (s + 5) = 0$$

$$\Rightarrow L[F(s)] = \frac{s+5}{s^2+3s+2}$$

Theorem-III

(13)

Statement: Let $F(t)$ and its derivative $F'(t), F''(t), \dots, F^{(n-1)}(t)$ be continuous $\forall t \geq 0$ and be of exponential order η as $t \rightarrow \infty$ and if $F^{(n)}(t)$ is of class A, then Laplace Transform of $F^{(n)}(t)$ is

Proof - 137

$$L[F^{(n)}(t)] = s^n L[F(t)] - s^{n-1} F(0) - s^{n-2} F'(0) - \dots - F^{(n-1)}(0) \text{ where}$$

$$F^{(n)}(t) = \frac{d^n F(t)}{dt^n}$$

Hint: Method of Pf: (1) Mathematical Induction, (2) Leibniz theorem

Problem: If $L[\sin \sqrt{t}] = \frac{\sqrt{\pi}}{2s^{3/2}} e^{-\frac{1}{4s}}$, then, P.T

(11)

$$L\left[\frac{\cos \sqrt{t}}{\sqrt{t}}\right] = \sqrt{\frac{\pi}{s}} e^{-\frac{1}{4s}}$$

Theorem-III (Laplace transform of the n th order derivative) General Case: —

* $y'' + 2y' + 2y = 0$; $y(0) = 1$; $y'(0) = 2$

H.W
 $L[\sin \sqrt{t}] = \frac{\sqrt{\pi}}{2s^{3/2}} e^{-1/4s}$

$L\left[\frac{\cos \sqrt{t}}{\sqrt{t}}\right] = ?$

$F(t) = \sin \sqrt{t}$

$F'(t) = \frac{1}{2} \frac{\cos \sqrt{t}}{\sqrt{t}}$

$L[F'(t)] = s \cdot L[F(t)] - F(0)$

$L\left[\frac{1}{2} \frac{\cos \sqrt{t}}{\sqrt{t}}\right] = s \cdot \frac{\sqrt{\pi}}{2 \cdot s^{3/2}} e^{-1/4s} - 0$

$\frac{1}{2} L\left[\frac{\cos \sqrt{t}}{\sqrt{t}}\right] = \frac{1}{2} \sqrt{\frac{\pi}{s}} e^{-1/4s}$

$L\left[\frac{\cos \sqrt{t}}{\sqrt{t}}\right] = \sqrt{\frac{\pi}{s}} e^{-1/4s}$

Proof: We shall use the principle of mathematical induction to prove for $n=1$.

$L\left[\frac{d^n}{dt^n} F(t)\right] = s^n L[F(t)] - s^{n-1} F(0) - s^{n-2} F'(0) - \dots - F^{(n-1)}(0)$
 — (1)

for $n=1$, eqⁿ (1) reduces to —

$L[F'(t)] = s L[F(t)] - F(0)$

Similarly for $n=2$ & $n=3$, then eqⁿ (1) reduces to

$L[F''(t)] = s^2 L[F(t)] - s F(0) - F'(0)$

$L[L F'''(t)] = s^3 L[F(t)] - s^2 F(0) - s F'(0) - F''(0)$

This shows that, the required result is true for $n=1, 2, 3$.

Now, we shall prove that the result is true for $n=k+1$.

Now, putting $n=k$, in (1), we get

$\Rightarrow L[F^{(k+1)}(t)] = s^{k+1} L[F(t)] - s^k F(0) - s^{k-1} F'(0) - \dots - F^{(k)}(0)$

Let $G(t) = F^{(k)}(t)$ $\Rightarrow G'(t) = F^{(k+1)}(t)$

$L[G'(t)] = s L[G(t)] - G(0)$

$L[F^{(k+1)}(t)] = s L[F^{(k)}(t)] - F^{(k)}(0)$

This shows that this result is true for $n=k+1$ & $n=k$ hence by the principle of mathematical induction the required result is proved.

Multiplication by positive integral power of t

Theorem: If $F(t)$ is a fⁿ of class A and
 $L[F(t)] = f(s)$, then
 $L[t^n F(t)] = (-1)^n \frac{d^n}{ds^n} f(s)$
 Pf: - let $tF(t) = u(t)$
 $\therefore u'(t) = F(t) + tF'(t)$

since $L[F(t)] = f(s)$
 $\therefore \frac{d}{ds} f(s) = \frac{d}{ds} L[F(t)] = \int_0^\infty \frac{\partial}{\partial s} (e^{-st} F(t)) dt$
 $= \int_0^\infty (-t \cdot e^{-st} \cdot F(t)) dt$
 $= - \int_0^\infty e^{-st} (tF(t)) dt$
 $= -L[tF(t)]$

$\therefore -f'(s) = L[tF(t)]$ (Prove)

Remark: Leibnitz Rule for
 $f(s) = \int_a^b F(s,t) dt = L[F(t)]$
 $\Rightarrow \frac{d}{ds} f(s) = \int_a^b \frac{\partial}{\partial s} F(s,t) dt$

Theorem: Let $F(t)$ be a function in a t int
 (16) $L[F(t)] = f(s)$; then $L[t^n F(t)] = (-1)^n \frac{d^n}{ds^n} f(s)$ for $n=1,2,3, \dots$

Proof: We prove this theorem by Mathematical Induction
 for $n=1$,
 $L[tF(t)] = -\frac{d}{ds} f(s)$

Similarly for $n=2, n=3$, eqn (1) reduces to -
 $L[t^2 F(t)] = \frac{d^2}{ds^2} f(s)$ and $L[t^3 F(t)] = -\frac{d^3}{ds^3} f(s)$

Suppose, this result is true for $n=k$.

Note: Differentiating both side w.r.t s and apply Leibnitz

$$\frac{d}{ds} \left[\int_0^\infty e^{-st} t^k F(t) dt \right] = (-1)^k \frac{d^{k+1}}{ds^{k+1}} f(s)$$

$$\int_0^\infty \frac{\partial}{\partial s} (e^{-st} t^k F(t)) dt = (-1)^k \frac{d^{k+1}}{ds^{k+1}} f(s)$$

$$\Rightarrow \int_0^\infty e^{-st} (-t) t^k F(t) dt = (-1)^k \frac{d^{k+1}}{ds^{k+1}} f(s)$$

$$\Rightarrow \int_0^\infty e^{-st} t^{k+1} F(t) dt = (-1)^{k+1} \frac{d^{k+1}}{ds^{k+1}} f(s)$$

$$\Rightarrow L[t^{k+1} F(t)] = (-1)^{k+1} \frac{d^{k+1}}{ds^{k+1}} f(s)$$

Hence the result is true for $n=k+1$, so it is true for all n. Hence result is true.
 — (Prove)

H.W

- (i) $L[t \sin at]$ (ii) $L[t \cos at]$
- (iii) $L[t^2 \cos at]$ ($(6(s^2 - a^2 + 1)/(s^2 + 1)^4)$)
- (f) Let $F(t) = \sin at$ then
 $L[F(t)] = \frac{a}{s^2 + a^2} = f(s)$

$$\therefore L[t \cdot \sin at] = -f'(s)$$

$$= -\left[\frac{-a \cdot 2s}{(s^2 + a^2)^2} \right]$$

$$= \frac{2as}{(s^2 + a^2)^2}$$

- (ii) $L[t \cos at]$
 let $F(t) = \cos at$ then
 $L[F(t)] = \frac{s}{s^2 + a^2} = f(s)$

$$\therefore L[t \cdot \cos at] = -f'(s) = \frac{1(s^2 + a^2) - s(2s)}{(s^2 + a^2)^2}$$

$$= \frac{1}{\beta} [\ln v - \ln(\beta v)]$$

$$= \frac{1}{\beta} [-\ln \frac{v+\beta}{v}] = \frac{1}{\beta} [-\ln(1 + \frac{\beta}{v})]$$

$$= \frac{1}{\beta} [0 - (-\ln(\beta+1))] = \frac{1}{\beta} \ln(\beta+1)$$

$$L[E(t)] = \frac{1}{\beta} \log(\beta+1)$$

II-Method

$$e^{-u} = 1 - u + \frac{u^2}{2!} - \frac{u^3}{3!} + \frac{u^4}{4!} - \frac{u^5}{5!} + \dots$$

$$= 1 - u + \frac{u^2}{2!} - \frac{u^3}{3!} + \frac{u^4}{4!} - \frac{u^5}{5!} + \dots$$

$$\int_1^{\infty} e^{-u} du = 1$$

5) Laguerre Polynomial

$$L_n(t) = \frac{e^t}{n!} \frac{d^n}{dt^n} (e^{-t} t^n), n=0,1,2,\dots$$

then Prove that

$$L[L_n(t)] = \frac{(\beta-1)^n}{\beta^{n+1}}$$

Here we have

$$L_n(t) = \frac{e^t}{n!} \frac{d^n}{dt^n} (e^{-t} t^n)$$

$$L[L_n(t)] = \int_0^{\infty} e^{-st} \left[\frac{e^t}{n!} \frac{d^n}{dt^n} (e^{-t} t^n) \right] dt$$

$$= \frac{1}{n!} \int_0^{\infty} \frac{e^{(1-s)t}}{e^{-t}} \frac{d^n}{dt^n} (e^{-t} t^n) dt$$

$$= \frac{1}{n!} \int_0^{\infty} e^{-(s-1)t} \frac{d^n}{dt^n} (e^{-t} t^n) dt$$

$$= \frac{1}{n!} [(1-\beta) \int_0^{\infty} e^{-(\beta-1)t} \frac{d^n}{dt^n} (e^{-t} t^n) dt]$$

$$\therefore e^{-t} = t e^{-t} + \int_0^t e^{-t} dt$$

$$L[L_n(t)] = 0 + \frac{s-1}{n!} \int_0^{\infty} e^{-(s-1)t} \frac{d^n}{dt^n} (e^{-t} t^n) dt \quad (\beta > 1)$$

Repeating the above process $\frac{d^n}{dt^n}$ (n-times)

$$L[L_n(t)] = \frac{(s-1)^n}{n!} \int_0^{\infty} e^{-(s-1)t} (e^{-t} t^n) dt$$

$$= \frac{(s-1)^n}{n!} \int_0^{\infty} e^{-st+t-t} t^n dt = \frac{(s-1)^n}{n!} \int_0^{\infty} e^{-st} t^n dt$$

$$= \frac{(s-1)^n}{n!} L[t^n]$$

$$= \frac{(s-1)^n}{n!} L[t^n]$$

$$L[L_n(t)] = \frac{(s-1)^n}{n!} \times \frac{n!}{\beta^{n+1}} = \frac{(s-1)^n}{\beta^{n+1}} \quad [\because L[t^n] = \frac{n!}{\beta^{n+1}}]$$

6) UNIT Step function: (or Heaviside Unit Function)

The unit step function $u(t-a)$ is defined as -

$$u(t-a) = \begin{cases} 0 & \text{when } t < a \\ 1 & \text{when } t \geq a \end{cases}$$

Example Express the following f^n in terms of unit step f^n and find its L.T

$$f(t) = \begin{cases} 8 & \text{when } t < 2 \\ 6 & \text{when } t \geq 2 \end{cases}$$

Prove that $L[J_0(t)] = \frac{1}{\sqrt{s^2+1}}$ and deduce that (i) $L[J_0(at)] = \frac{1}{\sqrt{s^2+a^2}}$

(ii) $L[t J_0(at)] = \frac{s}{(s^2+a^2)^{3/2}}$

(iii) $L[e^{at} J_0(at)] = \frac{1}{\sqrt{s^2+2as+2a^2}}$

$$\frac{s^2 + a^2 - 2s^2}{(s^2 + a^2)^2} = \frac{-s^2 + a^2}{(s^2 + a^2)^2} = \frac{s^2 - a^2}{(s^2 + a^2)^2}$$

(iii) $\therefore L[\cos at] = \frac{s}{s^2 + a^2} = f(s)$
 $\therefore L[t^3 \cos at] = (-1)^3 \frac{d^3}{ds^3} f(s)$

$$f(s) = \frac{s}{s^2 + a^2}$$

$$f'(s) = \frac{1 \cdot (s^2 + a^2) - s \cdot (2s)}{(s^2 + a^2)^2}$$

$$f''(s) = \frac{d}{ds} \left[\frac{s^2 + a^2 - 2s^2}{(s^2 + a^2)^2} \right] = \frac{d}{ds} \left[\frac{a^2 - s^2}{(s^2 + a^2)^2} \right]$$

$$= \frac{-2s(s^2 + a^2)^2 - (a^2 - s^2)(2(s^2 + a^2) \cdot 2s)}{(s^2 + a^2)^4}$$

$$= \frac{-2s(s^2 + a^2) [s^2 + a^2 + 2(a^2 - s^2)]}{(s^2 + a^2)^4}$$

$$= \frac{-2s[3a^2 - s^2]}{(s^2 + a^2)^3}$$

$$f'''(s) = \frac{[-6a^2 + 6s^2](s^2 + a^2)^3 - [3a^2 - s^2](3(s^2 + a^2)^2 \cdot 2s)}{(s^2 + a^2)^6}$$

Cancelled out $(s^2 + a^2)$ from numerator and denominator

$$= \frac{6[s^2 - a^2](s^2 + a^2) - 2s[3a^2 + 3a^2 - s^2]}{(s^2 + a^2)^4}$$

$$= \frac{6[s^4 - a^4 - 2s^2 a^2 - 6s^2 a^2]}{(s^2 + a^2)^4}$$

$$= \frac{-6[s^4 + a^4 + 8s^2 a^2]}{(s^2 + a^2)^4}$$

* Some special function:-

(1) The Sine integral:- The sine integral is defined as

$$Si(t) = \int_0^t \frac{\sin u}{u} du$$

To prove that $L[Si(t)] = \frac{1}{s} \tan^{-1} \left(\frac{1}{s} \right)$

Self
Pr:-

$$\therefore \sin u = u - \frac{u^3}{3!} + \frac{u^5}{5!} - \frac{u^7}{7!} + \dots$$

$$\therefore \frac{\sin u}{u} = 1 - \frac{u^2}{3!} + \frac{u^4}{5!} - \frac{u^6}{7!} + \dots$$

$$\therefore \int_0^t \frac{\sin u}{u} du = \left[u - \frac{u^3}{3 \cdot 3!} + \frac{u^5}{5 \cdot 5!} - \frac{u^7}{7 \cdot 7!} + \dots \right]_0^t$$

$$= \left[t - \frac{t^3}{3 \cdot 3!} + \frac{t^5}{5 \cdot 5!} - \frac{t^7}{7 \cdot 7!} + \dots \right]$$

Self

$$\therefore L[Si(t)] = L \left[t - \frac{t^3}{3 \cdot 3!} + \frac{t^5}{5 \cdot 5!} - \frac{t^7}{7 \cdot 7!} + \dots \right]$$

By linearity Prop. - $= L[t] - \frac{1}{3} L[t^3/3!] + \frac{1}{5} L[t^5/5!] - \dots$

$$= \frac{1}{s^2} - \frac{1}{3} \frac{t^3}{3! s^4} + \frac{1}{5} \frac{t^5}{5! s^6} - \frac{t^7}{7! s^8} + \dots$$

$\begin{matrix} \Gamma_1 = 3 \\ \Gamma_2 = 5 \\ \Gamma_3 = 7 \end{matrix}$

$$= \frac{1}{s} \left[\frac{1}{s} - \frac{1}{3} \left(\frac{1}{s} \right)^3 + \frac{1}{5} \left(\frac{1}{s} \right)^5 - \frac{1}{7} \left(\frac{1}{s} \right)^7 + \dots \right]$$

$$= \frac{1}{s} \left[\frac{1}{s} - \frac{1}{3} \left(\frac{1}{s} \right)^3 + \frac{1}{5} \left(\frac{1}{s} \right)^5 - \dots \right] = \frac{1}{s} \tan^{-1} \left(\frac{1}{s} \right)$$

(2) The Cosine integral is denoted by

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$$Ci(t) = \int_t^\infty \frac{\cos u}{u} du$$

To prove,

$$L[Ci(t)] = \frac{1}{2s} \log(s^2 + 1)$$

Self

$$\cos x = x - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots$$

$$L\{t^n f(t)\} = (-1)^n \frac{d^n f(s)}{ds^n}$$

$$\frac{\cos u}{u} = \frac{1}{4} \left[1 - \frac{u^2}{2!} + \frac{u^4}{4!} - \frac{u^6}{6!} + \frac{u^8}{8!} - \dots \right]$$

Self

$$= \left[\frac{1}{u} - \frac{u}{2!} + \frac{u^3}{4!} - \frac{u^5}{6!} + \frac{u^7}{8!} - \dots \right]$$

$$\therefore \int_t^\infty \frac{\cos u}{u} du = \left[\log u - \frac{u^2}{2!} + \frac{u^4}{4!} - \frac{u^6}{6!} + \dots \right]_t^\infty$$

$$\therefore L\{t \cdot F'(t)\} = - \int_t^\infty \frac{\cos u}{u} dt$$

As we know that $F(t) = \int_t^\infty \frac{\cos u}{u} du = - \int_\infty^t \frac{\cos u}{u} du$

$$F'(t) = -\frac{\cos t}{t} \Rightarrow t F'(t) = -\cos t$$

$$\Rightarrow L\{t F'(t)\} = L\{-\cos t\} = -\frac{\beta}{\beta^2 + 1}$$

$$\Rightarrow -\frac{d}{d\beta} L\{F(t)\} = -\frac{\beta}{\beta^2 + 1} \quad (\because L\{F(t)\} = \frac{d}{d\beta} f(s))$$

$$\Rightarrow \frac{d}{d\beta} L\{F(t)\} = \frac{\beta}{\beta^2 + 1} \quad (\because L\{F'(t)\} = -f(s))$$

$$\Rightarrow \frac{d}{d\beta} \left[\beta L\{F(t)\} - f(s) \right] = 0$$

$$\Rightarrow \beta L\{F(t)\} = f(s) + c \quad (\because f(0) \text{ is constant})$$

$$\Rightarrow \beta f(s) = \frac{1}{2} \log(\beta^2 + 1) + c \quad (\text{On integrating})$$

$$\lim_{\beta \rightarrow 0} \beta f(s) = \lim_{t \rightarrow \infty} f(t) = 0 \quad (\text{As } f(0) = 0, \beta \rightarrow 0 \Rightarrow 0 \cdot f(0) = 0)$$

$$\therefore f(s) = \frac{1}{2\beta} \log(\beta^2 + 1)$$

$$\therefore f(s) = L\{F(t)\} = L\left\{ \int_t^\infty \frac{\cos u}{u} du \right\}$$

$$\Rightarrow L\{F(t)\} = \frac{1}{2\beta} \log(\beta^2 + 1)$$

Initial value theorem:

Let $F(t)$ be continuous for all $t \geq 0$ and exponential order as $t \rightarrow \infty$ and $F(t)$ is of class A, then

$$I. \lim_{t \rightarrow 0} F(t) = \lim_{\beta \rightarrow \infty} \beta L\{F(t)\} = f(0)$$

$$\lim_{t \rightarrow \infty} F(t) = \lim_{\beta \rightarrow 0} \beta L\{F(t)\}$$

Final value theorem:

$$\lim_{t \rightarrow \infty} F(t) = \lim_{\beta \rightarrow 0} \beta L\{F(t)\} = \lim_{s \rightarrow 0} s f(s)$$

The Exponential Integral:

The exponential order is defined as

$$E(t) = \int_t^\infty \frac{e^{-u}}{u} du$$

then prove that

$$L\{E(t)\} = \frac{1}{\beta} \log(1 + \beta)$$

Proof: $L\{E(t)\} = L\left[\int_t^\infty \frac{e^{-u}}{u} du \right]$ let $u = tv \Rightarrow du = tv dv$

(Self) $\frac{1}{u} = \frac{1}{tv} = \frac{1}{t} \cdot \frac{1}{v}$

$$= \int_0^\infty e^{-\beta t} \left[\int_1^\infty \frac{e^{-tv}}{v} dv \right] dt$$

$$= \int_0^\infty \frac{1}{v} \left[\int_0^\infty e^{-t(\beta+v)} dt \right] dv \quad (\text{On changing the order of integration})$$

$$= \int_0^\infty \frac{1}{v} \left[\int_0^\infty e^{-t(\beta+v)} dt \right] dv$$

$$= \int_0^\infty \frac{1}{v} \left[-\frac{e^{-t(\beta+v)}}{\beta+v} \right]_0^\infty dv = \int_0^\infty \frac{1}{v} \left[0 - \left(-\frac{1}{\beta+v}\right) \right] dv$$

$$= \int_0^\infty \frac{1}{v} \left[\frac{1}{v} - \frac{1}{\beta+v} \right] dv = \frac{1}{\beta} \int_0^\infty \left[\frac{1}{v} - \frac{1}{\beta+v} \right] dv$$

(iv) $\int_0^{\infty} J_0(t) dt = 1$

(v) $\int_0^{\infty} t e^{-st} J_0(4t) dt = ?$

H.W (6) Unit step f^n

We have

$$u(t-a) = \begin{cases} 0 & t < a \\ 1 & t \geq a \end{cases} \quad \text{--- (1)}$$

$$\therefore L[u(t-a)] = \int_0^{\infty} e^{-st} u(t-a) dt$$

$$= \int_0^a e^{-st} \cdot 0 dt + \int_a^{\infty} e^{-st} \cdot 1 dt$$

$$= 0 + \int_a^{\infty} e^{-st} dt =$$

$$= \lim_{\beta \rightarrow \infty} \int_a^{\beta} e^{-st} dt = \left[\frac{-e^{-st}}{s} \right]_a^{\beta}$$

$$= \lim_{\beta \rightarrow \infty} \left[\frac{-e^{-s\beta} + e^{-as}}{s} \right]$$

$$= \lim_{\beta \rightarrow \infty} \left[\frac{-e^{-s\beta} + e^{-as}}{s} \right]$$

$$L[u(t-a)] = \left[\frac{0 + e^{-as}}{s} \right] = \frac{e^{-as}}{s}$$

(7) (i) T.S $L[J_0(at)] = \frac{1}{\sqrt{s^2+1}}$

$$\therefore L[J_0(t)] = L\left[1 + \frac{t^2}{2^2} + \frac{t^4}{2^2 \cdot 4^2} - \frac{t^6}{2^2 \cdot 4^2 \cdot 6^2} + \dots\right]$$

$$= L[1] - \frac{1}{2^2} L[t^2] + \frac{1}{2^2 \cdot 4^2} L[t^4] - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} L[t^6] + \dots$$

$$= \frac{1}{s} - \frac{1}{2^2} \frac{2!}{s^3} + \frac{1}{2^2 \cdot 4^2} \frac{4!}{s^5} - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} \frac{6!}{s^7} + \dots$$

$$= \frac{1}{s} \left[1 - \frac{1!}{2 \cdot s^2} + \frac{3 \cdot 1}{2 \cdot 4 \cdot s^4} - \frac{5 \cdot 3 \cdot 1}{2 \cdot 4 \cdot 6 \cdot s^6} + \dots \right]$$

$$= \frac{1}{s} \left[\left(1 + \frac{1}{s^2}\right)^{-1/2} \right] = \frac{1}{s} \left[\left(\frac{s^2+1}{s^2}\right)^{-1/2} \right]$$

$$= \frac{1}{s} \int \frac{s^2}{s^2+1} = \frac{1}{s} \frac{s}{\sqrt{s^2+1}}$$

$$= \frac{1}{\sqrt{s^2+1}}$$

$$\therefore L[J_0(t)] = \frac{1}{\sqrt{s^2+1}}$$

(ii) T.S $L[J_0(at)] = \frac{1}{\sqrt{s^2+a^2}}$

(i)

$$L[J_0(at)] = L\left[1 - \frac{(at)^2}{2^2} + \frac{(at)^4}{2^2 \cdot 4^2} - \frac{(at)^6}{2^2 \cdot 4^2 \cdot 6^2} + \frac{(at)^8}{2^2 \cdot 4^2 \cdot 6^2 \cdot 8^2} - \dots\right]$$

$$= L[1] - \frac{1}{2^2} L[(at)^2] + \frac{1}{2^2 \cdot 4^2} L[(at)^4] - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} L[(at)^6] + \dots$$

$$= \frac{1}{s} - \frac{1}{2^2} \frac{1}{s^3} + \frac{1}{2^2 \cdot 4^2} \frac{1}{s^5} - \frac{1}{2^2 \cdot 4^2 \cdot 6^2} \frac{1}{s^7} + \dots$$

Since $L[f(at)] = \frac{1}{a} f\left(\frac{s}{a}\right)$ (where $f(s) = L[F(t)]$)

$$\therefore L[J_0(at)] = \frac{1}{a} f\left(\frac{s}{a}\right) \text{ where } f(s) = \frac{1}{\sqrt{s^2+1}}$$

$$\therefore f\left(\frac{s}{a}\right) = \frac{1}{\sqrt{\left(\frac{s}{a}\right)^2+1}} = \frac{1}{\sqrt{\frac{s^2+a^2}{a^2}}} = \frac{a}{\sqrt{s^2+a^2}}$$

$$\therefore L[J_0(at)] = \frac{1}{a} f\left(\frac{s}{a}\right) = \frac{1}{\sqrt{s^2+a^2}}$$

(ii) (i)

As we know that

$$L[t F(t)] = -\frac{d}{ds} f(s)$$

where $L[F(t)] = f(s) = \frac{1}{\sqrt{s^2+1}}$

Similarly

$$L[t J_0(at)] = -\frac{d}{ds} f\left(\frac{\beta}{a}\right) = -\frac{d}{ds} \left[\frac{1}{\sqrt{s^2+a^2}} \right] = \frac{1}{\sqrt{s^2+a^2}}$$

$$= -\frac{d}{ds} \left(\frac{1}{\sqrt{s^2+a^2}} \right) = \frac{s}{(s^2+a^2)^{3/2}}$$

$-f(\beta) = f(\beta, a)$

$$L[t J_0(4t)] = \frac{0 - \frac{1}{2} \cdot 2s}{(s^2+4^2)^{3/2}} = \frac{s}{(s^2+16)^{3/2}}$$

(7)(2)(iii) As $L[F(t)] = f(s)$
we have $L[e^{-at} F(t)] = f(s+a)$

Here we have

$$L[F(t)] = f(s) = \frac{1}{\sqrt{s^2+a^2}}$$

$$\Rightarrow L[F(at)] = f(s+a) = \frac{1}{\sqrt{(s+a)^2+a^2}} = L[J_0(at)]$$

$$\therefore L[e^{-at} F(at)] = f(s+a)$$

$$\therefore L[e^{-at} J_0(at)] = \frac{1}{\sqrt{s^2+2a^2+a^2}}$$

(iv) T.P. $\int_0^\infty J_0(t) dt = 1$

$$\text{Aq } L[J_0(t)] = \int_0^\infty e^{-st} J_0(t) dt = \frac{1}{\sqrt{s^2+1}}$$

$$\Rightarrow \lim_{s \rightarrow 0} \int_0^\infty e^{-st} J_0(t) dt = \frac{1}{\sqrt{s^2+1}} \lim_{s \rightarrow 0} 1$$

$$\Rightarrow \int_0^\infty J_0(t) dt = 1 = \lim_{s \rightarrow 0} \frac{1}{\sqrt{s^2+1}}$$

$$(v) \int_0^\infty t \cdot e^{-st} J_0(4t) dt = \int_0^\infty e^{-st} (t \cdot J_0(4t)) dt$$

Here $\beta=3, a=4$

As we already proved that

$$L[e^{-at} J_0(at)] = \frac{1}{\sqrt{s^2+2a^2+a^2}}$$

Now putting $a=4, \beta=3$

$$L[e^{-4t} J_0(4t)] = \frac{1}{\sqrt{9+24+32}} = \frac{1}{\sqrt{65}}$$

$$= \frac{1}{\sqrt{65}} \text{ Ans}$$

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$$L[t J_0(at)] = \frac{s}{(s^2+a^2)^{3/2}}$$

$$= \frac{3}{(9+16)^{3/2}} = \frac{3}{5\sqrt{5}}$$

New Chapter

INVERSE Laplace Transform:

If $f(s)$ be the Laplace transform of a f'n $F(t)$ i.e

$$f(s) = L[F(t)]$$

then

$F(t)$ is called the Inverse Laplace transform of the f'n $f(s)$ and is written as

$$F(t) = L^{-1}\{f(s)\}$$

where L^{-1} is the inverse Laplace transform operator.

Formula for Inversion

$$f(s) \quad L^{-1}\{f(s)\} = F(t)$$

$f(s)$

$F(t) = L^{-1}$

- (i) $\frac{1}{s}$
- (ii) $\frac{1}{s^{n+1}}, n \geq 0$ $\frac{t^n}{n!}$
- (iii) $\frac{1}{s^{n+1}}, n \geq 0$ $\frac{t^n}{n!}$
- (iv) $\frac{1}{s-a}, s > a$ e^{at}
- (v) $\frac{1}{s+a}, s > a$ e^{-at}
- (vi) $\frac{1}{s^2+a^2}, s > 0$ $\frac{\sin at}{a}$
- (vii) $\frac{s}{s^2+a^2}, s > 0$ $\cos at$
- (viii) $\frac{1}{s^2-a^2}, s > |a|$ $\sinh at$
- (ix) $\frac{s}{s^2-a^2}, s > |a|$ $\cosh at$

Theorem: If $f_1(s)$ and $f_2(s)$ respectively be the L.T of the $f^n F_1(t)$ and $F_2(t)$ respectively. C_1 & C_2 be two constant, then

$$L^{-1}\{C_1 f_1(s) + C_2 f_2(s)\} = C_1 L^{-1}\{f_1(s)\} + C_2 L^{-1}\{f_2(s)\}$$

Proof: Here, we have given that $L[F_1(t)] = f_1(s)$ & $L[F_2(t)] = f_2(s)$
 By L.T.S

$$L^{-1}\{C_1 L[F_1(t)] + C_2 L[F_2(t)]\}$$

$$= L^{-1}\{L[C_1 F_1(t)] + L[C_2 F_2(t)]\}$$

$$= L^{-1}\{C_1 L[F_1(t)] + C_2 L[F_2(t)]\}$$

We know that the linearity property of L.T

$$L[C_1 F_1(t) + C_2 F_2(t)] = C_1 L[F_1(t)] + C_2 L[F_2(t)]$$

Given that

$$f_1(s) = L[F_1(t)] \quad \& \quad f_2(s) = L[F_2(t)]$$

From (1) & (2), we have

$$L[C_1 f_1(s) + C_2 f_2(s)] = C_1 f_1(s) + C_2 f_2(s)$$

$$\text{i.e. } L[C_1 F_1(t) + C_2 F_2(t)] = C_1 f_1(s) + C_2 f_2(s)$$

By defⁿ of L.T

$$\{F(t) = L^{-1}\{f(s)\}$$

$$C_1 F_1(t) + C_2 F_2(t) = L^{-1}\{C_1 f_1(s) + C_2 f_2(s)\}$$

$$\Rightarrow C_1 L^{-1}\{f_1(s)\} + C_2 L^{-1}\{f_2(s)\} = L^{-1}\{C_1 f_1(s) + C_2 f_2(s)\}$$

H.W

Example: Evaluate (i) $L^{-1}\left[\frac{3s+7}{s^2-2s+3}\right] = 4e^{2t} - e^{-t}$

(ii) $L^{-1}\left[\frac{s^2}{(s+1)(s+2)(s+3)}\right] = \frac{1}{2}e^{-t} - 4e^{-2t} + \frac{9}{2}e^{-3t}$

Solⁿ (i) $\frac{3s+7}{s^2-2s+3} = \frac{3s+7}{(s+1)(s-3)} = \frac{A}{s+1} + \frac{B}{s-3}$

(iii) $L^{-1}\left[\frac{s}{(s+a)(s^2+b^2)}\right] \rightarrow \frac{C \cos bt - C \sin bt}{a^2 + b^2}$

(iv) $L^{-1}\left[\frac{1}{(s+1)(s^2+1)}\right] \rightarrow \frac{1}{2}[e^{-t} - \cos t + \sin t]$

$$\frac{A s + B}{(s^2+a^2)} + \frac{C s + D}{(s^2+b^2)} = \frac{(A s + B)(s^2+b^2) + (C s + D)(s^2+a^2)}{(s^2+a^2)(s^2+b^2)}$$

* I. First Shifting theorem (For ILT)

If $L^{-1}\{f(s)\} = F(t)$ then

$$L^{-1}\{f(s-a)\} = e^{at} L^{-1}\{f(s)\}$$

Proof: Using $s \rightarrow s-a$

~~I~~ II. Second shifting theorem (I.L.T)

If $L^{-1}\{f(s)\} = F(t)$ then

$$L^{-1}\{e^{-as} f(s)\} = g(t)$$

$$\text{where } g(t) = \begin{cases} F(t-a) & t > a \\ 0 & t < a \end{cases}$$

Proof: By defⁿ, $L[g(t)] = \int_0^{\infty} e^{-st} g(t) dt$

$$= \int_0^a e^{-st} \cdot 0 dt + \int_a^{\infty} e^{-st} F(t-a) dt$$

$$= \int_a^{\infty} e^{-st} dt + \int_a^{\infty} e^{-st} F(t-a) dt$$

$$= \int_a^{\infty} e^{-st} F(t-a) dt$$

* ~~Second Shifting~~ $\Rightarrow dt = dx$

$t=a, x=0$
$t=\infty, x=\infty$
$t=a+x$

$$L[g(t)] = \int_0^{\infty} e^{-s(a+x)} F(x) dx$$

$$= \int_0^{\infty} e^{-sa} e^{-sx} F(x) dx$$

$$= e^{-sa} \int_0^{\infty} e^{-sx} F(x) dx$$

(By defⁿ of L)

$$= e^{-as} \cdot f(s)$$

$$\therefore g(t) = L^{-1}\{e^{-as} f(s)\}$$

Remark: Unit step fⁿ for I.L.T —

$$H(t-a) = \begin{cases} 1 & \text{when } t > a \\ 0 & t < a \end{cases}$$

Change of scale property: —

If $L^{-1}\{f(s)\} = F(t)$ then

$$L^{-1}\{f(as)\} = \frac{1}{a} F\left(\frac{t}{a}\right)$$

Proof: \therefore By defⁿ of L.T, we know that

$$L\left[\frac{1}{a} F\left(\frac{t}{a}\right)\right] = \int_0^{\infty} e^{-st} \cdot \frac{1}{a} F\left(\frac{t}{a}\right) dt$$

$$\text{Let } t/a = x \Rightarrow dt = a dx$$

$$= \frac{1}{a} \int_0^{\infty} e^{-sax} \cdot a F(x) dx$$

$$= \int_0^{\infty} e^{-sax} F(x) dx$$

$$= \int_0^{\infty} e^{-s(at)} F(t) dt$$

$$L\left[\frac{1}{a} F\left(\frac{t}{a}\right)\right] = f(as)$$

$$\Rightarrow L^{-1}\{f(as)\} = \frac{1}{a} F\left(\frac{t}{a}\right)$$

\rightarrow Next - Congruent $F(t) \cdot g(t) = \int_0^t$

Defⁿ:

Convolution:

Let $F(t)$ and $G(t)$ be the two fⁿ in t , then define the convolution

$$F(t) * G(t) = \int_0^t F(u) G(t-u) du$$

or $\int_0^t G(t) F(t-u) du$

This relation $F * G$ is also called as Resultant of $F(t)$ and $G(t)$ or as Faltung.

Properties of Convolution:-

(i) Commutative law of Convolution

$$F(t) * G(t) = G(t) * F(t)$$

Pf: $\because F(t) * G(t) = \int_0^t F(u) G(t-u) du$

Let $t-u = x$ & $u = t-x$
 $\Rightarrow dt = -dx$

$$\therefore F(t) * G(t) = \int_0^t F(t-x) G(x) dx$$

$$= \int_0^t G(x) F(t-x) dx = G(t) * F(t)$$

(ii) Let $F(t)$, $G(t)$ and $H(t)$ be the fⁿ in t , then

$$F(t) * [G(t) + H(t)] = F(t) * G(t) + F(t) * H(t)$$

(iii) $F(t) * [G(t) * H(t)] = [F(t) * G(t)] * H(t)$

(Verify)

$$\therefore G(t) * H(t)$$

$$(ii) J(t) = G(t) * H(t) = \int_0^t G(u) \cdot H(t-u) du$$

or $J(t) = H(t) * G(t) = \int_0^t G(t-u) \cdot H(u) du$ ✓

$$\therefore F(t) * [G(t) * H(t)] = F(t) * J(t)$$

$$\therefore F(t) * J(t) = \int_0^t F(w) J(t-w) dw$$

$$= \int_0^t F(w) \left\{ \int_0^{t-w} G(t-u-w) \cdot H(u) du \right\} dw$$

$$= \int_0^t F(w) \int_0^{t-w} G(t-u-w) \cdot H(u) du dw$$

$$= \int_0^t H(u) \int_0^{t-u} G(t-w-u) \cdot F(w) dw du$$

$$= \int_0^t H(u) \int_0^{t-u} F(w) \cdot G(t-u-w) dw du$$

$$= H(t) * [F(t) * G(t)]$$

$$\therefore F(t) * [G(t) * H(t)] = H(t) * [F(t) * G(t)]$$

$$= [F(t) * G(t)] * H(t)$$

Proof

$$F(t) * [G(t) + H(t)] = \int_0^t F(u) [G(t-u) + H(t-u)] du$$

$$= \int_0^t F(u) G(t-u) du + \int_0^t F(u) H(t-u) du$$

$$= F(t) * G(t) + F(t) * H(t)$$

$$\therefore F(t) * [G(t) * H(t)] = F(t) * G(t) + F(t) * H(t)$$

(Verify the convolution theorem)

The Convolution theorem:

(w) let $F(t)$ and $G(t)$ be the two fⁿ in t .
then define:

$$L^{-1}[f(s)] = F(t) \text{ and } L^{-1}[g(s)] = G(t)$$

Prove that

$$L^{-1}[f(s) \cdot g(s)] = \int_0^t F(u) \cdot G(t-u) du = F(t) * G(t)$$

Proof - let $H(t) = F(t) * G(t)$ (1)

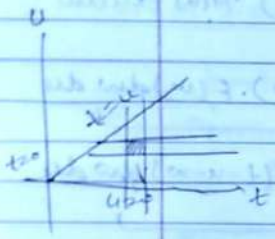
By def $L[H(t)] = \int_0^{\infty} e^{-st} H(t) dt$

$$= \int_0^{\infty} e^{-st} \{ F(t) * G(t) \} dt$$

$$= \int_0^{\infty} e^{-st} \left\{ \int_0^t F(u) G(t-u) du \right\} dt$$

$$= \int_0^{\infty} F(u) \left\{ \int_{t=u}^{\infty} e^{-st} G(t-u) dt \right\} du$$

$$= \int_0^{\infty} e^{-su} \cdot F(u) \left\{ \int_{t=u}^{\infty} e^{-s(t-u)} G(t-u) dt \right\} du$$



Example. Apply the Convolution theorem to prove that -

$$\beta(m, n) = \int_0^1 u^{m-1} (1-u)^{n-1} du = \frac{\Gamma(m) \Gamma(n)}{\Gamma(m+n)}$$

let $F(t) = t^{m-1}$ & $G(t) = t^{n-1}$

then by Convolution theorem, we have

$$L[\beta(m, n)] = L \left[\int_0^t u^{m-1} (t-u)^{n-1} du \right]$$

$$L[F(t) * G(t)] = L[F(t)] \cdot L[G(t)]$$

Remark: (1) The Convolution theorem can be written as -

$$L[F(t) * G(t)] = L[F(t)] \cdot L[G(t)]$$

(2) While using the Convolution theorem

We use one of the following two ways

$$L^{-1}\{f(s) \cdot g(s)\} = \int_0^t F(u) G(t-u) du = \int_0^t G(u) F(t-u) du$$

$$= F(t) * G(t)$$

$$f(s) \cdot g(s) = L[F(t) * G(t)]$$

$$\Rightarrow L[F(t) * G(t)] = [L[F(t)]] [L[G(t)]]$$

$$L \left[\int_0^t u^{m-1} (t-u)^{n-1} du \right] = L[t^{m-1}] L[t^{n-1}]$$

$$\int_0^t u^{m-1} (t-u)^{n-1} du = L^{-1} \left\{ \frac{\Gamma(m) \Gamma(n)}{s^m \cdot s^n} \right\}$$

$$= \Gamma(m) \Gamma(n) L^{-1} \left\{ \frac{1}{s^{m+n}} \right\}$$

$$= \Gamma(m) \Gamma(n) L^{-1} \left\{ \frac{1}{s^{m+n}} \right\}$$

$$= \frac{\Gamma(m) \Gamma(n)}{1} \cdot \frac{t^{m+n-1}}{\Gamma(m+n)}$$

$$\Rightarrow \int_0^t u^{m-1} (t-u)^{n-1} du = \frac{\Gamma(m) \Gamma(n)}{1} \cdot \frac{t^{m+n-1}}{\Gamma(m+n)}$$

putting $t=1$.

$$\int_0^1 u^{m-1} (1-u)^{n-1} du = \frac{\Gamma(m) \Gamma(n)}{\Gamma(m+n)} \cdot (1)$$

* * *

29/1/16 (1) Evaluate the value

(a) $L^{-1} \left[\frac{s^2}{(s^2+a^2)(s^2+b^2)} \right] = \frac{As + B}{(s^2+a^2)^2} + \frac{Cs + D}{s^2+b^2}$

(b) $L^{-1} \left[\frac{s^4}{(s^2+a^2)(s^2+b^2)} \right]$

using the Convolution theorem. Verify that the part is

$$p(u, n) + p(u, n+1) = p(u, n)$$

using the convolution theorem.

we know that

(6)

$$L[F(t)] \cdot L[g(t)] = L \int_0^t F(u) \cdot g(t-u) du$$

$$\text{Let } f(s) = \frac{1}{s^2 + a^2} \text{ and } g(s) = \frac{1}{s^2 + a^2}$$

$$\Rightarrow L^{-1}(f(s)) = \cos at \text{ and } L^{-1}(g(s)) = \cos at$$

$$\therefore L^{-1}[f(s)g(s)] = \cos at \cdot \cos at = \cos^2 at$$

$$\text{Let } F(t) = \cos at = g(t)$$

\therefore By convolution

$$L^{-1}[f(s)g(s)] = \int_0^t \cos au \cdot \cos a(t-u) du$$

$$= \int_0^t \cos au [\cos at \cdot \cos au + \sin at \cdot \sin au] du$$

$$= \int_0^t [\cos at \cdot \cos^2 au + \frac{\sin at \cdot \sin 2au}{2}] du$$

$$= \cos at \int_0^t \frac{\cos 2a\theta + 1}{2} du + \frac{\sin at}{2} \int_0^t \sin 2au du$$

$$= \frac{\cos at}{2} \left[\frac{\sin 2au}{2a} + u \right]_0^t + \frac{\sin at}{2} \left[-\frac{\cos 2au}{2a} \right]_0^t$$

$$= \frac{\cos at}{2} \left[\frac{\sin 2at}{2a} + t \right] + \frac{\sin at}{2} \left[-\frac{\cos 2at}{2a} + \frac{1}{2a} \right]$$

$$= \frac{\cos at}{2} \left[\frac{2 \sin at \cos at}{2a} + t \right] + \frac{\sin at}{2} \left[\frac{1 - \cos 2at}{2a} \right]$$

$$= \frac{\cos^2 at}{2a} [2 \sin at \cos at + t] + \frac{\cos at}{2} + \frac{\sin^2 at}{2}$$

$\cos^2 \theta = \frac{\cos 2\theta + 1}{2}$

$$f(s) = \frac{1}{s^n} \rightarrow F(t) = \frac{t^{n-1}}{(n-1)!}$$

(3) Problem. Find (i) $L^{-1} \left\{ \frac{1}{s(s+1)^3} \right\}$

(ii) $L^{-1} \left\{ \frac{1}{(s-2)(s+1)^2} \right\}$

— $f(s) = \frac{1}{s}$ and $g(s) = \frac{1}{(s+1)^3}$
 $\therefore L^{-1}\{f(s)\} = 1$ and $L^{-1}\{g(s)\} = \cos t$
 $= F(t) = 1$ and $G(t) = \frac{1}{2} e^{-t} t^2$

\therefore By convolution theorem

$$L^{-1}\{f(s)g(s)\} = \int_0^t F(u) \cdot G(t-u) du$$

$$= \int_0^t 1 \cdot \frac{1}{2} e^{-(t-u)} \cdot (t-u)^2 du$$

$$= \frac{1}{2} \int_0^t e^{-(t-u)} \cdot (u^2 + t^2 - 2tu) du$$

Let $t-u = p \Rightarrow du = -dp$
 at $u=0, p=t$
 at $u=t, p=0$

$$\therefore = -\frac{1}{2} \int_t^0 e^{-p} \cdot p^2 dp = \frac{1}{2} \int_0^t e^{-p} p^2 dp$$

$$= \frac{1}{2} \int_0^t e^{-p} p^2 dp = \frac{1}{2} \cdot \frac{1}{3} = \frac{1}{6}$$

$$= \left[\frac{1}{2} \right] \left[e^{-p} \cdot p^2 \right]$$

Application of Laplace Transform

The Laplace application, solⁿ of ODE with constant coefficient, solⁿ of PDE with constant and without and simultaneous diff eqn. Many more application in a classical mechanics, electrical circuit resolving solⁿ

Que: Using Laplace transform, determine the solⁿ

$$y'' + 3y' + 2y = e^t, \quad y(0) = y'(0) = 0$$

$$L[y'' + 3y' + 2y] = L[e^t]$$

$$\Rightarrow L[y''] + 3L[y'] + 2L[y] = L[e^t]$$

$$\Rightarrow \delta^2 L(y) - \delta y(0) - y'(0) + 3[\delta L(y)] - \delta y(0) + 2L(y) = \frac{1}{\delta + 1}$$

As given that $y(0) = y'(0) = 0$

$$\Rightarrow \delta^2 L(y) + 3\delta L(y) + 2L(y) = \frac{1}{\delta + 1}$$

$$\Rightarrow (\delta^2 + 3\delta + 2)L(y) = \frac{1}{\delta + 1}$$

$$\therefore L(y) = \frac{1}{(\delta + 1)(\delta^2 + 3\delta + 2)} = \frac{1}{(\delta + 1)(\delta + 1)(\delta + 2)}$$

$$L[y(t)] = \frac{1}{(\delta + 1)^2(\delta + 2)}$$

$$\therefore y(t) = L^{-1} \left[\frac{1}{(\delta + 1)^2(\delta + 2)} \right]$$

$$\text{let } f(s) = \frac{1}{s+2} \text{ \& } g(s) = \frac{1}{(s+1)^2}$$

$$\therefore L^{-1}(f(s)) = e^{-2t} \text{ \& } L^{-1}(g(s)) = e^{-t} \cdot t = F(t) = G(t)$$

\(\therefore\) By Convolution theorem

$$L^{-1}(f(s) \cdot g(s)) = \int_0^t e^{-2u} \cdot e^{-(t-u)} \cdot (t-u) du$$

$$= \int_0^t e^{-2u - t + u} \cdot (t-u) du$$

$$= \int_0^t e^{-t-u} \cdot (t-u) du$$

at $t \rightarrow u$

$$\left[\frac{e^{-(t+u)}}{-1} (t-u) + (t) \int_0^t \frac{e^{-(t+u)}}{-1} du \right]$$

$$= \left[e^{-(t+u)} (u-t) + \frac{e^{-(t+u)}}{-1} \right]_0^t$$

$$= \left[e^{-2t} (0) + \frac{e^{-2t}}{-1} - \left[e^{-t} (t-t) - \frac{e^{-t}}{-1} \right] \right]$$

$$= \left[+e^{-2t} + e^{-t} + e^{-t} \right]$$

$$= +e^{-2t} + (t-1)e^{-t} \quad \left(\frac{2t}{t+1} \right)$$

Que: Solve the DE

$$(\delta + 1)^2 y = t$$

given that

$$y(0) = -3 \text{ \& } y(1) = -1$$

next + Fourier series

Prepared by Kalika

$(D+1)^2 y = t, y(0) = -3, y(1) = -1$

$D^2 y + 2Dy + y = t$
i.e. $y'' + 2y' + y = t$

$\therefore L(y'' + 2y' + y) = L(t)$
By linearity property —
 $L(y'') + 2L(y') + L(y) = L(t)$

$\Rightarrow s^2 L(y) - sy(0) - y'(0) + 2[sL(y) - y'(0)] + L(y) = L(t)$

$\Rightarrow L(y)(s^2 + 2s + 1) - s(-3) - y'(0) - 2y'(0) = L(t)$

$\Rightarrow L(y)(s+1)^2 + 3s + 6 - y'(0) = L(t) = \frac{1}{s^2}$

Let $y'(0) = k$ (any const)

$L(y)(s+1)^2 = \frac{1}{s^2} - 3s - 6 + k$

$= \frac{1 - 3s^3 - 6s^2 + ks^2}{s^2}$

$L(y) = \frac{1 - 3s^3 - 6s^2 + ks^2}{s^2 (s+1)^2}$

$\frac{1}{s^2 (s+1)^2} = \frac{A}{s} + \frac{B}{s^2} + \frac{C}{s+1} + \frac{D}{(s+1)^2}$

$y = L^{-1} \left[\frac{1}{(s+1)^2 \cdot s^2} + \frac{3s+6+k}{(s+1)^2} \right]$

$= L^{-1} \left[\frac{1}{(s+1)^2 \cdot s^2} \right] + L^{-1} \left[\frac{3s+6+k}{(s+1)^2} \right]$

$y = L^{-1} \left[\frac{1}{(s+1)^2 \cdot s^2} \right] = L^{-1} \left[\frac{3}{(s+1)} + \frac{(3+k)}{(s+1)^2} \right]$

$= L^{-1} \left[\frac{1}{(s+1)^2 \cdot s^2} \right] - 3L^{-1} \left[\frac{1}{s+1} \right] - (3+k)L^{-1} \left[\frac{1}{(s+1)^2} \right]$

$(L f(s)) = \frac{1}{s^2}, g(s) = \frac{1}{(s+1)^2} = e^{-t}$

$\therefore L^{-1}[f(s)] = t = F(t)$

$L^{-1}[g(s)] = t \cdot e^{-t} = h(t)$

\therefore By convolution theorem —

$L^{-1}[f(s)g(s)] = \int_0^t F(u) \cdot h(t-u) du$

$I = \int_0^t [e^{-u} + t \cdot e^{-2u}] \cdot u \cdot e^{-4} du = \int_0^t (t-u) \cdot u \cdot e^{-4} du =$

$= \int_0^t \left[\frac{u^2 e^{-4}}{-1} + \frac{2 \cdot u e^{-4}}{-1} + \int_0^t e^{-4} \cdot 2 du \right] du = \int_0^t u^2 \cdot e^{-4} du$

$= \left[\frac{u^3 e^{-4}}{-3} - 2u e^{-4} - 2e^{-4} \right]_0^t = \frac{1}{3} = \frac{2}{3}$
 $= [-t^2 e^{-4} - 2t e^{-4} - 2e^{-4} + 2]$

$\therefore y = \frac{2}{3} - 3 \cdot e^{-t} - (3+k) \cdot t \cdot e^{-t}$

$\therefore y(1) = \frac{2}{3} - \frac{3}{e} - \frac{(3+k)}{e} = -1$

$t \int_0^t u \cdot e^{-4} du - \int_0^t u^2 \cdot e^{-4} du$

$\int_0^t [u \cdot e^{-4}] du + \int_0^t e^{-4} du - \int_0^t u^2 \left(\frac{e^{-4}}{-3} \right) - \int_0^t 2u \cdot e^{-4} du$

$= e^{-t} y = 1 - t + t^2 e^{-t} + t^3 e^{-t} + (3+k)$

Let $L^{-1} \left[\int_0^t F(u) G(t-u) du \right] = f(s) g(s)$

(I) Prove that $\int_0^t \sin u \cos(t-u) du = \frac{1}{2} t \sin t$.
using the convolution theorem.

Proof: Let $F(u) = \sin u$, $G(u) = \cos(t-u)$
given that $F(t) = \sin t$, $G(t) = \cos t$

$\therefore L[F(t)] = \frac{1}{s^2+1} = f(s)$

R $L[G(t)] = L[\cos t] = \frac{s}{s^2+1} = g(s)$

\therefore By convolution theorem —

$\therefore L^{-1} [f(s) g(s)] = \int_0^t F(t-u) G(u) du$

$\int_0^t F(t-u) G(u) du = L^{-1} \left[\frac{1}{s^2+1} \cdot \frac{s}{s^2+1} \right]$

$L^{-1} \left[\frac{1}{s^2+1} \cdot \frac{s}{s^2+1} \right] = \int_0^t \frac{1}{2} [\sin(t+u) + \sin(u-t)] du$

$= \frac{1}{2} \int_0^t (\sin t + \sin 2u - 1) du$

$= \frac{1}{2} t \sin t + \left[-\frac{\cos 2u}{2} \right]_0^t$

$+ \frac{1}{2} [-\cos t + \cos t]$

$= \frac{1}{2} t \sin t$ Δ

II. If $L^{-1} [f(s)] = F(t)$ then using convolution theorem, P.T.

$L^{-1} \left[\frac{f(s)}{(s-a)^2 + b^2} \right] = \frac{1}{b} e^{+at} \int_0^t f(z) e^{-az} \sin b(t-z) dz$

III. Solve the following Simultaneous eqⁿ.

(140) (i) $\frac{dx}{dt} = y + e^t$

(ii) $\frac{dy}{dt} + x = \sin t$. The IC $x(0) = 1$ & $y(0) = 0$

(i) $\frac{dx}{dt} + 5x - 2y = t$

(ii) $\frac{dy}{dt} + 2x + y = 0$ when $x=y=0$ then $t=0$.

* Dirichlet Conditions

* Fourier Integral Theorem:-

Consider a fⁿ f(x) which satisfies the Dirichlet's conditions in every interval $(-c, c)$, so, that

$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left(a_n \cos \frac{n\pi x}{c} + b_n \sin \frac{n\pi x}{c} \right)$ — (1)

Dirichlet Condition

(I). The fⁿ must be periodic.

(II). It must be single value and continuous, except possibly at a finite no. of finite discontinuity.

(III). It must have only a finite no. of maxima and minima within one periodic.

(IV). The integral over one period of $|f(x)|$ must be converge.

where

where $a_0 = \frac{1}{c} \int_{-c}^c f(t) dt$
 $a_n = \frac{1}{c} \int_{-c}^c f(t) \cos \frac{n\pi t}{c} dt$
 $b_n = \frac{1}{c} \int_{-c}^c f(t) \sin \frac{n\pi t}{c} dt$ (1)

Now, using eqn (1) and (2), we have

$$f(x) = \frac{1}{2c} \int_{-c}^c f(t) dt + \frac{1}{c} \sum_{n=1}^{\infty} \int_{-c}^c f(t) \cos \frac{n\pi(t-x)}{c} dt$$
 (3)

If we assume that $\int_{-\infty}^{\infty} |f(x)| dx$ converges, the first term on the right side of eqn (3) approaches 0 as $c \rightarrow \infty$. Since

$$\left| \frac{1}{2c} \int_{-c}^c f(t) dt \right| \leq \frac{1}{2c} \int_{-\infty}^{\infty} |f(t)| dt$$

the second term

$$\lim_{c \rightarrow \infty} \frac{1}{c} \sum_{n=1}^{\infty} \int_{-c}^c f(t) \cos \frac{n\pi(t-x)}{c} dt$$

Now putting $dt = \pi/c$

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$$\lim_{\delta \rightarrow 0} \frac{1}{\pi} \sum_{n=1}^{\infty} \delta \int_{-\infty}^{\infty} f(t) \cos n\delta(t-x) dt$$

This is of the form $\lim_{\delta \rightarrow 0} \sum_{n=1}^{\infty} F(n\delta)$ i.e. $\int_{-\infty}^{\infty} F(t) dt$

Thus as $c \rightarrow \infty$, (3) becomes

$$f(x) = \frac{1}{\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \cos \pi \delta(t-x) d\delta dt$$

which is known as the Fourier integral of $f(x)$.

Note: Note that, we have given a heuristic (involves, discovery and problems) demonstration of Fourier integral which help is deriving the result eqn(3). & when $f(x)$ at a pt. discontinuity, we replace $f(x)$ by $\frac{1}{2} [f(x+0) + f(x-0)]$ as in the case of Fourier series.

(1) For Fourier sine & Cosine integral expanding $\cos \pi \delta(t-x)$ may be written as

$$f(x) = \frac{1}{\pi} \left[\int_0^{\infty} \cos \pi x \left(\int_{-\infty}^{\infty} f(t) \cos \pi \delta t dt \right) d\delta + \int_0^{\infty} \sin \pi x \left(\int_{-\infty}^{\infty} f(t) \sin \pi \delta t dt \right) d\delta \right]$$
 (4)

(2) If $f(x)$ is an odd fn, $f(t) \cos \pi \delta t$ is also an odd fn while $f(t) \sin \pi \delta t$ is even then the second term on the right side of eqn(4) vanishes, and we get

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \sin \pi x \left[\int_{-\infty}^{\infty} f(t) \sin \pi \delta t dt \right] d\delta$$
 (5)

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \cos \pi x \left[\int_{-\infty}^{\infty} f(t) \cos \pi \delta t dt \right] d\delta$$
 (6)

Now, eqn (5) & (6) are known as Fourier sine integral & Fourier Cosine integral resp.

* A fn $f(x)$ defined in the interval $(0, \infty)$ is expressed either as a Fourier sine integral or as a Fourier Cosine integral.

Simply looking upon it as an odd or even fⁿ in $(-\infty, \infty)$ on the lines of half range Fourier series

(I) For complex form of Fourier integral eq (3) can be written as

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \cos \lambda(t-x) dt d\lambda \quad (7)$$

because $\cos \lambda(t-x)$ is an even fⁿ λ , Also $\sin \lambda(t-x)$ is an odd fⁿ of λ .
~~Now eq (7) + eq (8)~~, we have

$$0 = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \sin \lambda(t-x) dt d\lambda \quad (8)$$

Now, eq (7) + i eq (8) on both sides (side by side)

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) [\cos \lambda(t-x) + i \sin \lambda(t-x)] dt d\lambda$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) e^{i\lambda(t-x)} dt d\lambda \quad (9)$$

which is known as complex form of Fourier integral.

(II) For Fourier integral represent of a example fⁿ :-

$$f(x) = \begin{cases} 1 & \text{for } 0 < x < \pi \\ 0 & \text{for } x > \pi \end{cases}$$

By a Fourier sine integral and hence

$$\text{Evaluate } \int_0^{\infty} (1 - \cos \lambda t) \sin \lambda x d\lambda$$

Self:

The Fourier sine integral for $f(x) = \frac{2}{\pi} \int_0^{\infty} \sin \lambda(x) d\lambda \int_0^{\infty} f(t) \sin \lambda t dt$

$$= \frac{2}{\pi} \int_0^{\infty} \sin \lambda(x) \left[\int_0^{\pi} \sin \lambda t \cdot 1 dt \right] d\lambda$$

$$= \frac{2}{\pi} \int_0^{\infty} \sin \lambda x \left[\frac{-\cos \lambda t}{\lambda} \right]_0^{\pi} d\lambda$$

$$= \frac{2}{\pi} \int_0^{\infty} \frac{\sin \lambda x}{\lambda} [-1 + \cos \lambda \pi] d\lambda$$

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \frac{(1 - \cos \lambda \pi)}{\lambda} \sin \lambda x d\lambda$$

$$\Rightarrow \int_0^{\infty} \frac{(1 - \cos \lambda \pi)}{\lambda} \sin \lambda x d\lambda = \frac{\pi}{2} f(x)$$

$$= \frac{\pi}{2} \begin{cases} 1 & 0 < x < \pi \\ 0 & x > \pi \end{cases}$$

$$= \begin{cases} \pi/2 & 0 < x < \pi \\ 0 & x > \pi \end{cases}$$

at $x = \pi$, which is a pt. of discontinuity of $f(x)$, the value of the above integral

$$= \frac{\pi}{2} \left[\frac{f(\pi+0) + f(\pi-0)}{2} \right] = \frac{\pi}{4}$$

For Fourier Integral Representation of a fⁿ :-

Using (4), if a function $f(x)$, may be represented by a Fourier integral as

$$f(x) = \frac{1}{\pi} \int_{-\infty}^{\infty} \cos \lambda x \left(\int_{-\infty}^{\infty} f(t) \cos \lambda t dt \right) d\lambda + \int_0^{\infty} \sin \lambda x \left(\int_{-\infty}^{\infty} f(t) \sin \lambda t dt \right) d\lambda$$

$$= \frac{1}{\pi} \int_0^{\infty} A(\omega) \cos \omega x d\omega + \frac{1}{\pi} \int_0^{\infty} B(\omega) \sin \omega x d\omega$$

$$= \frac{1}{\pi} \left[\int_0^{\infty} A(\omega) \cos \omega x d\omega + \int_0^{\infty} B(\omega) \sin \omega x d\omega \right]$$

where

$$A(\omega) = \int_{-\infty}^{\infty} f(t) \cos \omega t dt = 2 \int_0^{\infty} f(t) \cos \omega t dt$$

$$B(\omega) = \int_{-\infty}^{\infty} f(t) \sin \omega t dt = 2 \int_0^{\infty} f(t) \sin \omega t dt$$

if $f(x)$ is an odd fⁿ then eqⁿ (5) is

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \sin \omega x \left[\int_0^{\infty} f(t) \sin \omega t dt \right] d\omega$$

$$f(x) = \frac{1}{\pi} \int_0^{\infty} B(\omega) \sin \omega x d\omega$$

$$f(x) = \frac{1}{\pi} \int_0^{\infty} A(\omega) \sin \omega x d\omega$$

Eqⁿ (9)

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \cdot e^{i\omega(t-x)} dt d\omega \quad \text{Eqⁿ (9)}$$

Now putting $\omega = \phi$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \cdot e^{i\phi(t-x)} dt d\phi$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(t) \cdot e^{i\phi t} \cdot e^{-i\phi x} d\phi dt$$

$$\left[F(\phi) = \int_{-\infty}^{\infty} f(t) e^{i\phi t} dt \right] \quad \text{(10)}$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\phi) e^{-i\phi x} d\phi \quad \text{(11)}$$

The $F(s)$ is defined by

$$F(s) = \int_{-i\infty}^{\infty} f(t) \cdot e^{ist} dt$$

Also this is called Fourier pair form.

Also the fⁿ $f(x)$ as given by eqⁿ (11) is called Inverse Fourier transform of $F(s)$.

For Fourier Sine & Cosine transform

Using eqⁿ (5),

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \sin \omega x \left[\int_0^{\infty} f(t) \sin \omega t dt \right] d\omega$$

define the fⁿ $F_s(\omega) = \int_0^{\infty} f(t) \sin \omega t dt$

$$f(x) = \frac{2}{\pi} \int_0^{\infty} F_s(\omega) \cdot \sin \omega x d\omega$$

$F_s(s)$
↑
for sine

And for cosine, using (6) —

$$f(x) = \frac{2}{\pi} \int_0^{\infty} \cos \omega x \left[\int_0^{\infty} f(t) \cos \omega t dt \right] d\omega$$

define the fⁿ $F_c(\omega) = \int_0^{\infty} f(t) \cos \omega t dt$

$$f(x) = \frac{2}{\pi} \int_0^{\infty} F_c(\omega) \cdot \cos \omega x d\omega$$

Example:
(I)

Find the Fourier transform of $f(x)$

$$f(x) = \begin{cases} 1-x^2 & |x| \leq 1 \\ 0 & |x| > 1 \end{cases}$$

Hence evaluate $\int_0^{\infty} (x \cos x - \sin x) \cos x dx$

* * * next - modulation

8/2/16

No class due to meeting with vc (mutual seminar)

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Example II

Division in last example, Evaluate the following

$$\int_0^{\infty} \frac{(x \cos x - \sin x)}{x^3} \sin \frac{x}{2} dx$$

Solⁿ * Here, $f(s) = \int_{-\infty}^{\infty} f(x) e^{isx} dx$ - Fourier transform

$$f(s) = \int_{-\infty}^{-1} 0 \cdot dx + \int_{-1}^1 (1-x^2) e^{isx} dx + \int_1^{\infty} 0 \cdot dx$$

$$= \int_{-1}^1 e^{isx} dx - \int_{-1}^1 x^2 e^{isx} dx$$

$$= \left[\frac{e^{isx}}{is} \right]_{-1}^1 - \left[\frac{x^2 e^{isx}}{is} - \int \frac{e^{isx}}{is} \cdot 2x dx \right]$$

$$= \left[\frac{e^{is} - e^{-is}}{is} \right] - \left[\frac{x^2 e^{isx}}{is} - \frac{e^{isx} \cdot 2x}{(is)^2} + \frac{e^{isx} \cdot 2}{(is)^2} \right]_{-1}^1$$

$$= \left[\frac{e^{is} - e^{-is}}{is} \right] - \left[\frac{x^2 e^{isx}}{is} - \frac{e^{isx} \cdot 2x}{(is)^2} + \frac{2e^{isx}}{(is)^2} \right]_{-1}^1$$

$$= \left[\frac{e^{is} - e^{-is}}{is} \right] - \left[\frac{e^{is} - e^{-is}}{is} - \frac{2(e^{is} + e^{-is})}{(is)^2} + \frac{2(e^{is} - e^{-is})}{(is)^2} \right]$$

$$= \frac{2 \cdot 2 (e^{is} + e^{-is})}{2(is)^2} - \frac{2 (e^{is} - e^{-is})}{(is)^2}$$

$$= -\frac{4 \cos 2s}{s^2} + \frac{4 \sin 2s}{s^3}$$

$$\therefore f(s) = \frac{4 \sin 2s}{s^3} - \frac{4 \cos 2s}{s^2}$$

Second part

Now evaluating $f(x)$:

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(s) e^{-isx} ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^2} (s \cos s - \sin s) e^{-isx} ds$$

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$$= (-1) e^{i0} = \cos 0 + i \sin 0$$

$$\therefore f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^3} (s \cos s - \sin s) (\cos sx - i \sin sx) ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^3} [s \cos s \cdot \cos sx - \sin s \cdot \cos sx - i (s \cos s \cdot \sin sx - \sin s \cdot \sin sx)] ds$$

Now, here term with i , that is we take only real part

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^3} [s \cos s \cos sx - \sin s \cdot \cos sx]$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^3} [s \cos s - \sin s] \cos sx ds$$

putting $x = 1/2$, we have

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} -\frac{4}{s^3} [s \cos s - \sin s] \cos \frac{s}{2} ds$$

$$\frac{3}{4} = f(x) = f(x) = -\frac{4}{2\pi} \int_{-\infty}^{\infty} \frac{(s \cos s - \sin s) \cos \frac{s}{2}}{s^3} ds$$

$$\frac{3}{4} \times \frac{\pi}{-4} = \int_0^{\infty} \frac{s \cos s - \sin s}{s^3} \cos \frac{s}{2} ds$$

$$\int_0^{\infty} \frac{s \cos s - \sin s}{s^3} \cos \frac{s}{2} ds = \frac{3\pi}{-16}$$

HW

* Properties of Fourier Transform

Sol: (I)

Linearity Property:

If $f(x)$ & $g(x)$ are Fourier transforms of $f(x)$ & $g(x)$ respectively, then

$$F[af(x) + bg(x)] = af(s) + bg(s)$$

i.e. $F[af(x) + bg(x)] = aF[f(x)] + bF[g(x)]$
 $= a f(s) + b g(s)$
 where a, b are constants

Pf: We know that

$$f(s) = F[f(x)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{isx} dx$$

$$g(s) = F[g(x)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} g(x) \cdot e^{isx} dx$$

$$\therefore F[af(x) + bg(x)] = \frac{1}{2\pi} \int_{-\infty}^{\infty} (af(x) + bg(x)) e^{isx} dx$$

$$= \frac{1}{2\pi} \left[\int_{-\infty}^{\infty} af(x) \cdot e^{isx} dx + \int_{-\infty}^{\infty} bg(x) \cdot e^{isx} dx \right]$$

$$= \frac{a}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{isx} dx + \frac{b}{2\pi} \int_{-\infty}^{\infty} g(x) \cdot e^{isx} dx$$

$$= a \cdot f(s) + b \cdot g(s)$$

$\Rightarrow F[af(x) + bg(x)] = a \cdot f(s) + b \cdot g(s)$ (Proved)

Sol: (II)

Change of Scale Property:

If $f(s)$ is the Fourier transform of

$f(x)$, then the Fourier transform of $f(ax) = \frac{1}{|a|} f\left(\frac{s}{a}\right)$, $a \neq 0$

Pf:

i.e. $F\{f(ax)\} = \frac{1}{|a|} \cdot f\left(\frac{s}{a}\right) = f\left(\frac{s}{a}\right)$

As we know that

$$F(s) = F(f(x)) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{isx} dx$$

$$\therefore F(f(ax)) = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(ax) \cdot e^{isx} dx$$

Now putting $ax = t \Rightarrow dx = \frac{dt}{a}$

$$F = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) \cdot e^{i\left(\frac{s}{a}\right)t} \frac{dt}{a}$$

$$= \frac{1}{a} \cdot \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) \cdot e^{i\left(\frac{s}{a}\right)t} dt$$

$$= \frac{1}{a} \cdot f\left(\frac{s}{a}\right)$$

$\Rightarrow F\{f(ax)\} = \frac{1}{|a|} f\left(\frac{s}{a}\right)$, $a \neq 0$

Sol: 3.

Shifting Property:

If $f(s)$ is the Fourier transform of $f(x)$, then the Fourier transform of $f(x-a) = e^{isa} \cdot f(s)$

i.e. $F\{f(x-a)\} = e^{isa} \cdot f(s)$

As we know that

$$F\{f(x)\} = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{isx} dx = f(s)$$

$$F\{f(x-a)\} = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x-a) e^{isx} dx$$

let $x-a = t \Rightarrow dx = dt$
 $\Rightarrow x = t+a$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) \cdot e^{is(t+a)} dt$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) e^{ist} \cdot e^{isa} dt$$

$$= e^{isa} \cdot \left[\frac{1}{2\pi} \int_{-\infty}^{\infty} f(t) \cdot e^{ist} dt \right]$$

$F\{f(x-a)\} = e^{isa} \cdot f(s)$ = $f(s)$
(proff)

Note: If $F_s(x)$ and $F_c(x)$ are the Fourier transform of the Fourier sine & Fourier cosine transform of $f(x)$ respectively, then,

$$F_s\{f(x)\} = \frac{1}{a} F_s\left(\frac{s}{a}\right) \quad \text{and}$$

$$F_c\{f(x)\} = \frac{1}{a} F_c\left(\frac{s}{a}\right)$$

(4) Modulation theorem:

If $f(s)$ is the complex Fourier transform of $f(x)$, then the Fourier transform of $f(x) \cos ax = \frac{1}{2} [f(s-a) + f(s+a)]$

$$F\{f(x) \cos ax\} = \frac{1}{2} [f(s-a) + f(s+a)]$$

self proff

We know that

$$F\{f(x)\} = \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{isx} dx$$

$$\therefore \cos ax = \frac{e^{iax} + e^{-iax}}{2}$$

$$\therefore F\{f(x) \cos ax\} = \frac{1}{2} \int_{-\infty}^{\infty} f(x) \cdot \left(\frac{e^{iax} + e^{-iax}}{2} \right) dx$$

$$= \frac{1}{2} \left[\int_{-\infty}^{\infty} f(x) \cdot \frac{e^{iax} \cdot e^{isx}}{2} dx + \int_{-\infty}^{\infty} f(x) \cdot \frac{e^{-iax} \cdot e^{isx}}{2} dx \right]$$

$$= \frac{1}{2} \left[\frac{1}{2} \int_{-\infty}^{\infty} f(x) \cdot e^{i(s+a)x} dx + \frac{1}{2} \int_{-\infty}^{\infty} f(x) \cdot e^{i(s-a)x} dx \right]$$

$$= \frac{1}{2} \left[\frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{ik_1x} dx + \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) \cdot e^{ik_2x} dx \right]$$

where $k_1 = s+a, k_2 = s-a$

$$= \frac{1}{2} [f(k_1) + f(k_2)]$$

$$= \frac{1}{2} [f(s+a) + f(s-a)]$$

$$\Rightarrow F\{f(x) \cos ax\} = \frac{1}{2} [f(s+a) + f(s-a)]$$

Note: (I) $F_s(x)$ & $F_c(x)$ are Fourier sine & cosine transform of $f(x)$

$$F_s\{f(x) \sin ax\} = \frac{1}{2} [F_s(s+a) + F_s(s-a)]$$

$$F_c\{f(x) \sin ax\} = \frac{1}{2} [F_c(s+a) + F_c(s-a)]$$

$$F_s\{f(x) \cos ax\} = \frac{1}{2} [F_s(s+a) + F_s(s-a)]$$

$$F_c\{f(x) \cos ax\} = \frac{1}{2} [F_c(s+a) + F_c(s-a)]$$

Examples

Find the Fourier sine transform & cosine transform of $x^n, n > 0$.

As we know that

$$f_c(\omega) = F_c(f(\omega)) = \int_0^{\infty} f(t) \sin \omega t dt$$

$$= \int_0^{\infty} x^{n+1} \sin \omega t dt \quad (2)$$

$$f_c(\omega) = F_c(f(\omega)) = \int_0^{\infty} x^n \cdot \cos \omega t dt \quad (5)$$

$$\int_0^{\infty} e^{-x} x^{n+1} dx \quad (1)$$

$$\left(\because \sin \omega t = \frac{e^{i\omega t} - e^{-i\omega t}}{2i} \quad \& \quad \cos \omega t = \frac{e^{i\omega t} + e^{-i\omega t}}{2} \right)$$

gn (1), putting

$$dx = i\omega dy \quad \Rightarrow \quad x = i\omega y \quad \& \quad i = \omega^{1/2}$$

$$\Gamma_n = \int_0^{\infty} e^{-i\omega y} (i\omega y)^{n+1} \cdot i\omega dy$$

$$= \int_0^{\infty} e^{-i\omega y} y^{n+1} \cdot (i\omega)^{n+2} dy$$

$$= (i\omega)^{n+2} \int_0^{\infty} e^{-i\omega y} y^{n+1} dy$$

$$= \frac{(-1)^{n+1}}{i^{n+2}} \int_0^{\infty} e^{-i\omega y} y^{n+1} dy$$

Find the Fourier sine & cosine transform of x^{n+1} , $0 < x < \infty$.

$$\cos 2\theta = 2\cos^2\theta - 1$$

Example: Solve the integral eqn

$$\int_0^{\infty} f_c(\omega) \cdot \cos \alpha \omega \cdot d\omega = \begin{cases} 1-x & 0 \leq x \leq 1 \\ 0 & x > 1 \end{cases}$$

Evaluate $\int_0^{\infty} \frac{\sin^2 t}{t^2} dt = \frac{\pi}{2}$

$$f_c(x) = \int_0^{\infty} f_c(\omega) \cdot \cos \alpha \omega \cdot d\omega = \int_0^{\infty} f_c(\omega) \cdot \frac{e^{i\alpha\omega} + e^{-i\alpha\omega}}{2} d\omega$$

By the Inversion formula, we have

$$f_c(x) = \frac{2}{\pi} \int_0^{\infty} F_c(\omega) \cdot \cos \alpha \omega \cdot d\omega$$

$$= \frac{2}{\pi} \int_0^1 (1-x) \cos \alpha \omega \cdot d\omega + \int_0^{\infty} 0 \cdot \cos \alpha \omega \cdot d\omega$$

$$= \frac{2}{\pi} \int_0^1 \cos \alpha \omega \cdot d\omega - \frac{2}{\pi} \int_0^1 x \cos \alpha \omega \cdot d\omega$$

$$= \frac{2}{\pi} \left[\left[\frac{\sin \alpha \omega}{\alpha} \right]_0^1 - \left[x \cdot \frac{\sin \alpha \omega}{\alpha} - \int \sin \alpha \omega \cdot dx \right]_0^1 \right]$$

$$= \frac{2}{\pi} \left[\frac{\sin \alpha}{\alpha} - \left[x \sin \alpha \omega + \frac{\cos \alpha \omega}{\alpha^2} \right]_0^1 \right]$$

$$= \frac{2}{\pi} \left[\frac{\sin \alpha}{\alpha} - \left[\sin \alpha + \cos \alpha - 1 \right] \right]$$

$$= \frac{2}{\pi} \left[1 - \cos \alpha \right] = \frac{2}{\pi} \left[2(1 - \cos^2 \frac{\alpha}{2}) \right]$$

$$= \frac{4}{\pi} \sin^2 \frac{\alpha}{2}$$

$$\int_0^{\infty} \frac{\sin^2 t}{t^2} dt = \frac{\pi}{2}$$

Prepared by Kalika

Internal

4 Given that: $\frac{\partial^4 y}{\partial x^4} = 3 \frac{\partial^2 y}{\partial x^2}$ — (1)

where $y(x, 0) = 0, \left(\frac{\partial y}{\partial x}\right)_{x=0} = 0$

$y(x, 0) = 30 \cos 5x$

In eqn (1) taking h.T on the both sides, we have

$\bar{y}(x, s) = Y(x, s) = 3 \frac{d^2 \bar{y}}{dx^2}$

$\frac{\bar{y}(x, s) - 30 \cos 5x}{3} = \frac{d^2 \bar{y}}{dx^2}$

$(s^2 - 5/3) \bar{y}(x, s) = -10 \cos 5x$

$f(s) \bar{y}(x, s) = X$ — (2)

Case-I, $x=0$
Case-II, $x \neq 0$

$\bar{y}(x, s) = C_1 e^{\sqrt{5/3} x} + C_2 e^{-\sqrt{5/3} x} + \frac{30 \cos 5x}{s^2 - 5/3}$

$y(x, t) = +30 e^{-5t} \cos 5x$

(2)

$L[y(x)] = 2 \int_0^x y(t) \cdot e^{-st} dt$
 $L[y(x)] = 2 [1] - L \left[\int_0^x y(t) \cdot e^{-st} dt \right]$
 $= \frac{2}{s} - L[F(x)] L[h(x)]$
 $= \frac{2}{s} -$

$F(x) = y(x)$
 $h(x) = e^{-st}$

Convolution:

The convolution of two function $f(x)$ & $g(x)$ over the interval $(-\infty, \infty)$ is defined as $f * g = \int_{-\infty}^{\infty} f(u)g(x-u) du$

$f * g = \int_{-\infty}^{\infty} f(u)g(x-u) du = h(x)$

$\Rightarrow \boxed{f * g = h(x)}$

Theorem: Convolution theorem for fourier transform

The fourier transform of the convolution of $f(x)$ and $g(x)$ is the product of their Fourier Transform.

$F[f(x) * g(x)] = F[f(x)] \cdot F[g(x)]$

Proof:

We know that Fourier transform

$F[f(x)] = \int_{-\infty}^{\infty} f(x) e^{isx} dx$

also let fourier transform of $g(x)$ be similar

so $F[f(x) * g(x)] = \int_{-\infty}^{\infty} (f(x) * g(x)) e^{isx} dx$

$= \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(u) * g(x-u) du \right] e^{isx} dx$

$= \int_{-\infty}^{\infty} f(u) \left[\int_{-\infty}^{\infty} g(x-u) e^{isx} dx \right] du$

(∵ changing the order of integration)

$= \int_{-\infty}^{\infty} f(u) \left[\int_{-\infty}^{\infty} e^{is(x-u)} \cdot g(x-u) dx \right] du$

$= \int_{-\infty}^{\infty} f(u) \cdot e^{isu} du * \int_{-\infty}^{\infty} e^{is(x-u)} g(x-u) dx$

$$F[f(x) \cdot g(x)] = F[f(x)] * F[g(x)]$$

* Parseval's Identity for Fourier transform

If the Fourier transform of $f(x)$ and $g(x)$ are respectively $F(s)$ and $\bar{G}(s)$ then

$$(i) \frac{1}{2\pi} \int_{-\infty}^{\infty} F(s) \bar{G}(s) ds = \int_{-\infty}^{\infty} f(x) g(x) dx$$

$$(ii) \frac{1}{2\pi} \int_{-\infty}^{\infty} [F(s)]^2 ds = \int_{-\infty}^{\infty} |f(x)|^2 dx$$

where bar implies the complex conjugate.

Verify

Proof: We have $\int_{-\infty}^{\infty} f(x) g(x) dx = \int_{-\infty}^{\infty} f(x) \left\{ \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{G}(s) e^{isx} ds \right\} dx$

We know that

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(s) e^{-isx} ds$$

$$\bar{f}(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{F}(s) e^{isx} ds$$

$$\therefore \int_{-\infty}^{\infty} f(x) \bar{g}(x) dx = \int_{-\infty}^{\infty} f(x) \left\{ \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{G}(s) e^{isx} ds \right\} dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{G}(s) \left\{ \int_{-\infty}^{\infty} f(x) e^{isx} dx \right\} ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{G}(s) \cdot F(s) ds$$

$$\int_{-\infty}^{\infty} f(x) g(x) dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(s) \bar{G}(s) ds$$

Tuesday

Proof: (11)

$$f(x) \cdot \bar{f}(x) = |f(x)|^2$$

$$\therefore \int_{-\infty}^{\infty} |f(x)|^2 dx = \int_{-\infty}^{\infty} f(x) \bar{f}(x) dx$$

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(s) e^{-isx} ds$$

$$\therefore \int_{-\infty}^{\infty} f(x) \bar{f}(x) dx = \int_{-\infty}^{\infty} f(x) \left\{ \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{F}(s) e^{isx} ds \right\} dx$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{F}(s) \left\{ \int_{-\infty}^{\infty} f(x) e^{isx} dx \right\} ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} \bar{F}(s) \cdot F(s) ds$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} [F(s)]^2 ds$$

$$\bar{F}(0) \cdot F(0) = |F(0)|^2$$

$$\int_{-\infty}^{\infty} |f(x)|^2 dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} [F(s)]^2 ds$$

NB: Note that: The following Parseval's identity for Fourier cosine & sine transform

$$(i) \frac{2}{\pi} \int_0^{\infty} F_c(s) G_c(s) ds = \int_0^{\infty} f(x) g(x) dx$$

$$(ii) \frac{2}{\pi} \int_0^{\infty} F_s(s) G_s(s) ds = \int_0^{\infty} f(x) g(x) dx$$

$$(iii) \frac{2}{\pi} \int_0^{\infty} [F_c(s)]^2 ds = \int_0^{\infty} [f(x)]^2 dx$$

$$(iv) \frac{2}{\pi} \int_0^{\infty} [F_s(s)]^2 ds = \int_0^{\infty} [f(x)]^2 dx$$

$$\frac{\partial^4 u}{\partial x^2} = \frac{\partial^2 (u_x)}{\partial x^2}$$

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II Fourier transform of the derivatives of function. The Fourier transform $F[u(x,t)]$ is defined as

$$F[u(x,t)] = \int_{-\infty}^{\infty} u e^{isx} dx \quad (1)$$

Now taking the derivative from $\frac{\partial^2 u}{\partial x^2}$

Replacing u by $\frac{\partial^2 u}{\partial x^2}$ in (1) we have

$$F\left[\frac{\partial^2 u}{\partial x^2}\right] = \int_{-\infty}^{\infty} \frac{\partial^2 u}{\partial x^2} \cdot e^{isx} dx$$

$$= \left[e^{isx} \frac{\partial u}{\partial x} - \int \frac{\partial u}{\partial x} \cdot e^{isx} \cdot is dx \right]_{-\infty}^{\infty}$$

How use
L.H.F.E

$\frac{\partial u}{\partial x} \rightarrow 0$
 $x \rightarrow \pm \infty$

$$= \left[e^{isx} \frac{\partial u}{\partial x} - is \int u \cdot e^{isx} dx \right]_{-\infty}^{\infty}$$

on applying limits of integration
Here, as $x \rightarrow \pm \infty$ then $\frac{\partial u}{\partial x} \rightarrow 0$, then

$$F\left[\frac{\partial^2 u}{\partial x^2}\right] = -s^2 F$$

Similarly, in the case of Fourier sine and cosine transform, we have

$$F_s\left[\frac{\partial^2 u}{\partial x^2}\right] = s(u)_{x=0} - s^2 F_s(u)$$

$$F_c\left[\frac{\partial^2 u}{\partial x^2}\right] = -\left(\frac{\partial u}{\partial x}\right)_{x=0} + s^2 F_c(u)$$

In general, the Fourier transform of the n^{th} derivative of $f(x)$ is given by

$$F\left[\frac{d^n f(x)}{dx^n}\right] = (-is)^n F[f(x)]$$

* * *

Friday

Application of Fourier Transform

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Application of Fourier transform for BVP

In one-dim BVP, the PDE can be easily transformed into ODE by applying a suitable transform. The required solⁿ is obtained by solving the eqⁿ and inverting by the complex inversion formula or any other method.

(1) If in a problem $u(x,t)$ or $u(x,0)$ is given, then we use infinite sine transform to remove $\frac{\partial^2 u}{\partial x^2}$ from the DE.

In case $\left[\frac{\partial u(x,t)}{\partial x}\right]_{x=0}$ is given then we

employ infinite Fourier cosine transform to remove $\frac{\partial^2 u}{\partial x^2}$

(2) If in a problem $u(0,t)$ & $u(l,t)$ are given, then we use finite sine transform to remove $\frac{\partial^2 u}{\partial x^2}$ from the DE.

In case $\left(\frac{\partial u}{\partial x}\right)_{x=0}$ and $\left(\frac{\partial u}{\partial x}\right)_{x=l}$ are given,

we employ finite Fourier cosine transform to remove $\frac{\partial^2 u}{\partial x^2}$.

Solve, $\frac{\partial u}{\partial t} = 2 \frac{\partial^2 u}{\partial x^2}$ if $u(0,t) = 0$,

$u(x,0) = e^{-x}$ ($x > 0$), $u(x,t)$ is b.d.f where $x > 0, t > 0$

Given that -

$$\frac{\partial u}{\partial t} = 2 \frac{\partial^2 u}{\partial x^2}, \quad x > 0, t > 0 \quad (1)$$

B.C $u(0,t) = 0$, $u(x) \text{ is b.d.f}$
I.C $u(x,0) = e^{-x}, x > 0$

(2)
(8)

∵ $u(x,0)$ is given then we take Fourier sine transform

$$\int_0^{\infty} \frac{\partial u}{\partial t} \sin px dx = 2 \int_0^{\infty} \frac{\partial^2 u}{\partial x^2} \sin px dx$$

$$\frac{d}{dt} \int_0^{\infty} u \sin px dx = 2 \left[\frac{\partial u}{\partial x} \sin px \Big|_0^{\infty} - \int_0^{\infty} \frac{\partial u}{\partial x} p \cos px dx \right]$$

$$\left(\bar{u}_s^2 \int_0^{\infty} u \sin px dx \right)$$

$$\frac{d\bar{u}_s}{dt} = 2 \left[0 - \int_0^{\infty} \frac{\partial u}{\partial x} p \cos px dx \right]$$

$$\frac{d\bar{u}_s}{dt} = -2p \left[u(x,t) \cos px \Big|_0^{\infty} - \int_0^{\infty} u(x,t) \{-p\} \sin px dx \right]$$

$$\frac{d\bar{u}_s}{dt} = -2p^2 \int_0^{\infty} u(x,t) \sin px dx$$

$$\frac{d\bar{u}_s}{dt} = -2p^2 \bar{u}_s$$

Integrating on the both sides, we have

$$\log \left(\frac{\bar{u}_s}{c} \right) = -2p^2 t$$

$$\bar{u}_s = c e^{-2p^2 t}$$

In eqn (3), $\int_0^{\infty} u(x,0) \sin px dx = \int_0^{\infty} e^{-x} \sin px dx$

$$\bar{u}_s(x,0) = \int_0^{\infty} e^{-x} \sin px dx$$

$$\bar{u}_s(p,0) = \int_0^{\infty} e^{-x} \sin px dx$$

$$\bar{u}_s(p,0) = \frac{p}{1+p^2} \quad \text{--- (5)}$$

In eqn (5), putting $c = \frac{p}{1+p^2}$

then,
$$\bar{u}_s(p,t) = \frac{p}{1+p^2} e^{-2p^2 t}$$

$$u(x,t) = \frac{2}{\pi} \int_0^{\infty} \left(\frac{p}{1+p^2} \right) e^{-2p^2 t} \sin px dp$$

Example: -1 Solve the eqn

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad (x > 0, u > 0)$$

Subject to the condition

(i) $u = 0$, where $x = 0, t > 0$

(ii) $u = \begin{cases} 1 & 0 < x < 1 \\ 0 & x \geq 1 \end{cases}$, where $t = 0$

(iii) $u(x,t)$ is bdd

* * *

(manoj)

Soln:

Here $u(0,t) = 0$

so, we take Fourier Sine transform of both sides of the given eqn.

$$\int_0^{\infty} \frac{\partial u}{\partial t} \sin px dx = \int_0^{\infty} \frac{\partial^2 u}{\partial x^2} \sin px dx$$

$$\Rightarrow \frac{d}{dt} \int_0^{\infty} u \sin px dx = \left[\frac{\partial u}{\partial x} \sin px \Big|_0^{\infty} - \int_0^{\infty} \frac{\partial u}{\partial x} p \cos px dx \right]$$

Here $\bar{u}_s = \int_0^{\infty} u \sin px dx$

$$\Rightarrow \frac{d\bar{u}_s}{dt} = \left[0 - \int_0^{\infty} p \frac{\partial u}{\partial x} \cos px dx \right]$$

$$= -p \left[u(x,t) \cos px + \int_0^{\infty} u(x,t) \sin px dx \right]$$

$$= -p \left[0 + p \int_0^{\infty} u(x,t) \sin px dx \right]$$

$$= -p^2 \bar{u}_s$$

Monday
29-2-16

$$1) \frac{\partial \bar{u}_g}{\partial t} = -P^2 \bar{u}_g$$

$$2) \frac{\partial \bar{u}_g}{\partial x} = -P^2 \bar{u}_g$$

$$\Rightarrow \bar{u}_g = c e^{-P^2 t}$$

By (ii) $\int_0^{\infty} u(x,0) \sin px dx = \int_0^a \sin px dx + \int_a^{\infty} 0 \sin px dx$

$$= \int_0^a \sin px dx$$

$$= \left[\frac{-\cos px}{p} \right]_0^a = \left[\frac{-\cos pa + 1}{p} \right]$$

$$\bar{u}(p,0) = \frac{1}{p} \cos pa$$

$$\Rightarrow \bar{u}(p,0) = c \cdot e^{-P^2 \cdot 0} = c$$

$$\Rightarrow c = \frac{1}{p} - \frac{\cos pa}{p}$$

$$\Rightarrow \bar{u}(p,t) = c \cdot e^{-P^2 t} = \left(\frac{1}{p} - \frac{\cos pa}{p} \right) e^{-P^2 t}$$

Now, inversion formula -

$$u(x,t) = \frac{2}{\pi} \int_0^{\infty} \left(\frac{1}{p} - \frac{\cos pa}{p} \right) e^{-P^2 t} \sin pxdp$$

Example 2 Solve the

$$\frac{\partial \theta}{\partial t} = c^2 \frac{\partial^2 \theta}{\partial x^2} \quad (t > 0) \quad \text{--- (1)}$$

If the initial temperature of an infinite bar is given by

$$\theta(x) = \begin{cases} \theta_0 & \text{for } |x| < a \\ 0 & \text{for } |x| > a \end{cases} \quad \text{--- (2)}$$

Prepared by Kalika

Determine temp. at any pt. x & at any instant time t, $\theta(x,t)$

Here $\bar{\theta} = \int \theta \sin px dx$

So, $\frac{d\bar{\theta}}{dt} = -c^2 p^2 \bar{\theta} \Rightarrow \bar{\theta} = d e^{-c^2 p^2 t}$

Here $\theta(x,0) = \begin{cases} \theta_0 & |x| < a \\ 0 & |x| > a \end{cases} \quad \text{--- (2)}$

So, no take fourier sine transform -

Now, $\int_0^{\infty} u(x,0) \sin px dx = \int_0^a \theta_0 \sin px dx + \int_a^{\infty} 0 \sin px dx$

$$= \int_0^a \theta_0 \sin px dx$$

$$= \theta_0 \left[\frac{-\cos px}{p} \right]_0^a = \theta_0 \left[\frac{-\cos pa + 1}{p} \right]$$

$$\bar{\theta}(p) = \frac{2\theta_0 \sin pa}{p} = \frac{2\theta_0}{p} \left[\frac{1 - \cos pa}{2} \right]$$

$$\bar{\theta}(p,0) = \frac{2\theta_0 \sin pa}{p} = \frac{2\theta_0}{p} \left[\frac{1 - \cos pa}{2} \right]$$

$$\Rightarrow \bar{\theta}(p,0) = d \cdot e^{-c^2 p^2 \cdot 0} = d$$

$$\Rightarrow d = \frac{2\theta_0}{p} \left[\frac{1 - \cos pa}{2} \right]$$

$$\bar{\theta}(p,t) = \frac{2\theta_0}{p} \left[\frac{1 - \cos pa}{2} \right] e^{-c^2 p^2 t}$$

By inversion formula -

$$\theta(x,t) = \frac{2}{\pi} \int_0^{\infty} \left[\frac{1 - \cos pa}{p} \right] e^{-c^2 p^2 t} \sin pxdp$$

$$\bar{Q} = \frac{2\theta_0 \cdot \sin ap}{p} \cdot e^{-c^2 p^2 t}$$

Now, taking its inverse Fourier transform
We get —

$$\theta = \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{2\theta_0 \sin ap}{p} \cdot e^{-c^2 p^2 t} \cdot e^{ipx} dp$$

$$= \frac{\theta_0}{\pi} \int_{-\infty}^{\infty} \frac{\sin ap}{p} \cdot e^{-c^2 p^2 t} (\cos px - i \sin px) dp$$

$$= \frac{\theta_0}{\pi} \int_{-\infty}^{\infty} \frac{(\sin ap \cos px - i \sin ap \cdot \sin px) \cdot e^{-c^2 p^2 t}}{p} dp$$

$$= \frac{\theta_0}{\pi} \int_{-\infty}^{\infty} \frac{e^{-c^2 p^2 t} \cdot \sin ap \cos px}{p} dp -$$

$$\frac{\theta_0}{\pi} \int_{-\infty}^{\infty} \frac{e^{-c^2 p^2 t} \sin ap \cdot \sin px}{p} dp$$

$$\theta = \frac{\theta_0}{2} \left[\operatorname{erf} \left(\frac{ax}{2\sqrt{t}} \right) + \operatorname{erf} \left(\frac{a-x}{2\sqrt{t}} \right) \right]$$

Example-3

In infinite string is initially at rest and that the initial displacement is $f(x)$ ($-\infty < x < \infty$). Determine the displacement $y(x,t)$ of the string.

Example-4 Solve the PDE

$$U_t(x,t) = U_{xx}(x,t) \Rightarrow \frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

then $u(0,t) = u(2^2,t) = 0$ and

$$u(x,0) = 2x \quad \text{where}$$

$$0 < x < 2^2, \quad t > 0.$$

80 14
self

Since $u(0,t) = 0$ is given

We take finite Fourier sine transform

$$\int_0^4 \frac{\partial u}{\partial t} \sin \left(\frac{n\pi}{4} \right) x dx = \int_0^4 \frac{\partial^2 u}{\partial x^2} \sin \left(\frac{n\pi}{4} \right) x dx$$

$$\frac{d \bar{u}_p}{dt} = -\beta^2 F_p \left(\frac{\partial^2 u}{\partial x^2} \right) = -\beta^2 \bar{u}_p$$

$$\Rightarrow \frac{d \bar{u}_p}{\bar{u}_p} = -\beta^2 dt = -\left(\frac{n\pi}{4} \right)^2 dt$$

$$\Rightarrow \boxed{\bar{u}_p = C e^{-\frac{n^2 \pi^2}{16} t}} \quad \text{--- (i)}$$

Putting $t=0$,

$$\bar{u}_p(x,0) = C$$

Now taking finite Fourier sine transform on second term

$$\begin{aligned} \bar{u}_p(x,0) &= \int_0^4 u(x,0) \sin \left(\frac{n\pi}{4} \right) x dx \\ &= \int_0^4 2x \sin \left(\frac{n\pi}{4} \right) x dx \end{aligned}$$

$$= \frac{-32}{n\pi} \cos n\pi$$

$$\Rightarrow C = \frac{-32}{n\pi} \cos n\pi$$

by (i) —

$$\bar{u}_p(x,t) = \frac{-32 \cos n\pi}{n\pi} e^{-\frac{n^2 \pi^2}{16} t}$$

taking finite inverse transform

$$u_p(x,t) = \int_0^4 \frac{-32 \cos n\pi}{n\pi} e^{-\frac{n^2 \pi^2}{16} t} \cdot \sin \left(\frac{n\pi}{4} \right) x dx$$

Review:

By pythagoras formula

$$\delta s = \sqrt{(\delta x)^2 + (\delta y)^2}$$

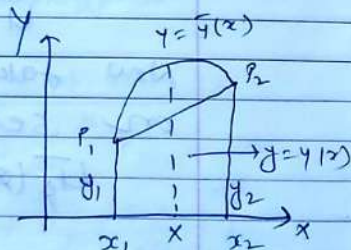
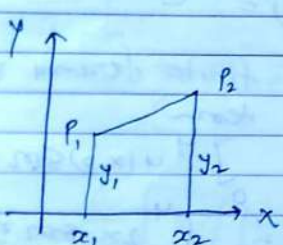
$$\delta f = \partial_x \sqrt{1 + \left(\frac{\partial y}{\partial x}\right)^2}$$

$$= \partial_x \sqrt{1 + |y'|^2}$$

$$\int_{x_1}^{x_2} \sqrt{1 + |y'|^2} dx \approx \int_{y_1}^{y_2} dy$$

* -

(29) EULER'S EQUATION or Euler's - Lagrange Eqⁿ: -



(30) 1-10/3

Let us examine the

$$I[y(x)] = \int_{x_1}^{x_2} F(x, y, y') dx \quad (1)$$

For an example - extreme value, the boundary point of the curve $y(x) = y_1$ & $y(x) = y_2$ — (2)

We assume that the given function - $F(x, y, y')$ (Max or min).

Let $y = y(x)$ be the curve which (Max or min). the eqⁿ (1) then assume that $y = \bar{y}(x)$ is an curve close to $y = y(x)$ st both $y(x)$ & $\bar{y}(x)$ are included in a one-parameter of curve

$$y(x, \beta) = y(x) + \beta \{ \bar{y}(x) - y(x) \} \quad (3)$$

For $\beta = 0$, then -

$$y = y(x)$$

For $\beta = 1$

$$y = \bar{y}(x)$$

$$\text{Now, } \bar{y}(x) - y(x) = \delta y \quad (4)$$

Where, δy is known as the variation of the function $y(x)$ then eqⁿ (3) reduces to

$$y(x, \beta) = y(x) + \beta \delta y \quad (5)$$

The variation δy is a function of x , In eqⁿ (4), we have -

$$\delta y' = \bar{y}'(x) - y'(x) = (\delta y)'$$

$$\text{In eqⁿ (5), } \left. \begin{aligned} \frac{\partial y}{\partial \beta} &= \delta y \end{aligned} \right\} (6)$$

$$\text{Similarly } \frac{\partial y'}{\partial \beta} = \delta y'$$

Now eqⁿ (1) & (5) reduces of β , say

$$\phi(\beta) = \int_{x_1}^{x_2} F[x, y(x, \beta), y'(x, \beta)] dx \quad (7)$$

For the extremum of the function,

$$\phi(\beta) \text{ for } \beta = 0$$

$$\text{then } \left(\frac{\partial \phi}{\partial \beta} \right)_{\beta=0} = 0 = \text{(next page)}$$

$$\left(\frac{\partial \Phi}{\partial \beta}\right)_{\beta=0} = \int_{x_1}^{x_2} \left[\frac{\partial F}{\partial x} \frac{\partial x}{\partial \beta} + \frac{\partial F}{\partial y} \frac{\partial y}{\partial \beta} + \frac{\partial F}{\partial y'} \frac{\partial y'}{\partial \beta} \right] dx$$

$$\Rightarrow \int_{x_1}^{x_2} \left[0 + \frac{\partial F}{\partial y} \delta y + \frac{\partial F}{\partial y'} \delta y' \right] dx = 0$$

$$\Rightarrow \int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \right) \delta y dx + \int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y'} \right) \delta y' dx = 0$$

$$\Rightarrow \int_{x_1}^{x_2} \left(\frac{\partial F}{\partial y} \right) \delta y dx + \left[\left(\frac{\partial F}{\partial y'} \right) \delta y \right]_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \delta y dx = 0$$

$$\delta y dx = 0$$

Note:-

$\because y = y(x)$ & $y = \bar{y}(x)$ both pass through $P_1(x_1, y_1)$ and $P_2(x_2, y_2)$

$\Rightarrow y(x_1) = \bar{y}(x_1)$ and $y(x_2) = \bar{y}(x_2)$ using eq (4)

$$\delta y \Big|_{x=x_1} = \bar{y}(x_1) - y(x_1) = 0$$

$$\delta y \Big|_{x=x_2} = \bar{y}(x_2) - y(x_2) = 0$$

$$\therefore \left[\frac{\partial F}{\partial y'} \delta y \right]_{x_1}^{x_2} = 0$$

$$\therefore \int_{x_1}^{x_2} \left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] \delta y dx = 0 \quad \text{--- (5)}$$

Now the first term (factor)

$$\left[\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right] = 0 \quad \text{on the extremizing curve}$$

$y = y(x)$ is a given c.p. fⁿ.
while the second factor $\delta y = 0$ because

of the arbitrary choice of the curve $y = \bar{y}(x)$ at $x = x_1$ & $x = x_2$

Hence, the function extremizing f^n $y = y(x)$ satisfy the differential eqⁿ

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0 \quad x_1 \leq x \leq x_2$$

which is known as Euler's Equation for functional for eqⁿ (1) and also

known as Euler-Lagrange's eqⁿ.

* * *

Tuesday

8/3/16

No class / shivnath
7/3/16

Example- Find the extremal for the functional

(P-147) Similar (i) $\int_0^{\pi/2} [y'^2 - y^2 + 2xy] dx$ with $y(0) = 0$, &

(148) (ii) $\int_0^{\pi/2} [y'^2 - y^2 - 2y \sin x] dx$ with $y(0) = 0$ & $y(\pi/2) = 1$

(iii) $\int_0^1 (xy + y^2 - 2y^2 y') dx$ with $y(0) = 1$ & $y(1) = 2$

(Not) $\leftarrow y(1) = 2$

Solⁿ
①

We know that

$$I[y(x)] = \int_{x_1}^{x_2} F[x, y, y'] dx$$

$$F[x, y, y'] = y'^2 - y^2 + 2xy$$

Using Euler's Equation -

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$$

$$2x - 2y - \frac{d}{dx} (2y') = 0 \Rightarrow y'' = -y + x$$

$$\text{or } (D^2 + 1)y = 0$$

2. y'' = x

$$\Rightarrow \frac{\partial}{\partial x} (2y' - 2y + 2x) - \frac{\partial}{\partial x} (2y') = 0$$

$$-2y + 2x - 2y'' = 0$$

$$\Rightarrow y'' + y = x$$

(D^2+1)y = x

Case-I, x=0 $\Rightarrow D = \pm i$

CF = C1 cos x + C2 sin x

Case-II, x ≠ 0

$$PI = \frac{1}{(D^2+1)} x = \frac{1}{(1+D^2)} x = x$$

$$PI = \frac{1}{(1+D^2+D^2+\dots)} x = x$$

$$y = x - \frac{\pi}{2} \sin x$$

y = CF + PI

$$y = C_1 \cos x + C_2 \sin x + x$$

y(0) = C1 cos 0 + C2 sin 0 + 0 = 0

$$= C_1 + 0 = 0 \Rightarrow C_1 = 0$$

y(pi/2) = C1 * 0 + C2 * 1 + pi/2 = 0

$$\Rightarrow C_2 = -\pi/2$$

$$y = x - \frac{\pi}{2} \sin x$$

* Some important cases in Euler's-Lagrange:-

Case-I, If F is independent of y' so that $\frac{\partial F}{\partial y'} = 0$, Hence the Euler's eqn $\frac{\partial F}{\partial y} - \frac{\partial}{\partial x} (\frac{\partial F}{\partial y'}) = 0$ reduces to $\frac{\partial F}{\partial y} = 0$ which is a first order ODE.

Case-II, If the f^n F is L.D on y' s.t $F(x, y, y') = M(x, y) + N(x, y)y'$

then $\frac{\partial F}{\partial y} = \frac{\partial M}{\partial y} + \frac{\partial N}{\partial y} y'$

$\frac{\partial F}{\partial y'} = N(x, y)$, $\frac{\partial}{\partial x} (\frac{\partial F}{\partial y'}) = \frac{d}{dx} N(x, y)$

Hence the Euler eqn $\frac{\partial F}{\partial y} - \frac{\partial}{\partial x} (\frac{\partial F}{\partial y'}) = 0$

$\frac{\partial M}{\partial y} + \frac{\partial N}{\partial y} y' - \frac{d}{dx} N(x, y) = 0$

$\frac{\partial M}{\partial y} + \frac{\partial N}{\partial y} y' - (\frac{\partial N}{\partial x} + \frac{\partial N}{\partial y} \frac{dy}{dx}) = 0$

$\frac{\partial M}{\partial y} + \frac{\partial N}{\partial y} y' - \frac{\partial N}{\partial x} - \frac{\partial N}{\partial y} y' = 0$

$\frac{\partial M}{\partial y} - \frac{\partial N}{\partial x} = 0$

Hence the expression Mdx + Ndy is an exact DE, then the functional

$$I[y(x)] = \int_{x_1}^{x_2} F(x, y, y') dx = \int_{x_1}^{x_2} (Mdx + Ndy)$$

Case-III, The functional F is independent of x and y i.e. F is dependent only on y'.
Special form

$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$ (1)

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Second form

So that $F = F(y)$

$\frac{dF}{dx} = \frac{\partial F}{\partial x} + \frac{\partial F}{\partial y} \frac{dy}{dx} + \frac{\partial F}{\partial y'} \frac{dy'}{dx}$ (2)

$\frac{dF}{dx} = \frac{\partial F}{\partial x} + \frac{\partial F}{\partial y} y' + \frac{\partial F}{\partial y'} y''$ (2)

$\frac{d}{dx} \left(y' \frac{\partial F}{\partial y'} \right) = y' \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) + \frac{\partial F}{\partial y'} y''$ (3)

Now by eqn (2) + (3)

$\frac{dF}{dx} - \frac{d}{dx} \left(y' \frac{\partial F}{\partial y'} \right) = \frac{\partial F}{\partial x} + y' \left\{ \frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) \right\}$

$\frac{d}{dx} \left(F - y' \frac{\partial F}{\partial y'} \right) = \frac{\partial F}{\partial x} = 0$

fund for

Since, $\frac{\partial F}{\partial y}$ is function of x, y & y'' , we have

$\frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = \frac{\partial}{\partial x} \left(\frac{\partial F}{\partial y'} \right) + \frac{\partial}{\partial y} \left(\frac{\partial F}{\partial y'} \right) \frac{dy}{dx} + \frac{\partial}{\partial y''} \left(\frac{\partial F}{\partial y'} \right) \frac{dy''}{dx}$

$\frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = \frac{\partial^2 F}{\partial x \partial y'} + y' \frac{\partial^2 F}{\partial y \partial y'} + y'' \frac{\partial^2 F}{\partial y''^2}$ (4)

Using eqn (4) from eqn (2), reduces to

$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$

$\frac{\partial F}{\partial y} - \frac{\partial^2 F}{\partial x \partial y'} - y' \frac{\partial^2 F}{\partial y \partial y'} - y'' \frac{\partial^2 F}{\partial y''^2} = 0$ (5)

The above expression is third form of Euler's Eqn.

In the case of III, $\frac{\partial F}{\partial y} = 0, \frac{\partial F}{\partial x} = 0$

$\frac{\partial^2 F}{\partial x \partial y'} = 0, \frac{\partial^2 F}{\partial y \partial y'} = 0$

then eqn (5) reduces

$-y'' \frac{\partial^2 F}{\partial y''^2} = 0 \Rightarrow y'' = 0$

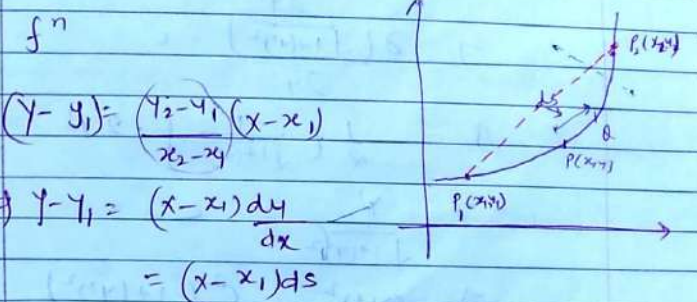
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$y = c_1 + c_2 x$

Example: (P-9)

Prove that, the shortest distance between the two points in a plane is st. line

Soln:



Soln:

Let $P_1(x_1, y_1)$ & $P_2(x_2, y_2)$ be two pts given on a plane.

Let $P(x, y)$ be any point on curve joining P_1, P_2 . let the arc $P_1 P_2 = s$, then arc length

$P_1 P_2 = \int_{x_1}^{x_2} ds$

$= \int_{x_1}^{x_2} \sqrt{1 + M'^2} dx$

$P_1 P_2 = \int_{x_1}^{x_2} f(y) dx, f(y) = \sqrt{1 + y'^2}$

The condition is that the curve b/w P_1 to P_2 be the shortest path.

Euler's formula

$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$

($\because f$ is independent of $y \rightarrow \frac{\partial F}{\partial y} = 0$)

$\Rightarrow -\frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$

self

$$\Rightarrow \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial y'} \right) = 0$$

on integrating both side, we have

$$\frac{\partial f}{\partial y'} = C$$

$$\Rightarrow \frac{\partial (\sqrt{1+(y')^2})}{\partial y'} = C$$

$$\Rightarrow \frac{1}{2} \left(\frac{2y'}{\sqrt{1+(y')^2}} \right) \cdot 2y' = C$$

$$\Rightarrow \frac{y'}{\sqrt{1+(y')^2}} = C$$

$$\Rightarrow (y')^2 = C^2 (1+(y')^2)$$

more info... $(y')^2 = C^2 (1+(y')^2)$... $\frac{\partial f}{\partial y'} = C$ (a is constant)

$$\Rightarrow y' = \frac{dy}{dx} = a$$

$$\Rightarrow dy = a dx$$

$$\Rightarrow y = ax + b, \text{ b is constant}$$

which is a line.

self

2) Show that the area of the surface of revolution of a curve $y = y(x)$ is

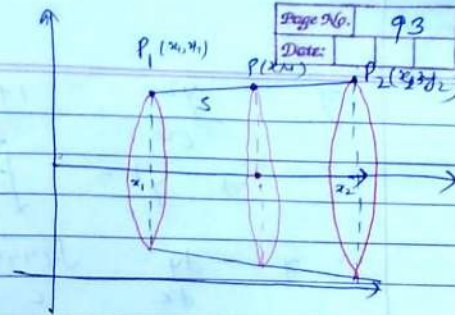
$$2\pi \int_a^b y \sqrt{1+(y')^2} dx$$

Hence, ST for f to be minimum, the curve must be catenary.

Let $P_1(x_1, y_1)$ & $P_2(x_2, y_2)$ be the extremities of any curve, now rotate the curve about x-axis.

self

Let $P(x, y)$ be any point on the curve. \therefore then the surface area of revolution is given by—



$$S = \int 2\pi y ds$$

$$= \int 2\pi y \frac{ds}{dx} dx$$

$$= 2\pi \int_{x_1}^{x_2} y \sqrt{1+(y')^2} dx$$

$$\because \frac{ds}{dx} = \sqrt{1+\left(\frac{dy}{dx}\right)^2} = \sqrt{1+(y')^2}$$

Proof

part-II, This has to be minimum

Since $f = y \sqrt{1+(y')^2}$ is independent of x . \therefore By Euler's Equation—

$$\frac{\partial f}{\partial y} - \frac{\partial}{\partial x} \left(\frac{\partial f}{\partial y'} \right) = 0$$

on integrating both side, we have

$$f - y' \frac{\partial f}{\partial y'} = C$$

$$\Rightarrow y \sqrt{1+(y')^2} - y' \frac{\partial}{\partial y'} (y \sqrt{1+(y')^2}) = C$$

$$\Rightarrow y \sqrt{1+(y')^2} - y' \left\{ y \cdot \frac{2y'}{2\sqrt{1+(y')^2}} \right\} = C$$

$$\Rightarrow y \sqrt{1+(y')^2} - \frac{y y'^2}{\sqrt{1+(y')^2}} = C$$

$$\Rightarrow y \left\{ \frac{1+(y')^2 - y'^2}{\sqrt{1+(y')^2}} \right\} = C$$

$$\Rightarrow \frac{y}{\sqrt{1+(y')^2}} = C$$

$$\frac{\partial f}{\partial y'} = y \frac{\partial}{\partial y'} \left(\frac{y \sqrt{1+(y')^2}}{y} \right)$$

$$\frac{y^2}{c^2} = 1 + y'^2$$

$$\Rightarrow y' = \sqrt{1 + \frac{y^2}{c^2}} = \frac{\sqrt{-c^2 + y^2}}{c}$$

$$\Rightarrow \frac{dy}{dx} = \frac{\sqrt{-c^2 + y^2}}{c} \Rightarrow \int \frac{dy}{\sqrt{-c^2 + y^2}} = \int \frac{dx}{c}$$

Now, integrating both sides, we have

$$\Rightarrow \int \frac{dy}{\sqrt{-c^2 + y^2}} = \int \frac{dx}{c}$$

$$\Rightarrow \operatorname{Cosh}^{-1}\left(\frac{y}{c}\right) = \frac{x+a}{c} = \frac{x+ac}{c}$$

$$y = c \operatorname{Cosh}\left(\frac{x+ac}{c}\right)$$

which is the eqⁿ of the curve of which
is catenary

* * *

Holi Vacation
19-3-16 to 27-3-16

28/2/16

No class due to Absent of Sir

29/3/16

paper discussion

$$= \int_{\theta_1}^{\theta_2} (r^2 + r'^2)^{1/2} d\theta, \quad r' = \frac{dr}{d\theta}$$

$$F(\theta, r, r') = (r^2 + r'^2)^{1/2}$$

$$\frac{\partial F}{\partial r} - \frac{d}{d\theta} \left(\frac{\partial F}{\partial r'} \right) = 0 \quad \text{--- (1)}$$

$$r^3 + 2rr'^2 - r^2 r'' = 0$$

$$r r'' - 2r'^2 - r^2 = 0 \quad \text{--- (2)}$$

putting $p = r' = \frac{dr}{d\theta}$

$$\frac{dp}{dr} = \frac{2p}{r} + \frac{r}{p}$$

putting $v = p/r$

$$\Rightarrow p = r v$$

$$\left(\frac{v}{1+v^2} \right) dv = \frac{dr}{r}$$

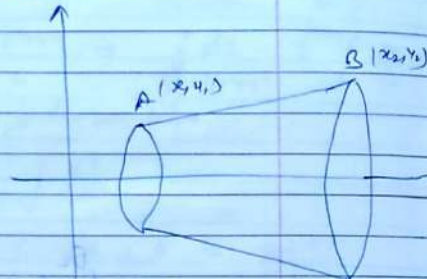
$$\Rightarrow v = \frac{\sqrt{r^2 - a^2}}{a}$$

$$p = \frac{r \sqrt{r^2 - a^2}}{a}$$

Prepared by Kalika

Exercise Find the curve passes through the pt. $A(x_1, y_1)$ & $B(x_2, y_2)$ which rotated about the x -axis gives a minimum surface area.

80/12
Date: 2/1/12



A curve joining $A(x_1, y_1)$ & $B(x_2, y_2)$ is revolved about x -axis. The area of surface of revolution

$$S[y(x)] = \int_{x_1}^{x_2} 2\pi y \, ds$$

$$= 2\pi \int_{x_1}^{x_2} y \sqrt{1+y'^2} \, dx$$

where $F(x, y, y') = y \sqrt{1+y'^2}$

Using Euler's - Lagrange Eqⁿ where x absent case (v).

$$\frac{d}{dx} \left(F - y' \frac{\partial F}{\partial y'} \right) = \frac{\partial F}{\partial x} = 0$$

$$\Rightarrow F - y' \frac{\partial F}{\partial y'} = C_1$$

$$\Rightarrow y = C_1 \sqrt{1+y'^2}$$

$$\Rightarrow y \sqrt{1+y'^2} = y' \left(\frac{y \sqrt{1+y'^2}}{2} \cdot 2y' \right) = C_1$$

$$\Rightarrow y(1+y'^2) - y'^2 y = C_1 \sqrt{1+y'^2}$$

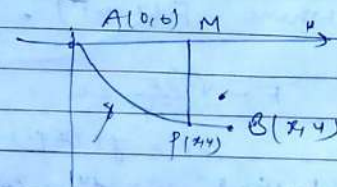
$$\Rightarrow y(1+y'^2 - y'^2) = C_1 \sqrt{1+y'^2}$$

$$\Rightarrow y = C_1 \sqrt{1+y'^2} \Rightarrow y' = \sqrt{\frac{y^2}{C_1^2} - 1}$$

$$\Rightarrow C_1 = \frac{y}{\sqrt{1+y'^2}}$$

(2) Find the curve connecting the pt of A & B. which is traversed by a particle sliding from A to B. in the shortest time (Friction & resistance of the medium are ignored) \rightarrow (Brachistochrone)

P-152



Let $A(0,0)$ be the origin, we minimize the time of descent from A to B at any time t using the principle of conservation of energy.

We have workdone in moving particle from A to B is equal to KE of P - PE of P. ($W = KE - PE$)

$$\Rightarrow mgy = \frac{1}{2}mv^2 - 0$$

$$v = \sqrt{2gy} \quad \left(v = \frac{ds}{dt} \right)$$

$$\Rightarrow dt = \frac{1}{\sqrt{2gy}} ds$$

the time spent in moving from A to B

$$t = \int_A^B \frac{ds}{\sqrt{2gy}}$$

$$t[y(x)] = \int_a^x \frac{\sqrt{1+y'^2}}{\sqrt{2gy}} dx$$

where $F(x, y, y') = \frac{\sqrt{1+y'^2}}{\sqrt{2gy}}$

$F(x, y, y')$ is independent of x , then Euler Eqn

$$\frac{d}{dx} \left(F - y' \frac{\partial F}{\partial y'} \right) - \frac{\partial F}{\partial x} = 0$$

2/3/16 ~~t[y(x)]~~ $t[y(x)] = \int_a^x \frac{\sqrt{1+y'^2}}{\sqrt{2gy}} dx$

where $F(x, y, y') = \frac{\sqrt{1+y'^2}}{y}$

$$F - y' \frac{\partial F}{\partial y'} = c$$

$$c_1 = \frac{1}{2} (1+y'^2)$$

putting $y' = \frac{dy}{dx} = \cot \theta$

$$y = \frac{c_1}{2} (1 - \cos 2\theta) \quad (*)$$

$$x = \frac{c_1}{2} (\theta - \sin 2\theta) + c_2$$

$$\Rightarrow (x - c_2) = \frac{c_1}{2} (\theta - \sin 2\theta)$$

putting $\psi = 2\theta, \theta = \frac{\psi}{2}$

Now by (*) and (**), we have $y = \frac{a}{2} (\cos \psi - 1) + c_2 \sin \psi$

Rewrite $x - c_2 = a(\psi - \sin \psi)$

$$y = a(1 - \cos \psi)$$

Now putting $A(0,0)$, where $x=0, y=0$ then $c_2 = 0, \psi = 0$

Hence, we get the eqn of cycloids in the st. form —

$$\left. \begin{aligned} x &= a(\psi - \sin \psi) \\ y &= a(1 - \cos \psi) \end{aligned} \right\} \text{ putting } y' = \frac{dy}{dx} = \cot \theta$$

where a is the radius of a rolling circle.

Isoperimetric Problems

(186)

In certain problem: It is necessary to make a given integral

$$I = \int_{x_1}^{x_2} f(x, y, y') dx \quad \text{--- (1)}$$

max. or min. while keeping other integral

$$J = \int_{x_1}^{x_2} g(x, y, y') dx \quad \text{--- (2)}$$

In such problems involves one or more constraints condition. Just as $J = a$, constant perimeter and surface area, we obtain refer to the problems of this type as Isoperimetric Problems.

Such problems are generally solved by the method of Lagrange Multiplier.

To extremize Eqn (1), multiply Eqn (2) by λ and adding to Eqn (1), where λ is the Lagrange multiplier.

Hence the necessary condition for this integral

$\int_{x_1}^{x_2} H dx$ to be an extremum is

$$\frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) = 0 \text{ where } H = f + \lambda g.$$

The value of the two constants of integration & the parameter λ .

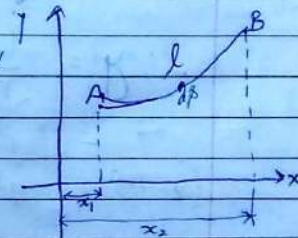
Example: Find the plane curve of fixed perimeter and Maximum Area.

(Self)

Solⁿ let l be the fixed perimeter of a plane curve b/w the points A, B , with x_1 & x_2

$$l = \int_{x_1}^{x_2} \sqrt{1+y'^2} dx \quad \text{--- (1)}$$

A = the area between the curve and x -axis.



$A = \int_a^x y dx$ — (2)
 Here, we have to maximize area, so max. @ & fo (without) (1)

Now, $H = f + \lambda g = y + \lambda \sqrt{1+y^2}$

Using Euler-Lagrange,

$$1 - \frac{d}{dx} \left(\frac{\lambda y'}{\sqrt{1+y^2}} \right) = 0 \quad \& \quad \frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) = 0$$

$$x - \frac{\lambda y'}{\sqrt{1+y^2}} = a \Rightarrow (x-a) = \frac{\lambda y'}{\sqrt{1+y^2}}$$

$$\left(\frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) \right) = 1 - \frac{d}{dx} \left(\frac{\lambda y'}{\sqrt{1+y^2}} \right) = 0$$

$$\Rightarrow \frac{d}{dx} \left(\frac{\lambda y'}{\sqrt{1+y^2}} \right) = 1$$

$$\Rightarrow \frac{\lambda y'}{\sqrt{1+y^2}} = x+a$$

$$\Rightarrow \frac{\lambda y'}{\sqrt{1+y^2}} = \frac{(x-a)^2}{(1+y^2)}$$

$$\Rightarrow (x-a)^2 = \frac{(\lambda y')^2}{(1+y^2)} \Rightarrow \frac{1+y^2}{y^2} = \frac{\lambda^2}{(x-a)^2}$$

$$\Rightarrow (y')^2 (x-a)^2 = - (x-a)^2$$

$$\Rightarrow y' = \frac{(x-a)}{\sqrt{\lambda^2 - (x-a)^2}}$$

$$y' = \frac{(x-a)}{\sqrt{\lambda^2 - (x-a)^2}} \quad \left(\begin{array}{l} \lambda^2 - (x-a)^2 = t \\ -2(x-a) dx = dt \end{array} \right)$$

$$7 = \frac{\sqrt{\lambda^2 - (x-a)^2}}{\lambda^2 - (x-a)^2} = \frac{(x-a) \cdot dt}{\sqrt{t} \cdot -2(x-a)}$$

$$(y-b)^2 = \lambda^2 - (x-a)^2$$

$$\Rightarrow (y-b)^2 + (x-a)^2 = \lambda^2$$

which is circle

$$y = -\sqrt{a^2 - (x-a)^2} + b$$

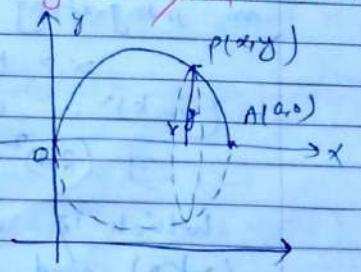
$$\Rightarrow (y-b)^2 = a^2 - (x-a)^2$$

$$\Rightarrow a^2 = (x-a)^2 + (y-b)^2$$

which represents a circle.

Example! Prove that the sphere is the solid of revolution which for a given surface area has max. volume.

Sol! Consider the arc DPA of the curve which rotate about the x-axis, then the surface area



$$S = \int_0^a 2\pi y \, ds$$

$$S = 2\pi \int_0^a y \, ds = 2\pi \int_0^a y \sqrt{1+y'^2} \, dx$$

And the volume of the surface solid, so formed - a

$$V = \int_0^a \pi y^2 \, dx \quad \text{where } f, g \text{ indep. of } x$$

then $H = \pi y^2 + \lambda \times 2\pi y \sqrt{1+y^2}$

Using Euler's Equation (Case: when the given H is independent of x).

(we use $\frac{d}{dx} \left(F - y' \frac{\partial F}{\partial y'} \right) - \frac{\partial F}{\partial x} = 0 - 1 - 0$)

$$\pi y^2 + 2\pi y \lambda \sqrt{1+y^2} - y' \left(2\pi y \lambda \cdot \frac{y'}{\sqrt{1+y^2}} \right) = C$$

$$\Rightarrow \pi y^2 + 2\pi y \lambda \sqrt{1+y^2} - \frac{2\pi y \lambda y'^2}{\sqrt{1+y^2}} = C$$

$$\Rightarrow \pi y^2 + \frac{2\pi y \lambda}{\sqrt{1+y^2}} = C$$

$$\Rightarrow \pi y^2 + \frac{2\pi y \lambda}{\sqrt{1+y^2}} = C$$

$$\pi y^2 + \frac{2\pi xy}{\sqrt{1+y^2}} = c \quad (1)$$

Since the curve passing through O and A for which $y=0$, then Eqn (1), $c=0$

$$\pi y^2 + \frac{2\pi xy}{\sqrt{1+y^2}} = 0 \Rightarrow y + \frac{2x}{\sqrt{1+y^2}} = 0$$

$$\frac{dy}{dx} = -\frac{2x}{y\sqrt{1+y^2}}$$

on integrating both sides, we have -

$$x = k - \sqrt{4x^2 - y^2}$$

$$(x-k)^2 = 4x^2 - y^2$$

$$\Rightarrow (2x)^2 = (x-k)^2 + y^2$$

which is a circle with centre $(k, 0)$ and radius $2x$.
(on putting $x=0, y=0$, we have $k=2x$)

Example Show that the functional

$$I[x(t), y(t)] = \int_0^{\pi/2} [2xy + (\frac{dx}{dt})^2 + (\frac{dy}{dt})^2] dt$$

s.t. $x(0)=0, x(\pi/2)=1, y(0)=0, y(\pi/2)=1$ is stationary for $x = -\sin t, y = \sin t$

pf: - Given that $F(x, y, y') = [2xy + (x')^2 + (y')^2]$

using Euler's Eqn,

$$\frac{\partial F}{\partial x} - \frac{d}{dt} \left(\frac{\partial F}{\partial x'} \right) = 0 \quad \text{and}$$

$$\frac{\partial F}{\partial y} - \frac{d}{dt} \left(\frac{\partial F}{\partial y'} \right) = 0$$

$$\Rightarrow 2y + 2x = 0 \Rightarrow y = -x$$

$$x'' = 2x \Rightarrow \frac{d^2 x}{dt^2} = 2x \quad (1)$$

$$\text{Similarly, } 2x - \frac{d^2 y}{dt^2} = 0 \Rightarrow \frac{d^2 y}{dt^2} = 2x \quad (2)$$

Q20. Solve the eqn (1) & (2),

then we have -

$$\frac{d^2}{dt^2} \left(\frac{d^2 x}{dt^2} \right) = 2x \Rightarrow \frac{d^4 x}{dt^4} - 2x = 0$$

$$\Rightarrow (D^4 - 2)x = 0$$

$$\Rightarrow (D^2 + 1)(D^2 - 1)x = 0 \Rightarrow D = \pm 1, \pm i$$

\Rightarrow G.F. is,

$$x = C_1 e^t + C_2 e^{-t} + C_3 \cos t + C_4 \sin t \quad (3)$$

$$y = x'' = C_1 e^t + C_2 e^{-t} - C_3 \cos t - C_4 \sin t \quad (4)$$

in eqn (4),

$$y(0) = 0 \Rightarrow C_1 + C_2 - C_3 = 0 \quad (5)$$

$$y(\pi/2) = 1 \Rightarrow C_1 e^{\pi/2} + C_2 e^{-\pi/2} - C_4 = 1 \quad (6)$$

in eqn (3),

$$x(0) = 0 \text{ then } C_1 + C_2 + C_3 = 0 \quad (7)$$

$$x(\pi/2) = -1 \Rightarrow C_1 e^{\pi/2} + C_2 e^{\pi/2} + C_4 = -1 \quad (8)$$

$$\text{then } x = -\sin t$$

$$y = \sin t$$

Example: Find the Extremal of the functional

$$I[x(t), y(t)] = \int_0^{\pi/4} [x^2 y' + 2x^2 + 2yy'] dt$$

where $x' = dx/dt, y' = dy/dt$ subject to the IC at $t=0$, then $x=y=0, t=\pi/4$

then $x=y=1$

$$\text{Ans } x=y = \frac{\sinh 2t}{\sinh \pi/2}$$

$$F = x^2 y' + 2x^2 + 2yy'$$

$$\frac{\partial F}{\partial x} - \frac{d}{dt} \left(\frac{\partial F}{\partial x'} \right) = 0$$

$$\Rightarrow 4x - \frac{d}{dt} y' = 0 \Rightarrow \frac{d^2 y}{dt^2} = 4x \quad (1)$$

$$\Rightarrow \frac{\partial F}{\partial y} - \frac{d}{dt} \left(\frac{\partial F}{\partial \dot{y}} \right) = 0$$

$$\Rightarrow 4y - \frac{d}{dt} (2x) = 0 \Rightarrow \frac{d^2 x}{dt^2} = 4y \Rightarrow y = \frac{1}{4} x''$$

$$\Rightarrow \frac{d^2}{dt^2} \left(\frac{1}{4} \frac{d^2 x}{dt^2} \right) = 4x \quad (\text{using (ii) in (i)})$$

$$\Rightarrow \frac{1}{4} \frac{d^4 x}{dt^4} = 4x \Rightarrow (D^4 - 16)x = 0$$

$$\Rightarrow (D^2 + 4)(D^2 - 4) = 0$$

$$\Rightarrow D = \pm 2i, \pm 2$$

$$x = c_1 e^{2t} + c_2 e^{-2t} + c_3 \cos 2t + c_4 \sin 2t$$

$$y = \frac{1}{4} x'' = \frac{1}{4} [4c_1 e^{2t} + 4c_2 e^{-2t} - 4c_3 \cos 2t - 4c_4 \sin 2t]$$

Now by, $x(0) = 0, c_1 + c_2 + c_3 = 0$
 $x(\pi/4) = 1, c_1 e^{\pi/2} + c_2 e^{-\pi/2} + c_4 = 1$

$y(0) = 0, c_1 + c_2 - c_3 = 0$
 $y(\pi/4) = 1, c_1 e^{\pi/2} + c_2 e^{-\pi/2} - c_4 = 1$

$\Rightarrow c_1 + c_2 = 0$ (a) $x(\pi/4) = 1 = c_1 e^{\pi/2} + c_2 e^{-\pi/2} + c_4$
 $c_1 e^{\pi/2} + c_2 e^{-\pi/2} = 2$ $x(0) = 0 = c_1 + c_2 + c_3$

$x = y = \frac{\sin 2t}{2}$

$y(0) = 0 = c_1 + c_2 - c_3 = 0$

(b) $y(\pi/4) = 1 = c_1 e^{\pi/2} + c_2 e^{-\pi/2} - c_4 = 0$

$a + b = 1 = 2c_4 \Rightarrow c_4 = \frac{1}{2}$

* Geodesic :-

A Geodesic on a surface is a curve along which the distance between two pt. on the surface is a minimum.

Example: Find the Geodesic of a Right Circular cylinder of radius a.

Sol:- In cylinder polar co-ordinate (ρ, ϕ, z) , $x = \rho \cos \phi, y = \rho \sin \phi, z = z$

The element of arc $d\rho$ on the surface of the circular cylinder of radius a is given then

$$(d\rho)^2 = (dz)^2 + (dy)^2 + (dx)^2$$

$$= (d\rho)^2 + \rho^2 (d\phi)^2 + (dz)^2$$

$$(d\rho)^2 = 0 + \rho^2 (d\phi)^2 + (dz)^2 \quad (\because \rho = a, d\rho = 0)$$

$$= (d\phi)^2 \{ \rho^2 + \left(\frac{dz}{d\phi} \right)^2 \}$$

$$d\rho = \sqrt{a^2 + \left(\frac{dz}{d\phi} \right)^2} d\phi$$

then minimum surface occupied

$$S = \int_{\phi_1}^{\phi_2} \sqrt{a^2 + (z')^2} d\phi \quad \text{--- (1)}$$

for the geodesics, given cylinder

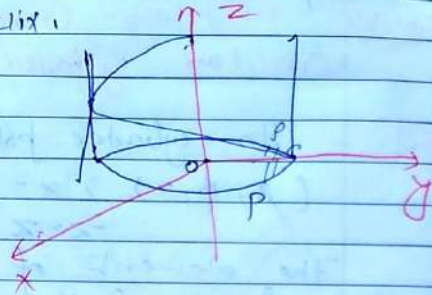
$$\frac{\partial F}{\partial z} - \frac{d}{d\phi} \left(\frac{\partial F}{\partial z'} \right) = 0$$

$$\Rightarrow -\frac{d}{d\phi} \frac{z z'}{\sqrt{a^2 + (z')^2}} = 0$$

$$\Rightarrow \frac{z'}{\sqrt{a^2 + (z')^2}} = c$$

$$\Rightarrow z = \left(\frac{a^2}{1-c^2} \right) \phi + c_2$$

The required Geodesics are Circular Helix. This curve is drawn on a circular cutting its generators at a constant angle and it is known as a Circular Helix.



Home work:

1) Prove that the Geodesics on a plane is straight line.

2) Find a function $y(x)$ for which $\int_0^1 (x^2 + y^2) dx$ is stationary.

Given that $\int_0^1 y^2 dx = 2$, $y(0) = 0$, $y(1) = 0$

3) Find the extremals of the isoperimetric problem

$x_0 = 0, x_1 = 2$
 $C = 1, y(0) = 0$
 $y(2) = 1$

$V[y(x)] = \int_{x_0}^{x_1} y^2 dx$, given that $\int_{x_0}^{x_1} y dx = C$, a constant.

4) Find the extremals of the functional $I = \int_0^\pi [y'^2 - y^2] dx$ under the condition $y(0) = 0, y(\pi) = 1$ and subject to the constraint $\int_0^\pi y dx = 1$

5) Determine the curve of length which passes through the points $(0,0)$ and $(1,0)$

and $(1,0)$ and for which the area between the curve and the x-axis is a maximum.

6) Find the curve of fixed length l that joins the point $P_1(x_1, 0)$ and $P_2(x_2, 0)$ lies above the x-axis and encloses the maximum area between itself and the x-axis.

7) Show that the curve c of given length l which minimize the curved surface area of solid generated by the revolution of c about x-axis is a catenary.

Tuesday
5/4/16

1) goto p-91
2) $I = \int (x^2 + y^2) dx$

Here $F(x,y) = x^2 + y^2$
Now using Euler's Eqn
 $\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$

Wrong Method

$0 = \frac{d}{dx} (2y') = 0$
 $2y' = C$
 $y' = \frac{C}{2} = d$
 $\Rightarrow y = dx + k$
 $\Rightarrow y^2 = (dx+k)^2$

II-Method

we have given that $I = \int_0^1 (x^2 + y^2) dx$
 $f = x^2 + y^2$
 $g = y^2 + \lambda(x - 1)$

By Euler's Eqn - $H = f + \lambda g = x^2 + y^2 + \lambda y^2$
 $\frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) = 0 \Rightarrow 2y + \frac{d}{dx} (2\lambda y) = 0$

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 Date: 9/4/16

$\Rightarrow \frac{d}{dx} (2\lambda y) = -2\lambda y$
 $\Rightarrow \lambda y' = -\lambda x + C$
 $\Rightarrow \lambda y = -\frac{\lambda}{2} x^2 + Cx + D$
 $\Rightarrow y'' = -\lambda$
 $(D^2 + 1)y = 0$
 $\Rightarrow D = \pm \sqrt{-1}$
 $D = \pm i\sqrt{1}$
 $D^2 - \lambda = 0$
 $D^2 = \lambda$
 $D = \pm \sqrt{\lambda}$

$y = C_1 e^{\sqrt{\lambda} t} + C_2 e^{-\sqrt{\lambda} t}$
 Now using given BC
 $y(0) = 0 = C_1 + C_2$
 $y(1) = 0 = C_1 e^{\sqrt{\lambda}} + C_2 e^{-\sqrt{\lambda}}$
 $C_1 (e^{\sqrt{\lambda}} - e^{-\sqrt{\lambda}}) = 0$
 $\Rightarrow C_1 = 0 = C_2$

3) given that d
 $v(x) = \int_0^x y^2 dx$ and $\int_0^x y dx = C$
 Here take $f = y^2 + g = y$
 Let $H = f + \lambda g = y^2 + \lambda y$
 Now using Euler's Eqn
 $\frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) = 0$
 $\Rightarrow 1 - \frac{d}{dx} (2y') = 0$

if $f = x^2 + y^2$
 then $H = x^2 + y^2 + \lambda y^2$
 By Euler's Eqn $\left(\frac{\partial H}{\partial y} - \frac{d}{dx} \left(\frac{\partial H}{\partial y'} \right) \right) = 0$
 $2\lambda y - \frac{d}{dx} (-2\lambda y) = 0$
 $\Rightarrow y'' + \lambda y = 0$
 $(D^2 + 1)y = 0 \Rightarrow D = \pm \sqrt{-1}$
 $\Rightarrow y = C_1 \cos \lambda t + C_2 \sin \lambda t$
 $y(0) = 0 = C_1 \Rightarrow C_1 = 0$
 $y(1) = 0 = C_2 \sin \lambda t$
 $\Rightarrow \sin \lambda t = 0$
 $\Rightarrow \sqrt{\lambda} = n\pi$
 $\Rightarrow y = C_2 \sin n\pi t$
 By condition
 $\int_0^1 C_2 \sin^2 n\pi t dt = 2$
 $\Rightarrow C_2 \int_0^1 \frac{1 - \cos 2n\pi t}{2} dt = 2$
 $\Rightarrow C_2 \left[\frac{t}{2} - \frac{\sin 2n\pi t}{2n\pi} \right]_0^1 = 4$
 $\Rightarrow C_2 \left[\frac{1}{2} - \frac{\sin 2n\pi}{2n\pi} \right] = 4$
 $C_2 = \frac{8}{2}$

$\Rightarrow \frac{d}{dx} (2y') = 1$
 $\Rightarrow 2y'' = 1 \Rightarrow 2D^2 y = 1$
 $D = \pm \sqrt{1/2}$
 $y = C_1 e^{\sqrt{1/2} x} + C_2 e^{-\sqrt{1/2} x}$
 Now using given condition
 $D^2 y = 1/2$
 on integrating both side twice
 $y = \frac{1}{9} x^2 + Cx + d$

Note: Suppose that intng
 given, it is given as
 $I = \int_0^2 y^2 dx + J = \int_0^2 y dx = 1 = C$
 BC, $y(0) = 0, y(2) = 1$

$y(0) = 0 = d$
 $y(2) = 1 = d + 2x \Rightarrow C = (1-d)/2$
 \therefore (1) becomes
 $y = \frac{1}{4} x^2 + \frac{(1-d)}{2} x$

4) given that $I = \int_0^\pi (y^2 - y) dx$
 and $J(y) = \int_0^\pi y dx = 1$
 Here we have
 $f = (y^2 - y) + g = y$

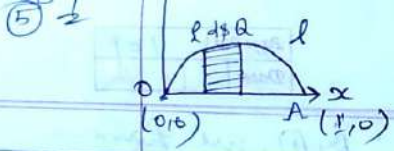
Let $F = f + \lambda g$
 $F = (y^2 - y) + \lambda y$
 Now using Euler's theorem
 $\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$
 $\Rightarrow (2y + \lambda) - \frac{d}{dx} (2y') = 0$
 $\frac{d}{dx} (2y') = (2y + \lambda)$

$y'' = \frac{(2y + \lambda)}{2}$
 On integrating twice,
 $y = \frac{(2y + \lambda)}{2} + Cx + d$
 Now using the condition
 $y(0) = 0 = d \Rightarrow d = 0$
 $y(\pi) = 1 = \frac{(2-d)\pi^2}{4} + C\pi$
 $\Rightarrow C = \frac{4 - (1-d)\pi^2}{4\pi}$

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By (1), we have
 $y = \frac{(1-d)x^2}{4} + \frac{(4-d)\pi^2 x}{4\pi}$
 Now substituting these values in (2), we have
 $\int_0^\pi \left(\frac{(1-d)x^3}{4} + \frac{(4-d)\pi^2 x^2}{4\pi} \right) dx - \int_0^\pi \left(\frac{(1-d)x^3}{4} + \frac{(4-d)\pi^2 x^2}{4\pi} \right) dx = 1$
 $= \left[\frac{(1-d)x^4}{16} + \frac{(4-d)\pi^2 x^3}{12} \right]_0^\pi - \left[\frac{(1-d)x^4}{16} + \frac{(4-d)\pi^2 x^3}{12} \right]_0^\pi = 1$
 $\Rightarrow \frac{(1-d)\pi^4}{16} + \frac{(4-d)\pi^3}{12} - \left(\frac{(1-d)\pi^4}{16} + \frac{(4-d)\pi^3}{12} \right) = 1$
 $\Rightarrow (1-d) = \frac{(2-\pi)\pi^2}{3}$
 $\Rightarrow d = 1 + \frac{(2-\pi)\pi^2}{3}$
 Putting in (2), we get desired result.

Prepared by



Here length of curve is given (ie perimeter) = l which passes through O and A . Let S be the area enclosed by the curve & x -axis. Then

$$S = \int_0^1 y dx$$

Subject to constraints

$$\int_0^1 ds = l$$

$$ie \int_0^1 \sqrt{1+y'^2} dx = l$$

BCs are $y(0) = y(1) = 0$

$$Let F = f + \lambda g$$

$$F = y + \lambda \sqrt{1+y'^2}$$

Now using Euler's Eqn

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$$

$$\Rightarrow 1 - \frac{d}{dx} \left[\frac{\lambda y'}{\sqrt{1+y'^2}} \right] = 0$$

$$\Rightarrow \frac{d}{dx} \left[\frac{\lambda y'}{\sqrt{1+y'^2}} \right] = 1$$

Now integrating both side

$$\frac{\lambda y'}{\sqrt{1+y'^2}} = x + C$$

$$\lambda y'^2 = (x+C)^2$$

$$\frac{1+y'^2}{y'^2} = \frac{\lambda^2}{(x+C)^2}$$

$$\Rightarrow \frac{1}{y'^2} + 1 = \frac{\lambda^2}{(x+C)^2}$$

$$\Rightarrow \frac{1}{y'^2} = \frac{\lambda^2 - (x+C)^2}{(x+C)^2}$$

$$\Rightarrow y' = \frac{x+C}{\sqrt{\lambda^2 - (x+C)^2}}$$

on integrating both side, we get

$$y = \int \frac{(x+C) dx}{\sqrt{\lambda^2 - (x+C)^2}}$$

$$Let \lambda^2 - (x+C)^2 = t^2$$

$$\Rightarrow -2(x+C) dx = 2t dt$$

$$\Rightarrow y = \int \frac{-t dt}{t} = -\int \frac{t}{t} dt = -t$$

$$\Rightarrow y = -\sqrt{\lambda^2 - (x+C)^2} + d$$

Now putting BC

$$y(0) = -\sqrt{\lambda^2 - C^2} + d = 0$$

$$y(1) = -\sqrt{\lambda^2 - (1+C)^2} + d = 0$$

$$\sqrt{\lambda^2 - (1+C)^2} = \sqrt{\lambda^2 - C^2}$$

$$\Rightarrow \lambda^2 - (1+C)^2 = \lambda^2 - C^2$$

$$C+1 = C$$

$$(y-d)^2 = (\lambda^2 - (x+C)^2)$$

$$(y-d)^2 + (x+C)^2 = \lambda^2$$

which is equation of circle with radius λ .

Method Now using BC

$$y(0) = 0 \Rightarrow (-d)^2 + C^2 = \lambda^2$$

$$\Rightarrow C^2 = \lambda^2 - d^2$$

$$y(1) = 0 \Rightarrow (1+C)^2 = \lambda^2 - d^2$$

② Find the functional for the B.V.P defined by $\frac{dy}{dx} = f(x)$, $y(a) = y(b) = 0$

$$dy = f(x) dx \Rightarrow y = \int f(x) dx + C$$

$$\int_a^b f(x) dy dx = \int_a^b \frac{d^2 y}{dx^2} dy dx$$

Now differentiating integrate by part



$$= \left[\frac{d^2 y}{dx^2} dy \right]_a^b - \int_a^b \frac{dy}{dx} \cdot \frac{d}{dx} \left(\frac{dy}{dx} \right) dx$$

$$\int_a^b f(x) dy dx = 0 - \int_a^b \frac{dy}{dx} \cdot f \left(\frac{dy}{dx} \right) dx$$

$$\left(\frac{d}{dx} dy \right) = d \left(\frac{dy}{dx} \right)$$

$$\left(\int y(a) = \int y(b) = y(a) = y(b) = 0 \right)$$

$$= -\frac{1}{2} \int_a^b \left(\frac{dy}{dx} \right)^2 dx$$

$$\Rightarrow \int_a^b f(x) dx + \frac{1}{2} \int_a^b \left(\frac{dy}{dx} \right)^2 dx = 0$$

$$\Rightarrow I(y) = \int_a^b \left[f(x) + \frac{1}{2} \left(\frac{dy}{dx} \right)^2 \right] dx$$

It follows that a unique solⁿ of the problem exist at a min. value of the integral (1). Therefore, the integral in eqn (2) represents the required f^oal of the problem.

In similar way, functional of other BVP & initial BVP can be derived.

(i) $\frac{d^2y}{dx^2} = f(x), y(a) = y(b) = 0$

$I(V) = \int_a^b V(2f - V'') dx$

(ii) $\frac{d^2y}{dx^2} + ky = x^2, 0 < x < 1, y(0) = 0,$

$(\frac{dy}{dx})_{x=1} = 1, I(V) = \frac{1}{2} \int_a^b [(\frac{dv}{dx})^2 - kv^2 + 2vx^2] dx - v(1)$

(iii) $x^2y'' + 2xy' = f(x), y(a) = y(b) = 0,$

$I(V) = \int_a^b V(2f - \frac{d}{dx}(x^2y')) dx$

(iv) $\nabla^2 u = 0, u = 0$ on the boundary C to R

$I(V) = \iint_R \frac{1}{2} [(\frac{\partial u}{\partial x})^2 + (\frac{\partial u}{\partial y})^2] dx dy$

(v) $\nabla^2 u = -f, u = 0$ on the boundary C to R

$I(V) = \iint_R \frac{1}{2} [(\frac{\partial u}{\partial x})^2 + (\frac{\partial u}{\partial y})^2] - uf$

(vi) $EI \frac{d^4y}{dx^4} + ky = f(x), 0 < x < l$

$y = 0 = \frac{d^2y}{dx^2}$ at $x = 0,$

$I(V) = \frac{1}{2} \int_0^l [EI (\frac{dv}{dx})^2 + kv^2 - 2vf] dx$

Base function :-

Suppose that to approximate a real valued f^n over a finite interval $[a, b]$. A usual approach is to define $[a, b]$ into a no. of

Subintervals $[x_i, x_{i+1}], i = 0, 1, 2, \dots, n-1$.
where $x_0 = a$ & $x_n = b$ & to interpolate linearity betⁿ the value of $f(x)$ at the end points of each sub-intervals.

In $[x_i, x_{i+1}]$, the linear approximation f^n is given by

$L_i(x) = \frac{1}{h_i} [(x_{i+1} - x)f_i + (x - x_i)f_{i+1}]$

where $h_i = x_{i+1} - x_i$, from this to construct the piecewise linear interpolating f^n over $[x_0, x_n]$.

by the formula $p(x) = \sum_{i=0}^n \phi_i(x) f_i$

where $\phi_0(x) = \begin{cases} \frac{(x_1 - x)}{h_0} & , x_0 \leq x \leq x_1 \\ 0 & , x_1 \leq x \leq x_2 \end{cases}$

$\phi_i(x) = \begin{cases} \frac{(x - x_{i-1})}{h_{i-1}} & , x_{i-1} \leq x \leq x_i \\ \frac{(x_{i+1} - x)}{h_i} & , x_i \leq x \leq x_{i+1} \\ 0 & , x > x_{i+1} \end{cases}$

$\phi_{n-1}(x) = \begin{cases} 0 & , x_0 \leq x \leq x_{n-1} \\ \frac{(x - x_{n-1})}{h_{n-1}} & , x_{n-1} \leq x \leq x_n \end{cases}$

The $f^n \phi_i(x), i = 1, 2, \dots, n$ are called Base function or Shape function. It is easily seen that the base $f^n \phi_i(x)$ are identically zero except for the range $[x_{i-1}, x_{i+1}]$ with $\phi_i(x_i) = 1$.

Rayleigh-Ritz Method

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Consider second order BVP defined by

$$y'' + p(x)y' + q(x)y = 0, \quad y(a) = y(b) = 0 \quad (1)$$

$$I[y] = \int_a^b \left[\left(\frac{dy}{dx} \right)^2 - p(x)y' - 2q(x)y \right] dx \quad (2)$$

If $y(x)$, the solⁿ eqⁿ (1) is substituted in eqⁿ (2), then the integral I will be min.

Now let $y(x) = \sum_{i=1}^n \alpha_i \phi_i(x)$ be an approximate solⁿ

$$I(\alpha_1, \alpha_2, \dots, \alpha_n) = \int_a^b \left[\left(\frac{d}{dx} \sum_{i=1}^n \alpha_i \phi_i \right)^2 - p \left(\sum_{i=1}^n \alpha_i \phi_i \right)' - 2q \sum_{i=1}^n \alpha_i \phi_i \right] dx \quad (3)$$

The base fun^s $\phi_i(x)$ are L.I. & satisfy the B.C. in eqⁿ (1)
 $\phi_i(a) = 0$ & $\phi_i(b) = 0$. (4)

For minimum, we have $\frac{\partial I}{\partial \alpha_1} \cdot d\alpha_1 + \frac{\partial I}{\partial \alpha_2} \cdot d\alpha_2$

$$+ \dots + \frac{\partial I}{\partial \alpha_n} \cdot d\alpha_n = 0$$

Since the $d\alpha_i$ are arbitrary eqⁿ (5)

$$\frac{\partial I}{\partial \alpha_i} = 0, \quad i = 1, 2, 3, \dots, n \quad (6)$$

If I is a quadratic of y & $\frac{dy}{dx}$, the eqⁿ (6) will be linear in α_i & can be solved easily.

* * *

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* Integral Equation: -

A Eqⁿ in which an unknown function appears under one, two or more integral signs is called integral equation.

If the derivative & the integral of the unknown f^n is in the Eqⁿ then this type of Eqⁿ is called Integro-Diff Eqⁿ.

Types of linear integral Eqⁿ: -

$$\alpha(x)u(x) + F(x) + \lambda \int_D k(x,t)u(t)dt = 0 \quad (1)$$

An integral Eqⁿ is called linear if only linear operator (linear operations) are performed upon the unknown function.

Thus an Eqⁿ of the form (1) is known as the linear integral equation.

where $u(x)$ is the unknown function, $\alpha(x)$, $F(x)$ and $k(x,t)$ are known functions. λ is a $n-2$ real or complex parameter.

The integration extends over the domain D of the variable t .

Linear integral Eq^s are classified into two types.

→ (I) Volterra Linear Integral Eqⁿ

The eqⁿ of the form

$$\alpha(x)u(x) = F(x) + \lambda \int_a^x k(x,t)u(t)dt$$

is a Volterra linear

integral Eqⁿ. The eqⁿ in which one of the limits is a variable, is called Volterra linear integral Eqⁿ.

Particular cases:

(i) When $\alpha(x) = 0$, Then the eqⁿ involves unknown f^n is only under the integral sign and such Eqⁿ is known as Volterra linear integral of the first kind.

$$F(x) = \int_a^x k(x+t)u(t)dt, \quad a > -\infty$$

(ii) When $\alpha(x) = 1$, Then the eqⁿ involves the unknown f^n is only under the integral sign & such Eqⁿ is known as Volterra linear integral of the second type

$$V(x) = F(x) + \int_a^x k(x+t)u(t)dt$$

(iii) When $\alpha(x) = 1, F(x) = 0$

$$U(x) = \int_a^x k(x+t)u(t)dt$$

is known as Volterra Homogeneous linear integral Eqⁿ.

Example: If $f(x) = \frac{1}{\pi\sqrt{x}}$ is a solⁿ of $\int_0^x \frac{f(s)}{\sqrt{x-s}} ds = 1$

Let $y(x) = f(x)$

$\Rightarrow y(s) = f(s)$

$$\int_0^x \frac{f(s)}{\sqrt{x-s}} ds = 1 \quad \frac{1}{\pi} \int_0^x \frac{ds}{\sqrt{x-s}}$$

Let $I = \int_0^x \frac{1}{\pi\sqrt{s}} \frac{1}{\sqrt{x-s}} ds$

$$= \frac{1}{\pi} \int_0^x \frac{1}{\sqrt{x-s} \sqrt{s}} ds = \frac{1}{\pi} \int_0^x \frac{ds}{\sqrt{\left(\frac{x}{2}\right)^2 - \left(s - \frac{x}{2}\right)^2}}$$

$$= \frac{1}{\pi} \left[\sin^{-1} \left(\frac{s - \frac{x}{2}}{\frac{x}{2}} \right) \right]_0^x$$

$$= \frac{1}{\pi} \left[\sin^{-1}(1) - \sin^{-1}(-1) \right]$$

$$= \frac{1}{\pi} \left[\frac{\pi}{2} - \left(-\frac{\pi}{2}\right) \right] = 1$$

Example: Solve the $y'' + 2y' + y = x^2$ with $y(0) = 1, y'(0) = 0$

Let $y'' = \phi(x)$ ← Integrating both

$$\int_0^x d\left(\frac{dy}{dx}\right) = \int_0^x \phi(x) dx$$

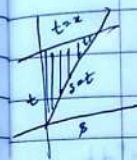
$$\left(\frac{dy}{dx}\right) - 0 = \int_0^x \phi(s) ds$$

$$y(x) - y(0) = \int_0^x \int_0^s \phi(s) ds dt$$

$$= \int_0^x \int_s^x \phi(s) ds dt$$

$$\Rightarrow y(x) = 1 + \int_0^x (x-s) \phi(s) ds$$

$$\phi(x) = 1 - x^2 - \int_0^x (2+x-s) \phi(s) ds$$



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Transformation of DE into Integral Eqn.

Example Consider the IVP consisting, 2nd order linear DE

$$\frac{d^2 y}{dx^2} + a_1 \frac{dy}{dx} + a_2 y(x) = F(x)$$

with $y(0) = y_0$ & $y'(0) = y_1$

Solⁿ define $\frac{dy}{dx} = \phi(x)$

Now eqn (1) can be written as

$$y'' + a_1 y' + a_2 y = F(x)$$

$$\int_0^x \frac{d^2 y}{dx^2} dt = \int_0^x F(t) dt - \int_0^x a_1 y' dt - \int_0^x a_2 y dt$$

$$\frac{dy}{dx} - y_1 = \int_0^x F(t) dt - \int_0^x a_1(t) \cdot y(t) dt$$

$$- \left\{ [a_1(t)y(t)]_0^x - \int_0^x \frac{d}{dt} [a_1(t)y(t)] dt \right\}$$

$$y(x) = y_0 + [a_1(0)y_0 + y_1]x + \int_0^x \left\{ \frac{d}{dt} [a_1(t) - a_2(t)] \right\}$$

$$(x-t) - a_1(t) \} y(t) dt + \int_0^x (x-t) F(t) dt$$

Now $y(x) = G(x) + \int_0^x k(x,t) y(t) dt$

where $G(x) = y_0 + [a_1(0)y_0 + y_1]x + \int_0^x (x-t) F(t) dt$

$$k(x,t) = \left\{ \frac{d}{dt} [a_1(t) - a_2(t)] \right\} (x-t) - a_1(t)$$

Unable to understand

$$\int_0^x y''(t) dt = \int_0^x u(t) dt$$

$$y'(x) - y'(0) = \int_0^x u(t) dt$$

$$y'(x) = y_1 + \int_0^x u(t) dt$$

Now again integrate both side

$$y(x) - y(0) = y_1 x + \int_0^x \left(\int_0^t u(t) dt \right) dt$$

$$y(x) = y_0 + y_1 x + \int_0^x (x-t) u(t) dt$$

Now putting these values in eqn (1)

$$y'' + a_1 y' + a_2 y = F(x)$$

$$u + a_1 \left(y_1 + \int_0^x u(t) dt \right) + a_2 \left[y_0 + y_1 x + \int_0^x (x-t) u(t) dt \right] = F(x)$$

$$U(x) = F(x) - y_0 a_2(x) - y_1 a_1(x) - \int_0^x (a_1 + a_2(x-t)) u(t) dt$$

Example:

Show that $f^n u(x) = x e^x$ is a solⁿ of the linear integral Eqn.

$$u(x) = \sin x + 2 \int_0^2 \log(x-t) u(t) dt$$

take R.H.S

given $u(x) = x e^x$ or $u(t) = t e^t$

$$\sin x + 2 \int_0^2 \log(x-t) \cdot t e^t dt$$

$$= \sin x + 2 \int_0^2 (\log x - \log t) t e^t dt$$

$$= \sin x + 2 \cos x \left[\int_0^x e^t \cos t dt \right] + 2 \sin x \left[\int_0^x e^t \sin t dt \right]$$

Formula

We know
 Let $I_1 = \int e^t \cos t dt$
 $= \frac{e^t}{2} (\sin t + \cos t)$
 $I_2 = \int e^t \sin t dt = \frac{e^t}{2} (\sin t - \cos t)$
 (simplify by self)

self

$$= \sin x + 2 \cos x \left[\frac{e^x}{2} (\sin x + \cos x) \right] - \int \frac{e^t}{2} (\sin t + \cos t) dt + 2 \sin x \left[\frac{e^x}{2} (\sin x - \cos x) \right] - \int \frac{e^t}{2} (\sin t - \cos t) dt$$

$$= \sin x + \cos x \left[x e^x (\sin x + \cos x) \right] + \sin x \left[x e^x (\sin x - \cos x) \right]$$

$$\cos x \left[- \int e^t \sin t dt - \int e^t \cos t dt \right] + \sin x \left[- \int e^t \sin t dt + \int e^t \cos t dt \right]$$

$$= \sin x + x e^x \left[\cos x \cdot \sin x + \cos^2 x + \sin^2 x - \sin x \cos x \right]$$

$$- \int e^t \sin t dt (\cos x + \sin x) - \int e^t \cos t dt (\cos x - \sin x)$$

$$= \sin x + x e^x - \left[\frac{e^x}{2} (\sin x - \cos x) + \frac{1}{2} x e^x (\sin x + \cos x) \right]$$

$$+ \left[\frac{e^x}{2} (\sin x + \cos x) + \frac{1}{2} x e^x (\cos x - \sin x) \right]$$

$$= \sin x + x e^x -$$

Example

Obtain Fredholm integration of second kind corresponding to the BVP

$$\frac{d^2 u}{dx^2} + u = x, \quad u(0) = 0, u(1) = 0$$

Sol

Let $\frac{d^2 u}{dx^2} = \phi(x)$
 $\int_0^x \frac{d^2 u}{dx^2} dx = \int_0^x \phi(x) dx$

$$\Rightarrow u'(x) - u'(0) = \int_0^x \phi(x) dx$$

$$\Rightarrow u'(x) = u'(0) + \int_0^x \phi(x) dx$$

(Given that $u' = x - u$ with $u(0) = 0, u(1) = 0$)
 Integrating w.r.t from 0 to 1

$$\int_0^x u'(t) dt = \int_0^x (x - u(t)) dt$$

$$u'(x) - u'(0) = \frac{x^2}{2} - \int_0^x u(t) dt \quad \text{--- (I)}$$

Again integrating both side

$$u(x) - u(0) = Ax + \frac{x^3}{6} - \int_0^x (x-t) u(t) dt \quad \text{--- (II)}$$

where $u'(0) = A$

Now putting $x=1$ in eqn (I)

$$A = -\frac{1}{2} + \int_0^1 u(t) dt \quad \text{--- (3)}$$

Now putting A in (II)

$$u(x) = \left[-\frac{1}{2} + \int_0^1 u(t) dt \right] x + \frac{x^3}{6} - \int_0^x (x-t) u(t) dt$$

$$= -\frac{x}{2} + \int_0^1 x u(t) dt + \frac{x^3}{6} - \int_0^x (x-t) u(t) dt$$

$$= -\frac{x}{2} + \frac{x^3}{6} + \int_0^x x u(t) dt + \int_x^1 x(u(t) dt) - \int_0^x (x-t) u(t) dt$$

Solution of Volterra's integral Eqⁿ of the IInd kind by the method of successive approximation

The Volterra's integral Eqⁿ of second kind

$$U(x) = F(x) + \lambda \int_0^x k(x,t) U(t) dt \quad \text{--- (1)}$$

Where $k(x,t)$ is a ctg fⁿ $0 \leq x \leq a$
 $0 \leq t \leq a$ and $F(x)$ is ctg for $0 \leq x \leq a$.

The solⁿ of the integral Eqⁿ (1) in the form of an infinite Power series is

$$U(x) = U_0(x) + \lambda U_1(x) + \lambda^2 U_2(x) + \lambda^3 U_3(x) + \dots + \lambda^n U_n(x) \quad \text{--- (2)}$$

Substituting this series in Eqⁿ (1) we have

$$U_0(x) + \lambda U_1(x) + \lambda^2 U_2(x) + \dots + \lambda^n U_n(x) = F(x) + \lambda \int_0^x k(x,t) \{ U_0(t) + \lambda U_1(t) + \lambda^2 U_2(t) + \dots + \lambda^n U_n(t) \} dt \quad \text{--- (3)}$$

Equating the coefficient of equal power of λ , we have

$$U_0(x) = F(x)$$

$$U_1(x) = \int_0^x k(x,t) U_0(t) dt$$

$$U_2(x) = \int_0^x k(x,t) F(t) dt$$

$$U_2(x) = \int_0^x k(x,t) U_1(t) dt = \int_0^x k(x,t) \left[\int_0^t k(t,t_1) F(t_1) dt_1 \right] dt \quad \text{--- (4)}$$

The relation Eqⁿ (4) yields a method for a successive determination of the fⁿ $U_n(x)$.

It may be shown that the series Eqⁿ (2) is convergent uniformly in x and λ for any λ and $x \in [0, a]$ under assumption with respect $F(x)$ & $k(x,t)$, its sum is unique solⁿ of Eqⁿ (1).

$$U_1(x) = \int_0^x k(x,t) F(t) dt$$

$$U_2(x) = \int_0^x k(x,t) \left[\int_0^t k(t,t_1) F(t_1) dt_1 \right] dt$$

By interchanging the order of integration we have

$$U_2(x) = \int_0^x F(t_1) dt_1 \int_{t_1}^x k(x,t) k(t,t_1) dt$$

$$U_2(x) = \int_0^x k_2(x,t) F(t) dt \quad \text{--- (5)}$$

where $k_2(x,t) = \int_t^x k(x,t) k(t,t) dt$

Similarly in general

$$U_n(x) = \int_0^x k_n(x,t) F(t) dt \quad \text{--- (6)}$$

The fⁿ $k_n(x,t)$ is called iterated Kernel.

By eqn (6), we have

$$U_1(x) = \int_0^x k_1(x,t) F(t) dt$$

$$U_2(x) = \int_0^x k_2(x,t) F(t) dt$$

Hence $k_1(x,t) = k(x,t)$

And for determination $k_2(x,t)$, $k_3(x,t)$ etc, we have Recursion formula

$$K_{n+1}(x,t) = \int_t^x k(x,z) k_n(z,t) dz \quad (7)$$

$n=1, 2, \dots$

Using eqn (6) & (7), the relation eqn (2) may be written as

$$U(x) = F(x) + \sum_{n=1}^{\infty} \lambda^n \int_0^x k_n(x,t) F(t) dt$$

The fn $R(x,t,d)$ defined as

$$R(x,t,d) = k_1(x,t) + \lambda k_2(x,t) + \lambda^2 k_3(x,t) + \dots + \lambda^n k_{n+1}(x,t) + \dots$$

$$= \sum_{m=0}^{\infty} \lambda^m k_{m+1}(x,t)$$

is called the Resolvent kernel.

(or Reciprocal kernel) for the integral.

eqn (1)

$$U(x) = F(x) + \lambda \int_0^x R(x,t,d) F(t) dt$$

Example:

Prove that

$$R(x,t,d) = k(x,t) + \lambda \int_t^x k(x,z) R(z,t,d) dz$$

Where $R(x,t,d)$ be the Resolvent kernel of a Volterra integral eqn

$$U(x) = F(x) + \lambda \int_0^x k(x,t) U(t) dt$$

Proof: We have $R(x,t,d) = \sum_{m=1}^{\infty} \lambda^m k_m(x,t)$ (1)

Where iterated kernel are given by

$$k_1(x,t) = k(x,t)$$

$$k_m(x,t) = \int_t^x k(x,z) k_{m-1}(z,t) dz$$

$$R(x,t,d) = k_1(x,t) + \sum_{m=2}^{\infty} \lambda^m k_m(x,t)$$

$$= k_1(x,t) + \sum_{m=2}^{\infty} \lambda^{m-1} \int_t^x k(x,z) k_{m-1}(z,t) dz$$

$$= k_1(x,t) + \sum_{n=1}^{\infty} \lambda^n \int_t^x k(x,z) k_n(z,t) dz$$

(putting $n=m-1$)

$$R(x,t,d) = k(x,t) + \lambda \sum_{n=1}^{\infty} \lambda^n \int_t^x k_n(z,t) k(x,z) dz$$

$$= k(x,t) + \lambda \int_t^x \left\{ \sum_{n=1}^{\infty} \lambda^n k_n(z,t) \right\} k(x,z) dz$$

then

$$R(x,t,d) = k(x,t) + \lambda \int_t^x k(x,z) R(z,t,d) dz$$

Example:

Find the Resolvent kernel of the Volterra integral eqn with $k(x,t) = 1$.

Solⁿ

Given that $k(x,t) = 1$

We know that $f_1(x,t) = k(x,t) = 1$

$$k_2(x,t) = \int_t^x k(x,z) k_1(z,t) dz$$

$$= \int_t^x 1 \cdot 1 dz = (x-t)$$

$$k_3(x,t) = \int_t^x k(x,z) k_2(z,t) dz$$

$$= \int_t^x 1 \cdot (z-t) dz = \frac{(x-t)^2}{2!}$$

Similarly for

$k_4, k_5, k_6 \dots$

$$k_m(x,t) = \int_t^x k(x,z) k_{m-1}(z,t) dz$$

$$= \int_t^x 1 \cdot \frac{(z-t)^{m-2}}{(m-2)!} dz$$

$$= \frac{(x-t)^{m-1}}{(m-1)!}$$

Hence the Resolvent Kernel

$$R(x,t,t) = \sum_{m=0}^{\infty} \lambda^m k_{m+1}(x,t)$$

$$= \sum_{m=0}^{\infty} \lambda^m \frac{(x-t)^m}{m!}$$

$$= 1 + \frac{\lambda(x-t)}{1!} + \frac{\lambda^2(x-t)^2}{2!} + \frac{\lambda^3(x-t)^3}{3!} + \dots$$

$$= e^{\lambda(x-t)}$$

Note:-

(1) The VIE of 2nd kind

$$u(x) = F(x) + \lambda \int_a^x k(x,t) u(t) dt$$

(ii) $u(x) = F(x) + \lambda \int_0^x R(x,t,t) F(t) dt$

Where $R(x,t,t) = \sum_{m=0}^{\infty} \lambda^{m+1} k_{m+1}(x,t)$

$k_n(x,t)$ is the n^{th} iteration taking $k_1(x,t) = k(x,t)$

(i) $u(x) = \lim_{n \rightarrow \infty} u_n(x)$

Example: (i) $k(x,t) = \frac{2+t \cos x}{2+t \cos t}$ (Similar prob. in exam)

(ii) $k(x,t) = e^{x-t}$

(iii) $k(x,t) = e^{x^2-t^2}$

(iv) $k(x,t) = a^{x-t} \quad a > 0$

* Solution of Fredholm integral Eqn by the method of successive approximation.

The Fredholm integral Eqn of 2nd

$$u(x) = F(x) + \lambda \int_a^b k(x,t) u(t) dt \quad \text{where}$$

$k(x,t) \neq 0$ is real & ctp in the Rectangle R , $a \leq x \leq b$ & $a \leq t \leq b$, $F(x) \neq 0$

has the solⁿ $u(x) = F(x) + \lambda \int_a^b R(x,t,t) F(t) dt$

where $R(x,t,t) = \sum_{n=0}^{\infty} \lambda^n k_{n+1}(x,t)$ is called

Resolvent-Kernel

Let the solⁿ of Eqn (2) in the form of a Power series is

$$u(x) = u_0(x) + \lambda u_1(x) + \lambda^2 u_2(x) + \dots + \lambda^n u_n(x)$$

(3)

$$U_0(x) = F(x)$$

$$U_1(x) = \int_a^b K(x,t) U_0(t) dt \\ = \int_a^b K(x,t) F(t) dt$$

$$U_2(x) = \int_a^b K(x,t) U_1(t) dt$$

Similarly $U_n(x) = \int_a^b K(x,t) U_{n-1}(t) dt$

If we define the iterated kernel

$$K_1(x,t) = K(x,t) \\ K_2(x,t) = \int_a^b K(x,t_1) K_1(t_1,t) dt_1$$

$$K_3(x,t) = \int_a^b K(x,t_1) K_2(t_1,t) dt_1$$

$$K_3(x,t) = \int_a^b \int_a^b K(x,t_1) K_2(t_1,t_2) K_1(t_2,t) dt_1 dt_2$$

Hence, in the general \rightarrow on

$$K_n(x,t) = \int_a^b \dots \int_a^b K(x,t_1) K_1(t_1,t_2) \dots K_{n-1}(t_{n-1},t) K(t_{n-1},t) dt_1 dt_2 \dots dt_{n-1}$$

It follows at once that iterated kernel satisfies $K_n(x,t) = \int_a^b K_n(x,z) K_{n-1}(z,t) dz$

where $r < n$ 4

Now consider the series

$$R(x,t;d) = K_1(x,t) + d K_2(x,t) + \dots \\ + d^{n-1} K_n(x,t) + \dots$$

$$= K(x,t) + d K_2(x,t) + \dots \\ = \sum_{n=0}^{\infty} d^n K_{n+1}(x,t) \quad \text{--- (5)}$$

This series is called the Neumann series of the kernel $K(x,t)$.

Suppose that $F(x)$ & $K(x,t)$ are c/s function in the range of

$$|F(x)| \leq m, \quad a \leq x \leq b$$

$$|K(x,t)| \leq M \quad (a \leq x \leq b, a \leq t \leq b)$$

This series where m, M are positive constant. Now $K_1(x,t) = K(x,t) \leq M$

$$K_2(x,t) = \int_a^b K(x,t_1) K_1(t_1,t) dt_1 \\ \leq \int_a^b M \cdot M dt_1 = M^2(b-a)$$

$$K_3(x,t) = \int_a^b K(x,t_1) K_2(t_1,t) dt_1$$

$$K_3(x,t) \leq M^3 \cdot (b-a)^2$$

in general $K_{n+1}(x,t) \leq M^{n+1} (b-a)^n$ 6

Since $\left| \sum_{n=0}^{\infty} d^n K_{n+1}(x,t) \right| \leq \sum_{n=0}^{\infty} |d|^n M^{n+1} (b-a)^n$
where $|d| < \frac{1}{M(b-a)}$

Hence the solⁿ of eqn (5).

$$U(x) = F(x) + d \int_a^b R(x,t;d) F(t) dt$$

where $R(x,t;d) = \sum_{n=0}^{\infty} d^n K_{n+1}(x,t)$

Find the resolvent kernel of the following

(i) $k(x,t) = (1+x)(1-t)$ with $a=-1, b=0$

(ii) $k(x,t) = e^{(x+t)}$, with $a=0, b=1$

\Rightarrow We have $k_1(x,t) = k(x,t) = (1+x)(1-t)$

①
$$k_2 = \int_0^1 \underbrace{(1+x)(1-z)}_{a} \underbrace{(1+z)(1-t)}_{b} dz$$

$$= (1+x)(1-t) \int_0^1 (1-z^2) dz$$

$$= (1+x)(1-t) \left[z - \frac{z^3}{3} \right]_0^1$$

$$= (1+x)(1-t) \cdot \left[\frac{2}{3} \right]$$

$k_3 = \left(\frac{2}{3}\right)^2 (1+x)(1-t)$

$$k_n(x,t) = \int_0^1 k(x,z) k_{n-1}(z,t) dz$$

$$= \left(\frac{2}{3}\right)^{n-1} (1+x)(1-t)$$

$$R(x,t,t) = \sum_{n=0}^{\infty} d^n k_{n+1}(x,t)$$

$$= \sum_{n=0}^{\infty} d^n \cdot \left(\frac{2}{3}\right)^n (1+x)(1-t)$$

$$= (1+x)(1-t) \left[1 + \left(\frac{2}{3}\right)d + \left(\frac{2}{3}\right)^2 d^2 + \dots \right]$$

$$= \frac{(1+x)(1-t) \cdot 3}{(3-2d)} \quad \text{when } |d| < \frac{3}{2}$$

Solve the integral Eqn

$$u(x) = 1 + d \int_0^1 (1-3zt) u(t) dt$$

Evaluate the resolvent kernel.

\Rightarrow We have $k_1(x,t) = k(x,t) = 1-3xt$

$$k_2(x,t) = \int_0^1 k(x,z) k_1(z,t) dz$$

$$= \int_0^1 (1-3xz)(1-3zt) dz$$

$$= \int_0^1 [1-3z(x+t) + 9xz^2t] dz$$

$$= \left[z - \frac{3}{2}z(x+t) + 3xz^2t \right]$$

$$k_3(x,t) = \int_0^1 k(x,z) k_2(z,t) dz$$

$$= \int_0^1 (1-3xz) \left[z - \frac{3}{2}z(x+t) + 3zt \right] dz$$

$$= \int_0^1 \left[\left(1 - \frac{3}{2}t\right)z - 3z\left(\frac{1}{2}t + x - \frac{3}{2}xt\right) + 9xz^2\left(\frac{1}{2}t\right) \right] dz$$

$$= \frac{1}{4} - \frac{3}{4}xt = \frac{1}{4}(1-3xt)$$

$$\boxed{k_3(x,t) = \frac{1}{4} k_1(x,t)}$$

$$k_4(x,t) = \int_0^1 k(x,z) k_3(z,t) dz$$

$$= \frac{1}{4} \int_0^1 (1-3xz)(1-3zt) dz$$

$$= \frac{1}{4} \int_0^1 [1-3z(x+t) + 9xz^2t] dz$$

$$= \frac{1}{4} k_2(x,t)$$

Similarly

$k_5(x,t) = \left(\frac{1}{4}\right)^2 k_1(x,t)$

$k_6(x,t) = \left(\frac{1}{4}\right)^2 k_2(x,t)$

Now the Resolvent Kernel

$$R(x,t) = \sum_{n=1}^{\infty} A^{n+1} k_n(x,t)$$

Prepared by Kalika

Course finished
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2/4/16

SLEF SECTION

REVISION

Rayleigh Ritz Method for

Boundary Value Problem

Consider the second order LODE —

$$A(x) \frac{d^2 y}{dx^2} + B(x) \frac{dy}{dx} + C(x)y = f(x) \quad (1)$$

with BC $y(x_1) = y_1$ & $y(x_2) = y_2$

then we can construct an integral in the function (*) for which Euler's Equation satisfies

(*) = [Considers a functional along with non-homog. fixed BC or $I = \int_{x_1}^{x_2} F(x, y, y') dx$ when $\frac{dy}{dx} = y'$, $y(x_1) = y_1$, & $y(x_2) = y_2$]

Functional for DE (1)

$$F = p(x)y'^2 + q(x)y^2 + r(x)y$$

By Euler's Eqⁿ. $\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$

$$\Rightarrow q(x) \cdot 2y + r(x) - \frac{d}{dx} [2y' \cdot p(x)] = 0$$

$$\Rightarrow 2p(x) \cdot y'' - 2q(x)y - r(x) + 2y' \cdot p'(x) = 0$$

$$\Rightarrow 2 [p(x)y'' + p'(x)y' - q(x)y] = r(x)$$

$$\text{or } p(x)y'' + p'(x)y' - q(x)y = \frac{r(x)}{2} \quad (2)$$

$$\text{or } y'' + \frac{p'(x)}{p(x)} y' - \frac{q(x)}{p(x)} y = \frac{r(x)}{2p(x)}$$

$$\therefore y' = e^{\int \frac{p'(x)}{p(x)} dx} = e^{\log(p(x))} = p(x)$$

Now multiply eqⁿ (1) by arbitrary $f^n h(x)$ then compare with (2), we get —

$p(x) = A(x) h(x)$	$q(x) = -C(x) \cdot h(x)$
$p'(x) = B(x) h(x)$	$r(x) = 2 f(x) \cdot h(x)$

Example: Solve the BVP —

$$x^2 \frac{d^2 y}{dx^2} + x \frac{dy}{dx} - y = 2x^2 \quad \text{(By Ritz Method)}$$

Subject to the conditions $y(0) = 0$ & $y(1) = 1$

Solⁿ (1) $\Rightarrow y'' + \frac{1}{x} y' - \frac{y}{x^2} = 2$

\therefore one of its solⁿ is given by $IF = e^{\int \frac{1}{x} dx} = x$
by $u = x$
 $u' = 1$

Solving by Ritz Method

The corresponding functional is given by

(2) $I = \int_0^1 (p y'^2 + q y^2 + r y) dx$, $y(0) = 0$, $y(1) = 1$

Now comparing (1) with

$$a(x) \frac{d^2 y}{dx^2} + b(x) \frac{dy}{dx} + c(x) y = f(x)$$

We have $a(x) = x^2$, $b(x) = x$, $c(x) = -1$, $f(x) = 2x^2$

Now using

$$h(x) = \frac{p(x)}{a(x)} = \frac{x}{x^2} = \frac{1}{x}$$

$$p(x) = a(x) h(x)$$

$$p'(x) = b(x) h(x)$$

$$q(x) = -c(x) h(x)$$

$$r(x) = 2 f(x) h(x)$$

$$p(x) = x^2 \cdot \frac{1}{x} = x$$

$$q(x) = 1 \cdot h(x) = \frac{1}{x}$$

$$r(x) = 4x^2 \cdot h(x) = 4x$$

then $p(x) = e^{\int \frac{b(x)}{a(x)} dx} = e^{\int \frac{1}{x} dx} = e^{\ln x} = x$

\therefore (2) becomes

$$I = \int_0^1 (x y'^2 + \frac{y^2}{x} + 4x y) dx$$

$y(0) = 0$
 $y(1) = 1$

Let $\phi_1(x) = x(1-x)$

$\phi_2(x) = x$ be two solⁿs of (1)

then approximate solⁿ is given by

$$y = u_0(x) + c_1(x-0)(x-1), \quad u_0(0) = 0, u_0(1) = 1$$

(1) ST $\phi(x) = x e^x$ is a solⁿ of the Volterra integral eqⁿ

$$\phi(x) = \sin x + 2 \int_0^x \cos(x-t) \phi(t) dt$$

Solⁿ taking R.H.S $\phi(t) = t e^t$

$$\sin x + 2 \int_0^x (\cos x \cdot \cos t + \sin x \cdot \sin t) (t \cdot e^t) dt$$

let $x-t = p \Rightarrow t = x+p$
 $\Rightarrow dt = -dp$

$$\sin x + 2 \int_0^x \cos p (x+p) \cdot e^{(x+p)} \cdot (-dp)$$

$$= \sin x + 2 \int_0^x \cos(x+p) (x+p) \cdot e^x \cdot e^p dp$$

$$= \sin x + 2 \int_0^x e^x \int_0^p (x+p) \cos(x+p) e^p dp$$

$$= \sin x + 2 e^x \left[x \int_0^p \cos p \cdot e^p dp + \int_0^p p \cos p \cdot e^p dp \right]$$

Formulas $\int e^t \sin t dt = \frac{e^t}{2} (\sin t - \cos t)$
 $\int e^t \cos t dt = \frac{e^t}{2} (\sin t + \cos t)$

$$= \sin x + 2 e^x \left\{ x \cdot \frac{e^p}{2} (\sin p + \cos p) + p \cdot \frac{e^p}{2} (\sin p + \cos p) - \int_0^p \frac{e^p}{2} (\sin p + \cos p) dp \right\}$$

$$= \sin x + e^x \left\{ x e^p (\sin p + \cos p) + p e^p (\sin p + \cos p) - \frac{e^p}{2} (\sin p - \cos p) - \frac{e^p}{2} (\sin p + \cos p) \right\}$$

$$= \sin x + e^x \left[(\sin p + \cos p) (x e^p + p e^p + \frac{e^p}{2}) - \frac{e^p}{2} (\sin p - \cos p) \right]$$

$$= \sin x + e^x \left[\frac{1}{2} \left(x + \frac{1}{2} \right) + \frac{1}{2} \right] - \left(\cos x - \sin x \right) \frac{x e^x}{2} - \frac{x e^x + e^x}{2} + \frac{e^{-x}}{2} (\sin x + \cos x)$$

$$= \sin x + e^x \left[x + 1 - \frac{e^{-x}}{2} \cos x + \frac{e^{-x}}{2} \sin x - \frac{e^{-x}}{2} \cos x \right]$$

$$= \sin x + e^x [x + 1 - e^{-x} \cos x]$$

$$= \sin x + x e^x + e^x - \cos x e^x$$

Pending

Prepared by Kalika

Laplace transform of the nth order
fun: derivative of f(x)

Statement: Let f(x) & its derivatives f'(x), f''(x), ..., f⁽ⁿ⁻¹⁾(x) are ctg for t=0 & of exponential order as t → ∞. and if f⁽ⁿ⁾(t) is of class A & if h(t) = f(x) then Laplace transform of f⁽ⁿ⁾(t) exists when s > a & given by —

from p-29

$$L\{f^{(n)}(t)\} = s^n L\{f(t)\} - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - s f^{(n-1)}(0)$$

Pf: We have for f'(t)
L[f'(t)] = s L[f(t)] - f(0) — (1)

then for f''(t), we have

$$L[f''(t)] = s^2 L[f(t)] - s f'(0) - f(0)$$

$$i.e. L[f''(t)] = s L[f'(t)] - f'(0)$$

$$= s [s L[f(t)] - f(0)] - f'(0)$$

$$= s^2 L[f(t)] - s f(0) - f'(0)$$

for f'''(t),

$$L[f'''(t)] = s L[f''(t)] - f''(0)$$

$$= s [s^2 L[f(t)] - s f(0) - f'(0)] - f''(0)$$

$$= s^3 L[f(t)] - s^2 f(0) - s f'(0) - f''(0)$$

proceeding in similar way, we have

$$L\{f^{(n)}(t)\} = s^n L\{f(t)\} - \sum_{r=0}^{n-1} s^{n-1-r} f^{(r)}(0)$$

② cosine integral

$p = -3/2$

$$C(t) = \int_1^{\infty} \frac{\cos ut}{u} du, \quad L[C(t)] = ??$$

$$\Rightarrow C'(t) = -\frac{\cos t}{t}$$

$$\Rightarrow -t C'(t) = \cos t$$

$$\Rightarrow -L[t C'(t)] = L[\cos t] = \frac{s}{s^2+1}$$

$$\Rightarrow L[t C'(t)] = -\frac{s}{s^2+1} \quad \text{--- (1)}$$

Now

$$\therefore L[t f(t)] = -\frac{d}{ds} f(s)$$

where $L[f(t)] = f(s)$

$$= -\frac{d}{ds} L[f(t)]$$

$$\therefore L[t C'(t)] = -\frac{d}{ds} L[C(t)]$$

$$= -\frac{d}{ds} \left[\frac{s}{s^2+1} \right] \quad \text{--- (1)}$$

by (1)

$$\frac{s}{s^2+1} = \frac{d}{ds} [s L[C(t)] - C(0)]$$

$$= \frac{d}{ds} [s L[C(t)]] - 0$$

Because $C(0) = \int_1^{\infty} \frac{\cos u}{u} du = 0$

$$\Rightarrow \frac{d}{ds} [s L[C(t)]] = \frac{s}{s^2+1}$$

Now integrating both side, we get

$$s L[C(t)] = \frac{1}{2} \int \frac{2s}{s^2+1} ds$$

$$= \frac{1}{2} \log(s^2+1) + C \quad \text{--- (1)}$$

$$L[C(t)] = \frac{1}{2s} \log(s^2+1) \quad \text{--- (1)}$$

in (1) & y final value

$$\lim_{s \rightarrow 0} s f(s) = \lim_{s \rightarrow 0} f(s) = 0$$

$$\therefore C(0) = 0$$

$$L[C(t)] = \frac{1}{2s} \log(s^2+1)$$

③ To find inverse Laplace transform by convolution

① Use convolution to evaluate

$$L^{-1} \left[\frac{s}{(s^2+4)^2} \right] = L^{-1} \left[\frac{1}{(s^2+4)} \cdot \frac{s}{(s^2+4)} \right]$$

$$\text{let } f(s) = \frac{s}{2(s^2+4)}, \quad g(s) = \frac{s}{(s^2+4)}$$

$$\therefore L^{-1}[f(s)] = \frac{1}{2} \sin 2t$$

$$L^{-1}[g(s)] = \cos 2t$$

∴ A/c to convolution theorem

$$L^{-1}[f(s)g(s)] = \int_0^t f(u)g(t-u)du$$

$$= \frac{1}{2} \int_0^t \sin 2u \cdot \cos 2(t-u) du$$

$$= \frac{1}{2} \int_0^t [\sin 2u \cos 2t - 2u \sin 2u \sin 2t + \cos 2u \cos 2t] du$$

$$= \frac{1}{4} \left[\int_0^t \sin 2u du + \int_0^t \sin(4u-2t) du \right]$$

$$= \frac{1}{4} \left[u \sin 2t + \cos(4u-2t)/4 \right]_0^t$$

$$\frac{1}{4} \sin 2t = \frac{1}{4} [t \sin 2t + \cos(2t) - \cos(2t)]$$

② $L^{-1} \left[\frac{1}{(s^2+1)^2} \right]$

let $g(s) = \frac{1}{(s^2+1)}$ = $f(s)$

$g(t) = L^{-1}[g(s)] = \frac{1}{1} \sin t = L^{-1}[f(s)] = f(t)$

∴ By convolution theorem

$L^{-1}[f(s)g(s)] = \int_0^t \frac{1}{a} \sin at \cdot \frac{1}{a} \sin a(t-u) du$

$= \frac{1}{a^2} \left[\int_0^t \sin at dt \right] = \frac{1}{a^2} \int_0^t (1 - \cos at) dt$

③

$\frac{dx}{dt} - y = e^t$ — (1) ($x(0) = 1, y(0) = 2$)

$L[x(t)] - L[y] = L[e^t]$

$\Rightarrow sL[x(t)] - x(0) - L[y] = \frac{1}{s-1}$

$sL[x(t)] - 1 - L[y] = \frac{1}{s-1}$ — (1) $\times s$

$\frac{dx}{dt} + x = \sin t \Rightarrow s\bar{x} - \bar{y} = \frac{1}{s-1} + 1$

$L[y(t)] + L[x] = L[\sin t] = \frac{1}{s+1}$

$\Rightarrow sL[y(t)] + y(0) + L[x] = \frac{1}{s+1}$

$\Rightarrow sL[y(t)] + L[x] = \frac{1}{s^2+1}$ — (11) $\times s$

$L[x] (s^2+1) - 1 = \frac{1}{s} + \frac{1}{s+1}$

$L[x] =$

$\frac{s\bar{x} - 1}{s^2+1} = \frac{1}{s} + \frac{1}{s+1}$

$\bar{x} + s\bar{y} = \frac{1}{s+1}$

$\bar{x}(s^2+1) = \frac{1}{s} + \frac{1}{s+1}$

$\bar{x} = \frac{1}{(s+1)(s^2+1)} + \frac{1}{(s^2+1)^2}$

$\bar{y} = \frac{1}{(s^2+1)^2} - \frac{1}{(s-1)(s^2+1)}$

From here, we solve \bar{x} & \bar{y} by taking inverse Laplace transform and get x & y , that is the solⁿ of the given qn.

④ Solve the DE using Laplace transform method.

$\frac{\partial y}{\partial t} = 3 \frac{\partial^2 y}{\partial x^2}$ when

$y(0, t) = 0,$
 $\left. \frac{\partial y}{\partial x} \right|_{x=0} = 0,$
 $y(x, 0) = 30 \cos 5x$

solⁿ given equation is

$\frac{\partial y}{\partial t} = 3 \frac{\partial^2 y}{\partial x^2}$

taking Laplace transform on both sides, we

$L[y(t)] = 3 L[y''(t)]$

$L[y] = \bar{y}$ $\Rightarrow s\bar{y} - y(0) = 3 [s^2\bar{y} - s y(0) - y'(0)]$

$\Rightarrow s\bar{y} - y(x, 0) = L \left[3 \frac{\partial^2 y}{\partial x^2} \right] = 3 \frac{\partial^2 L[y]}{\partial x^2}$

$s\bar{y} - y(x, 0) = 3 \frac{\partial^2 \bar{y}}{\partial x^2}$

$\Rightarrow 3 \frac{d^2 \bar{y}}{dx^2} = s\bar{y} = -y(x, 0) = -30 \cos 5x$

$\Rightarrow (3D^2 - s)\bar{y} = -30 \cos 5x$

$$(D^2 - 8/3)y = -10 \cos 5x$$

Solⁿ is given as
 $y = CF + PI$

$$m^2 - 8/3 = 0 \Rightarrow m = \pm \sqrt{8/3}$$

$$\therefore CF = C_1 e^{x\sqrt{8/3}} + C_2 e^{-x\sqrt{8/3}}$$

$$P.I = \frac{(-10 \cos 5x)}{(D^2 - 8/3)} = \frac{30 \cos 5x}{75 + 9}$$

$$\bar{y} = C_1 e^{x\sqrt{8/3}} + C_2 e^{-x\sqrt{8/3}} + \frac{30 \cos 5x}{75 + 9} \quad \text{--- (i)}$$

Now

$$\left. \frac{\partial y}{\partial x} \right|_{x=0} = 0 \quad \text{--- (ii)}$$

$$L\left[\frac{\partial y}{\partial x}\right] = h[0] \Rightarrow \frac{\partial \bar{y}}{\partial x} = 0$$

Again, $y(\pi/2, t) = 0 \Rightarrow h[y(\pi/2, t)] = 0$

$$\Rightarrow \bar{y}(\pi/2, 0) = 0 \quad \text{--- (iii)}$$

diff (i) w.r.t x,

$$\frac{\partial y}{\partial x} = C_1 \sqrt{8/3} e^{x\sqrt{8/3}} - C_2 \sqrt{8/3} e^{-x\sqrt{8/3}} - \frac{150 \sin 5x}{(75+9)}$$

$$= \sqrt{8/3} \left[C_1 e^{x\sqrt{8/3}} - C_2 e^{-x\sqrt{8/3}} \right] - \frac{150 \sin 5x}{75+9}$$

putting $x=0 \Rightarrow \frac{\partial \bar{y}}{\partial x} = 0 \quad \text{(by (ii))}$

$$0 = \sqrt{8/3} [C_1 - C_2] - 0 \Rightarrow C_1 = C_2$$

$$\bar{y} = C_1 \left[e^{x\sqrt{8/3}} + e^{-x\sqrt{8/3}} \right] - \frac{150 \sin 5x}{75+9}$$

Now subject to the cond.

$$\bar{y}(\pi/2, 0) = 0 \quad \text{(by iii)}$$

by (i) $0 = C_1 \left[e^{\pi/2 \sqrt{8/3}} + e^{-\pi/2 \sqrt{8/3}} \right] + \frac{30 \cos(15\pi/2)}{75+9}$

$$- C_1 \left[\dots \right] = 0 \Rightarrow C_1 = 0$$

$$L[y] = \bar{y} = \frac{30 \cos 5x}{75+9} \quad \text{(by (i))}$$

taking inverse lap

$$y = L^{-1} \left[\frac{30 \cos 5x}{75+9} \right] = 30 \cos 5x \cdot e^{-75t}$$

II - UNIT (Also p-84)

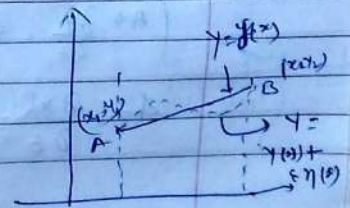
(1) Euler's Equation

The necessary condition for $f(x, y, z)$ to be an extremum is that $\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0$

which is called Euler's Equation.

pf.

Let $y = y(x)$ be a curve joining the pts. (x_1, y_1) & (x_2, y_2) .



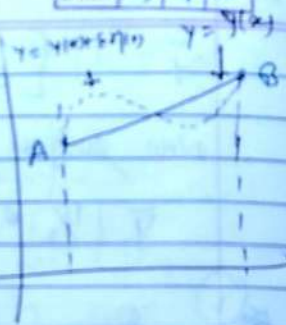
$$I = \int_a^b f(x, y, y') dx$$

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which makes I an extremum or stationary.

Let $y = y(x) + \epsilon \eta(x)$ be a neighbouring curve joining these points.

where ϵ is small and $\eta(x)$ is an arbitrary differentiable fⁿ of x satisfying $\eta(x_1) = 0$ & $\eta(x_2) = 0$ at A & B respectively.



From the value along $y = y(x) + \epsilon \eta(x)$

$$I(\epsilon) = \int_{x_1}^{x_2} f[x, y(x) + \epsilon \eta(x), y'(x) + \epsilon \eta'(x)] dx \quad \text{--- (1)}$$

this being a fⁿ of ϵ is maximum or minimum for $\epsilon = 0$. When

$$\frac{dI(\epsilon)}{d\epsilon} = 0 \text{ at } \epsilon = 0 \quad \text{--- (2)}$$

Now differentiating $I(\epsilon)$ under the integral sign by Leibnitz's rule, we have

$$\frac{dI(\epsilon)}{d\epsilon} = \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial \epsilon} + \frac{\partial f}{\partial y} \frac{\partial y}{\partial \epsilon} + \frac{\partial f}{\partial y'} \frac{\partial y'}{\partial \epsilon} \right] dx \quad \text{--- (3)}$$

Since ϵ is independent of x .

$$\therefore \frac{\partial x}{\partial \epsilon} = 0 \text{ Also } \frac{\partial y}{\partial \epsilon} = \eta(x)$$

$$\left(\text{As } \frac{\partial y}{\partial \epsilon} = \frac{\partial y(x)}{\partial \epsilon} + \frac{\partial (\epsilon \eta(x))}{\partial \epsilon} = \eta(x) \right)$$

$$\therefore \frac{\partial y'}{\partial \epsilon} = \eta'(x) \text{ with } y = y(x) + \epsilon \eta(x)$$

Substituting these values in eqn. (3), we get

$$\begin{aligned} \frac{dI(\epsilon)}{d\epsilon} &= \int_{x_1}^{x_2} \left[0 + \frac{\partial f}{\partial y} \eta(x) + \frac{\partial f}{\partial y'} \eta'(x) \right] dx \\ &= \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial y} \eta(x) + \left[\frac{\partial f}{\partial y'} \eta'(x) - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \eta(x) \right] \right] dx \\ &= \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial y} \eta(x) \right] dx + 0 - \int_{x_1}^{x_2} \left[\frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \eta(x) \right] dx \end{aligned}$$

$$\left[\because \eta(x_1) = 0 \text{ \& } \eta(x_2) = 0 \text{ at A \& B} \right]$$

$$= \int_{x_1}^{x_2} \left[\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \right] \eta(x) dx \quad \text{--- (4)}$$

Now from eqn (2) & (4), we get

$$\left[\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) \right] = 0 \quad \text{--- (5)}$$

Which is the desired Euler's Eqn. Second form of Euler's Eqn

II. Another form of Euler's Equation

the Euler's Equation is

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \quad \text{(by (5))}$$

$$\Rightarrow y' \frac{\partial f}{\partial y} - y' \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \quad \text{(multiply by } y')$$

$$\Rightarrow y' \frac{\partial f}{\partial y} - \left[\frac{d}{dx} (y' \frac{\partial f}{\partial y'}) - \left(\frac{\partial f}{\partial y'} \frac{dy'}{dx} \right) \right] = 0$$

$$\Rightarrow \frac{\partial f}{\partial x} + y' \frac{\partial f}{\partial y} + \frac{\partial f}{\partial y'} \frac{dy'}{dx} - \frac{d}{dx} \left(y' \frac{\partial f}{\partial y'} \right) - \frac{\partial f}{\partial x}$$

$$\Rightarrow \frac{\partial f}{\partial x}$$

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$f(x, y, y')$ -

$$\frac{df}{dx} = \frac{\partial f}{\partial x} + \frac{\partial f}{\partial y} \frac{dy}{dx} + \frac{\partial f}{\partial y'} \frac{dy'}{dx}$$

$$\Rightarrow \frac{df}{dx} - \frac{d}{dx} \left(y' \frac{\partial f}{\partial y'} \right) = \frac{\partial f}{\partial x}$$

$$\Rightarrow \frac{d}{dx} \left[f - y' \frac{\partial f}{\partial y'} \right] = \frac{\partial f}{\partial x}$$

$$\Rightarrow \frac{d}{dx} \left[f - y' \frac{\partial f}{\partial y'} \right] - \frac{\partial f}{\partial x} = 0 \quad \text{--- (6)}$$

Which second form of Euler's Equation

III. Particular Cases of Euler's Eqⁿ

The Euler's Equation is

$$\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \quad \text{--- (1)}$$

Every solⁿ of Euler's Eqⁿ which satisfy the BC is called a stationary fⁿ or ~~extremal~~ extremal of the problem

(a) If f is independent of y ,

then $\frac{\partial f}{\partial y} = 0$, so by (1),

$$\frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \Rightarrow \frac{\partial f}{\partial y'} = c$$

(b) If f is independent of y' ,

then $\frac{\partial f}{\partial y'} = 0$, then by eqⁿ (1),

which gives solⁿ directly.

(iii) If f is independent of x .

then $\frac{\partial f}{\partial x} = 0$, then eqⁿ (6) gives

$$\frac{d}{dx} \left[f - y' \frac{\partial f}{\partial y'} \right] = 0$$

$$\Rightarrow f - y' \frac{\partial f}{\partial y'} = c$$

for this we use second form

(2) Find the extremals of the following functions

(i) $\int_{x_0}^{x_1} (x+y)y' dx$

Here $f(x, y, y') = (x+y)y'$ which is independent of y .

so $\frac{\partial f}{\partial y} = 0$ $\left(\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0 \right)$

$$0 - \frac{d}{dx} (x + y') = 0$$

$$\Rightarrow x + y' = c \Rightarrow y' = \frac{c-x}{2}$$

on integrate

$$\Rightarrow y = \frac{c}{2}x - \frac{x^2}{4} + d$$

(ii) $\int_{x_0}^{x_1} (y^2 + y' - 2y \sin x) dx$

Here $f = y^2 + y' - 2y \sin x$

Euler's Equation $\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0$ --- (1)

$$(2y - 2 \sin x) - \frac{d}{dx} (2y') = 0$$

$$\Rightarrow 2 \frac{d}{dx} (y') = 2y - 2 \sin x$$

$$\Rightarrow y'' - y = -\sin x$$

$$(D^2 - 1)y = -\sin x$$

Solⁿ $y = CF + PI$

$\therefore D^2 - 1 = 0 \Rightarrow D = \pm 1$

CF = $C_1 e^x + C_2 e^{-x}$

P.I = $\frac{-\sin x}{(D^2 - 1)} = \frac{1}{2} \sin x$

$(1 - D^2)^{-1}(\sin x) = (1 + D^2 + D^4 + \dots) \sin x = \sin x$

$\therefore d(\sin x) = \cos x$	$d(-\sin x) = -\cos x$
$d(\cos x) = -\sin x$	$d(-\cos x) = \sin x$

$\therefore P.I = \frac{\sin x + \sin x + \sin x - (\cos x) + \dots}{2}$

$\therefore Y = C_1 e^x + C_2 e^{-x} + \frac{1}{2} \sin x$

(c) $\int_{x_0}^{x_1} \frac{1+y^2}{y^2} dx$, $F(x, y, y') = \frac{1+y^2}{y^2}$

1-Method \Rightarrow second form of Euler's eqn

$\frac{d}{dx} \left(f - y' \frac{\partial f}{\partial y'} \right) - \frac{\partial f}{\partial x} = 0$ ($\because \frac{\partial f}{\partial x} = 0$)

$\Rightarrow \frac{1+y^2}{y^2} - y' \left((1+y^2) \cdot (-2) \cdot y^{-3} \right) = 0$

$\Rightarrow \frac{1+y^2}{y^2} + y' \frac{2(1+y^2)}{y^3} = 0$

$\frac{(1+y^2)}{y^2} = 0 \Rightarrow y' = -\frac{1+y^2}{y}$

$\Rightarrow \int \frac{dy}{\sqrt{1+y^2}} = \frac{1}{\sqrt{c}} \int dx = \frac{x}{\sqrt{c}} + d$

$\Rightarrow \sinh^{-1}(y) = \frac{x}{\sqrt{c}} + d$

$y = \sinh^{-1} \left(\frac{x}{\sqrt{c}} + d \right)$

(d) $\int_{x_0}^{x_1} [y^2 - y'^2 - 2y \sin x] dx$, $y(0) = 0$
 $y(\pi/2) = 1$

$f = y^2 - y'^2 - 2y \sin x$

\therefore By Euler's equation $\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0$

$(2y - 2 \sin x) - \frac{d}{dx} (-2y') = 0$

$\Rightarrow [y'' + y = \sin x] \Rightarrow (D^2 + 1)y = \sin x$

$\therefore y = CF + PI$ $(D^2 + 1) = 0$

$\Rightarrow D = \pm i$

CF, $y = C_1 \cos x + C_2 \sin x$

PI $\Rightarrow \frac{1}{(D^2 + 1)} \sin x = \frac{-x}{2} \cos x$

$\therefore y = C_1 \cos x + C_2 \sin x + \frac{1}{2} \sin x - \frac{x}{2} \cos x$

$y(0) = 0 = C_1 \Rightarrow C_1 = 0$

$y(\pi/2) = 1 = C_2 + 1/2 \Rightarrow C_2 = 1/2$

$y = \sin x + \frac{x}{2} \cos x$

(e) $\int_0^1 (xy + y^2 - 2y^2 y') dx$, $y(0) = 1$, $y(1) = 2$

$f(x, y, y') = xy + y^2 - 2y^2 y'$

$\frac{\partial f}{\partial y} - \frac{d}{dx} \left(\frac{\partial f}{\partial y'} \right) = 0$

$\Rightarrow (x + 2y - 4y y') - \frac{d}{dx} (-2y^2) = 0$

$\Rightarrow x + 2y - 4y y' + 2 \cdot 2y \cdot y' = 0$

$\Rightarrow 4y y' + 2y = 0$

$y(0) = 1 \Rightarrow y = -x/2$

Since B.C not satisfied.

Hence extremum can't be achieved.

III. 3rd form of Euler's eqn

(90) Since $\frac{\partial F}{\partial y'}$ is a fⁿ of x, y, y' , we have

$$\frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = \frac{\partial}{\partial x} \left(\frac{\partial F}{\partial y'} \right) + \frac{\partial}{\partial y} \left(\frac{\partial F}{\partial y'} \right) \frac{dy}{dx} + \frac{\partial}{\partial y'} \left(\frac{\partial F}{\partial y'} \right) \frac{dy'}{dx}$$

$$\frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = \frac{\partial^2 F}{\partial x \partial y'} + y' \frac{\partial^2 F}{\partial y \partial y'} + y'' \frac{\partial^2 F}{\partial y'^2}$$

Since Euler's Equation is -

$$\frac{\partial F}{\partial y} - \frac{d}{dx} \left(\frac{\partial F}{\partial y'} \right) = 0$$

$$\Rightarrow \frac{\partial F}{\partial y} - \frac{\partial^2 F}{\partial x \partial y'} - y' \frac{\partial^2 F}{\partial y \partial y'} - y'' \frac{\partial^2 F}{\partial y'^2} = 0$$

which is third form of Euler's eqn

(3) If (x_1, θ_1) & (x_2, θ_2) are two pts in the plane then the length of arc of the curve joining these pts is given by

$$I[x, \theta] = \int_{\theta_1}^{\theta_2} (r^2 + r'^2)^{1/2} d\theta, \quad r' = \frac{dr}{d\theta}$$

by minimizing the above integral obtain the eqn of pt. line in the polar coordinates

solⁿ Comparing the given functional with $\int_{\theta_1}^{\theta_2} F(\theta, r, r') d\theta$, we get

$$F(\theta, r, r') = (r^2 + r'^2)^{1/2}$$

By Euler's eqn

$$\frac{\partial F}{\partial r} - \frac{d}{d\theta} \left(\frac{\partial F}{\partial r'} \right) = 0 \quad \text{--- (1)}$$

$$\frac{\partial F}{\partial r} = \frac{1}{2} \cdot \frac{2r}{\sqrt{r^2 + r'^2}}$$

$$+ \frac{d}{d\theta} \left(\frac{\partial F}{\partial r'} \right) = \frac{r'}{\sqrt{r^2 + r'^2}}$$

By (1)

$$\frac{r}{\sqrt{r^2 + r'^2}} - \frac{d}{d\theta} \left(\frac{r'}{\sqrt{r^2 + r'^2}} \right) = 0$$

$$\frac{r}{\sqrt{r^2 + r'^2}} - \frac{r''(\sqrt{r^2 + r'^2}) - r' \cdot \frac{1}{2} \cdot \frac{2r'r'}{\sqrt{r^2 + r'^2}}}{(r^2 + r'^2)^{3/2}} = 0$$

$$\frac{r}{\sqrt{r^2 + r'^2}} - \frac{r''(\sqrt{r^2 + r'^2}) - r'r' \cdot \frac{r'r'}{\sqrt{r^2 + r'^2}}}{(r^2 + r'^2)^{3/2}} = 0$$

$$r^3 + r'^2 r - [r^2 r'' + r'^2 r'' - r'r' \cdot \frac{r'r'}{\sqrt{r^2 + r'^2}} - r'r'^2 r''] = 0$$

$$\Rightarrow r^3 + r'^2 r - r^2 r'' - r'^2 r'' + r'r' \cdot \frac{r'r'}{\sqrt{r^2 + r'^2}} + r'r'^2 r'' = 0$$

$$\Rightarrow r^2 + 2r'r'^2 - r^2 r'' = 0$$

$$\Rightarrow r^2 + 2r'r'^2 - r'r'' = 0$$

$$\text{let } p = r' \Rightarrow r'' = \frac{d^2 r}{d\theta^2} = \frac{d}{d\theta} \left(\frac{dr}{d\theta} \right) = \frac{dp}{d\theta} \cdot \frac{d\theta}{dr} = \frac{dp}{dr} \cdot \frac{dr}{d\theta} = \frac{dp}{dr} \cdot p$$

$$\Rightarrow r^2 + 2r p^2 - r p \frac{dp}{dr} = 0$$

$$\frac{dp}{dr} = \frac{r^2 + 2r p^2}{r p} = \frac{r}{p} + 2p/r \quad \text{--- (5)}$$

$$\text{let } p/r = x \Rightarrow p = rx$$

$$\Rightarrow \frac{dp}{dr} = x + r \frac{dx}{dr}$$

By (5) =

$$x + r \frac{dx}{dr} = 2x + \frac{1}{x} \Rightarrow r \frac{dx}{dr} = x + \frac{1}{x}$$

$$\Rightarrow \frac{dr}{r} = \frac{x dx}{x^2 + 1} \Rightarrow \log r = \frac{1}{2} \log(x^2 + 1) + \log c$$

$$\Rightarrow \log r = \frac{1}{2} \log(x^2 + 1) + \log c$$

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$$\frac{v^2}{c^2} = (x^2 + 1) \Rightarrow r = c^2 \sqrt{x^2 + 1} = \sqrt{p^2 + c^2}$$

$$x = \frac{\sqrt{r^2 - c^2}}{r} \quad r = c^2 \sqrt{x^2 + 1} \Rightarrow p = \sqrt{\frac{r^2 - c^2}{c^2}} = 1$$

$$\frac{p}{\delta} = \frac{\sqrt{r^2 - c^2}}{c} \Rightarrow p = \frac{r^2 - c^2}{c}$$

$$p = \frac{r^2 - c^2}{c} \Rightarrow \frac{dr}{dt} = \frac{c^2}{\sqrt{r^2 - c^2}}$$

$$\frac{dr}{dt} = \frac{c^2}{\sqrt{r^2 - c^2}} \Rightarrow \frac{dr}{\sqrt{r^2 - c^2}} = \frac{dt}{c} \Rightarrow \log \left| \frac{r + \sqrt{r^2 - c^2}}{c} \right|$$

$$\Rightarrow \frac{dr}{r\sqrt{r^2 - c^2}} = \frac{1}{c} dt$$

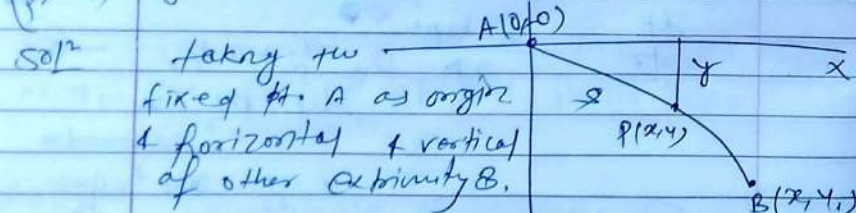
$$\Rightarrow \frac{1}{c} + d = c \log \left(\frac{r + \sqrt{r^2 - c^2}}{c} \right)$$

$$\Rightarrow \sec(\theta + d) = r/c$$

$$r = c \sec(\theta + d)$$

or $c = r \cos(\theta + d)$
When $c + d \neq$ arbitary constant

Q. Find the curve connecting the pts A & B which is traversed by a particle sliding from A to B, in the shortest time.



At time t , let v be the velocity in the position P and actual dist. from origin is ρ .

By Energy Eqⁿ

Change in KE = Work done in moving particle from A to B,

$$W = KE + PE \Rightarrow \text{avg } y = \frac{1}{2} vt + t_0$$

$$\Rightarrow v = \sqrt{2gy} \quad (\because dv = \frac{dy}{dt})$$

$$\Rightarrow dt = \frac{1}{\sqrt{2gy}} dy$$

Thus time taken by the particle from A to B is

$$t = \int_0^x \frac{dy}{\sqrt{2gy}} = \frac{1}{\sqrt{2g}} \int_0^x \frac{1}{\sqrt{y}} dy$$

$$= \int_0^x f(x,y) dx \text{ when } f = \frac{1}{\sqrt{2g}} \frac{1}{\sqrt{y}}$$

Since f is independent of $x \Rightarrow \frac{\partial F}{\partial x} = 0$
By Second form of Euler's eqⁿ

$$f - y \frac{\partial f}{\partial y} = C$$

$$\Rightarrow \frac{1}{\sqrt{2g}} \frac{1}{\sqrt{y}} - y \left(\frac{1}{\sqrt{2g}} \cdot \frac{1}{2} \frac{1}{y^{3/2}} \right) = C$$

$$\Rightarrow \frac{1}{\sqrt{2g}} \left(\frac{1 + y^2 - y^2}{\sqrt{1 + y^2}} \right) = C$$

$$\Rightarrow \frac{1}{\sqrt{1 + y^2}} = C \cdot \sqrt{2g}$$

$$\Rightarrow y(1 + y^2) = \frac{1}{c^2 2g} = 2k, \quad k = \text{const}$$

$$\Rightarrow y(1+y') = 2k$$

$$\left(\therefore \text{then we have } \frac{dy}{dx} = y' = \tan \theta \right)$$

$$\Rightarrow y(1 + \tan^2 \theta) = 2k$$

$$\Rightarrow y \sec^2 \theta = 2k$$

$$\Rightarrow y = 2k \cdot \cos^2 \theta = k(1 + \cos 2\theta)$$

Now $\therefore \frac{dx}{dy} = \frac{dy}{dx} = \tan \theta \Rightarrow \frac{dx}{dy} = \cot \theta$

$$\begin{aligned} \Rightarrow dx &= \cot \theta \cdot dy \\ &= \cot \theta \cdot k \cdot (-\sin 2\theta) \cdot 2 \\ &= -2k \cot \theta (2 \sin \theta \cos \theta) d\theta \\ &= -4k \cos^2 \theta d\theta \\ &= -2k(1 + \cos 2\theta) d\theta \end{aligned}$$

$$\therefore x = -2k \left[\theta + \frac{\sin 2\theta}{2} \right] + C$$

$$= C - 2k \left(\theta + \frac{\sin 2\theta}{2} \right)$$

(Here we find value of x & y in this problem)

(99) Isoperimetric problem

It is necessary to make a given integral

$$I = \int_{x_1}^{x_2} f(x, y, y') dx \text{ subject to}$$

$$\text{the constraint } J = \int_{x_1}^{x_2} g(x, y, y') dx = \text{const}$$

Method

Such problem involve one or more constraints. Conditions. This type of problems are called Isoperimetric problem.

Rayleigh-Ritz method

(1) Consider the two-pt. BVP -
 $y'' + x = 0, 0 \leq x < 1, y(0) = y(1) = 0$

Solⁿ

(2) Solve the BVP, defined below
 $y'' + y = -x, 0 < x < 1$ - (1)
with $y(0) = 0 = y(1)$

Solⁿ

To find the approximate solⁿ by Rayleigh Ritz method, we take functional

$$I(y) = \int_0^1 (yy'' + y^2 + 2yx) dx \quad \text{--- (2)}$$

Let an approximate solⁿ is

$$\text{given by } y(x) = \sum_{i=1}^n \alpha_i \phi_i(x) \quad \text{--- (3)}$$

where base fn ϕ_i satisfies
B.C. $\phi_i(0) = 0 = \phi_i(1) \quad \forall i$

Now substituting (3) in (2), we have

$$I(y) = \int_0^1 \left[\sum_{i=1}^n \alpha_i \phi_i(x) \right] \left[\sum_{i=1}^n \alpha_i \phi_i''(x) + \sum_{i=1}^n \alpha_i \phi_i^2(x) + 2x \right] dx \quad \text{--- (4)}$$

$$\text{let } p_i = \int_0^1 x \cdot \phi_i(x) dx \quad \text{--- (5)}$$

$$I_{ij} = \int_0^1 \phi_i(x) \phi_j''(x) dx \quad \text{--- (6)}$$

$$Q_{ij} = \int_0^1 \phi_i(x) \phi_j(x) dx = - \int_0^1 \phi_i'(x) \phi_j'(x) dx \quad \text{--- (7)}$$

Now, putting $\epsilon = 0$, we get

$$I(\alpha) = \sum_{i=1}^n \sum_{j=1}^n \alpha_i \alpha_j q_{ij} + 2 \sum_{i=1}^n \alpha_i \alpha_j r_{ij} + 2 \sum_{i=1}^n \alpha_i p_i$$

for minimum $\frac{\partial I}{\partial \alpha_i} = 0$, which gives

$$2 \sum_{j=1}^n \alpha_j q_{ij} + 2 \sum_{j=1}^n \alpha_j r_{ij} + 2 p_i = 0$$

$$\sum_{j=1}^n \alpha_j (q_{ij} + r_{ij}) + p_i = 0 \quad \text{--- (8)}$$

to obtain approximate solⁿ, then, we take $n=2$, then by eqⁿ (8)

$$\begin{cases} \alpha_1 (q_{11} + r_{11}) + p_1 = 0 \\ \alpha_1 (q_{12} + r_{12}) + \alpha_2 (q_{22} + r_{22}) = -p_2 \\ \alpha_1 (q_{21} + r_{21}) + \alpha_2 (q_{22} + r_{22}) = -p_2 \end{cases}$$

(Here $q_{21} = q_{12}$ by symmetry)

let $\phi_1(x) = x(1-x)$ & $\phi_2(x) = x^2(1-x)$ be our ϕ_i , then we have

$$p_1 = \int_0^1 x \phi_1(x) dx = \int_0^1 x^2(1-x) dx = \left[\frac{x^3}{3} - \frac{x^4}{4} \right]_0^1 = \left[\frac{1}{3} - \frac{1}{4} \right] = \frac{1}{12}$$

$$p_2 = \int_0^1 x \phi_2(x) dx = \int_0^1 x^4(1-x) dx = \left[\frac{x^5}{5} - \frac{x^6}{6} \right]_0^1 = \frac{1}{20}$$

$$q_{11} = - \int_0^1 \phi_1'(x) \phi_1'(x) dx = - \int_0^1 (1-2x)^2 dx = - \left[\frac{(1-2x)^3}{3(1-2)} \right]_0^1 = \frac{1}{6} [-1 - 1] = -\frac{1}{3}$$

$$q_{21} = q_{12} = - \int_0^1 \phi_1'(x) \phi_2'(x) dx = - \int_0^1 (1-2x)(2x-3x^2) dx = - \int_0^1 (2x - 4x^2 - 3x^3 + 6x^3) dx = - \left[x^2 - \frac{4x^3}{3} + \frac{3x^4}{2} \right]_0^1 = - \left[1 - \frac{4}{3} + \frac{3}{2} \right] = - \left[\frac{6 - 14 + 9}{6} \right] = -\frac{1}{6}$$

$$q_{22} = - \int_0^1 \phi_2'(x) \phi_2'(x) dx = - \int_0^1 (2x-3x^2)^2 dx = - \int_0^1 (4x^2 - 12x^3 + 9x^4) dx = - \left[\frac{4x^3}{3} - \frac{12x^4}{4} + \frac{9x^5}{5} \right]_0^1 = - \left[\frac{4}{3} - 3 + \frac{9}{5} \right] = - \left[\frac{20 - 45 + 36}{15} \right] = -\frac{2}{15}$$

$$\begin{cases} \alpha_{11} = 1/30 \\ \alpha_{12} = 1/60 \\ \alpha_{22} = 1/105 \end{cases} \begin{cases} q_{11} + r_{11} = -\frac{1}{3} + \frac{1}{30} \\ \phantom{q_{11} + r_{11}} = -\frac{9}{30} = -\frac{3}{10} \\ q_{12} + r_{12} = \frac{1}{60} + \frac{1}{6} \\ \phantom{q_{12} + r_{12}} = \frac{1}{60} + \frac{10}{60} = \frac{11}{60} \\ q_{22} + r_{22} = \frac{1}{105} + \frac{1}{20} \\ \phantom{q_{22} + r_{22}} = \frac{4}{420} + \frac{21}{420} = \frac{25}{420} = \frac{5}{84} \end{cases}$$

putting in (8), we get
now eqⁿ in form of α_i for $i=1,2$.

$$12d_1 + 9d_2 = 5 \times \frac{1}{2}$$

$$63d_1 + 52d_2 = 21 \times \frac{1}{2}$$

$$0 + \left(\frac{9}{2} + \frac{52}{7}\right)d = \frac{5}{2} + 3 = \frac{11}{2}$$

$$d_1 = \frac{11}{2} \times \frac{14}{(63+104)} = \frac{77}{167}$$

10-4

5

Solve $y'' - x = 0$, $y(0) = 0$ — (1)

by Ritz method $y'(1) = 1/2$

functional for this problem is given by

$$I(v) = \int_0^1 (v'^2 + 2vx) dx + v(1)$$

let

$$dv = d_1 \phi_1(x) + d_2 \phi_2(x)$$

$$v = d_1 x + d_2 x^2 \quad \text{--- (2)}$$

be an approximate soln so that

$v(x)$ satisfies the B.C

$$v(0) = 0$$

$$v'(1) = d_1 + 2d_2 = v'(1) = 1/2 = d_1 + d_2$$

Now putting in eqn (2), we have

$$I(v) = \int_0^1 (d_1 + 2d_2 x)^2 + 2(d_1 x + d_2 x^2) x dx + v(1)$$

Now for min!

$$\frac{\partial I}{\partial d_1} = 0 = \int_0^1 2(d_1 + 2d_2 x) \cdot 1 + 2x dx + 1$$

$$\frac{\partial I}{\partial d_2} = 0 = \int_0^1 2(d_1 + 2d_2 x) \cdot 2x + 2x(x^2) dx + 1$$

$$\Rightarrow \left[2d_1 x + 4d_2 \frac{x^2}{2} + \frac{2}{3} x^3 \right]_0^1 + \frac{1}{2} = 0$$

$$\Rightarrow 2(d_1 + d_2) + \frac{2}{3} + \frac{1}{2} = 0 \Rightarrow d_1 + d_2 = -\frac{5}{3+2}$$

$$4 \left[4 \frac{d_1 x^2}{2} + 8d_2 \frac{x^3}{3} + 2 \frac{x^4}{4} \right]_0^1 + 1 = 0$$

$$\Rightarrow 2d_1 + \frac{8}{3}d_2 + \frac{1}{2} + 1 = 0$$

$$\Rightarrow d_1 + \frac{4}{3}d_2 = -\frac{3}{2} = -\frac{3}{4}$$

$$\Rightarrow d_1 = -\frac{13}{12}, d_2 = \frac{1}{4}$$

\(\therefore\) Approximate soln is given by

$$y(x) = -\frac{13}{12}x + \frac{1}{4}x^2$$

Formula

$(1-x)^{-1} = 1+x+x^2+x^3+\dots$ $(1+x)^{-1} = 1-x+x^2-x^3+\dots$

(1) $\frac{1}{x} - \frac{(1/x)^3}{3} + \frac{(1/x)^5}{5} - \frac{(1/x)^7}{7} + \dots$
 $= \tan^{-1}(1/x)$

ie $\frac{p}{1} - \frac{p^3}{3} + \frac{p^5}{5} - \frac{p^7}{7} + \dots = \tan^{-1}(p)$

(2) $\int_0^x (x) = \frac{1-x^2}{2} + \frac{x^4}{2 \cdot 4} - \frac{x^6}{2 \cdot 4 \cdot 6} + \dots$

(3) $\left(1 + \frac{1}{p^2}\right)^{-1/2} = 1 - \frac{1}{2} \left(\frac{1}{p^2}\right) + \frac{1 \cdot 3}{2 \cdot 4} \left(\frac{1}{p^4}\right) - \frac{1 \cdot 3 \cdot 5}{2 \cdot 4 \cdot 6} \left(\frac{1}{p^6}\right) + \dots$

(4) $e^x = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \frac{x^4}{4!} + \dots$

(5) $\sin x = x - \frac{x^3}{3!} + \frac{x^5}{5!} - \frac{x^7}{7!} + \dots$

$\cos x = 1 - \frac{x^2}{2!} + \frac{x^4}{4!} - \frac{x^6}{6!} + \dots$

(6) $\tan^{-1} x = x - \frac{x^3}{3} + \frac{x^5}{5} - \frac{x^7}{7} + \dots$

(7) $\Gamma n = \int_0^{\infty} e^{-x} x^{n-1} dx$

Laplace transform of special function

(1) $H(t-a) = \frac{e^{-as}}{s} \quad (e^{-as} f(s))$

(2) $\text{Exp}(t) = \frac{\log(1+t)}{s}$

(3) $\text{Sine}, \text{Si}(t) = \frac{1}{s} \cdot \tan^{-1}(1/s)$

(4) $\text{Cosine}, \text{Ci}(t) = \frac{1}{s} \cdot \log(1+t)$

(5) $\text{erf} \sqrt{t} = \frac{1}{\sqrt{s}} \cdot \frac{1}{\sqrt{s^2+1}}$

(6) $\sin ft = \frac{\sqrt{\pi/2}}{s} (e^{-k/s})$

(7) $\frac{\log t}{\sqrt{t}} = \sqrt{\pi/4} \cdot e^{-1/4}$

(8) $\int_0^{\infty} (at) = \frac{1}{(-1)(s^2+a^2)}$

Inverse formula - 44

$F(s) \quad f(t) = L^{-1}[F(s)]$

1. 1 $\dots 1/s$

2. t $1/s^2$

3. $t^n = \frac{n!}{s^{n+1}} = \frac{\Gamma(n+1)}{s^{n+1}}, \quad s > 0$

4. $e^{at} = \frac{1}{s-a}, \quad s > a$

5. $t^n e^{at} = \frac{1}{(s-a)^{n+1}}$

6. $\sin at = \frac{a}{s^2+a^2}, \quad \text{Cos} at = \frac{s}{s^2+a^2}, \quad s > 0$

7. $\text{Sin} hat = \frac{a}{s^2-a^2}, \quad \text{Cos} hat = \frac{s}{s^2-a^2}, \quad s > |a|$

8. $e^{bt} \sin at = \frac{a}{(s-b)^2+a^2}$

9. $e^{bt} \text{Cos} at = \frac{s-b}{(s-b)^2+a^2}$

(1) $\frac{1}{s^2+a^2} \cdot \sin ax = \frac{-2}{2a} \text{Cos} ax$ for P.I of non-homog eqn

(10) $L[f(at)] = \frac{1}{a} f(s/a)$

(11) $L[e^{at} f(t)] = f(s-a)$ (shifting)

(12) $L[f(t)] = \int_0^T e^{-st} f(t) dt$
 $\frac{1}{1-e^{-sT}}$ when $f(t) = f(t+T)$ periodic function

(13) $L[f'(t)] = -s f(s) - f(0)$

(14) $L[f^n(t)] = s^n L[f(t)] - s^{n-1} f(0) - s^{n-2} f'(0) - \dots - f^{(n-1)}(0)$

(15) $L[t^n f(t)] = (-1)^n \frac{d^n}{ds^n} f(s)$

$L[t f(t)] = -\frac{d}{ds} f(s)$